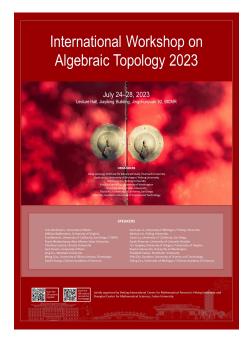
Moduli, moduli, moduli

Yifei Zhu

A moduli space is a space of parameters that label a certain family of structured objects we are interested in. I'll report on using methods of algebraic topology to understand aspects of a diverse set of moduli problems: (1, joint with Guozhen Wang et al.) in connection with p-adic arithmetic geometry, a filtered equivariant quasi-syntomic sheaf of Koszul complexes for computing unstable chromatic homotopy of spheres, over moduli spaces that parametrize deformations of a formal group with level structures; (2, joint with Hongwei Jia et al.) in connection with condensed-matter physics and materials science, monodromy of stratified vector bundles as moduli for gapless quantum mechanical systems, which arise from non-Hermitian symmetries; and (3, joint with Pingyao Feng et al.) in connection with data science, topological distribution spaces for image and speech signals, as revealed from persistent homology, and applied to the design of convolutional layers for deep learning. For each, I will introduce the context of study and describe the mathematical objects in question, with all technical terms above explained.



Moduli, moduli, moduli: Portraits of moduli spaces

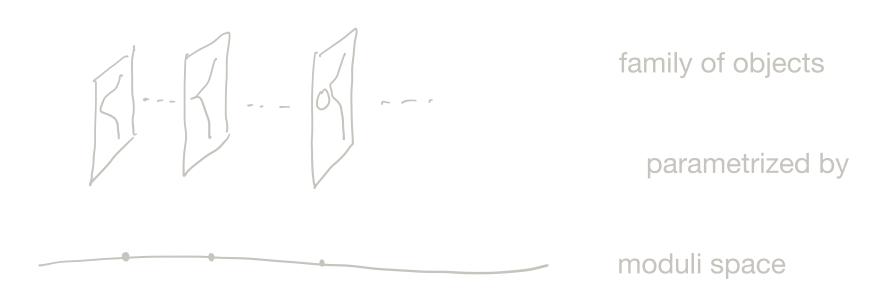


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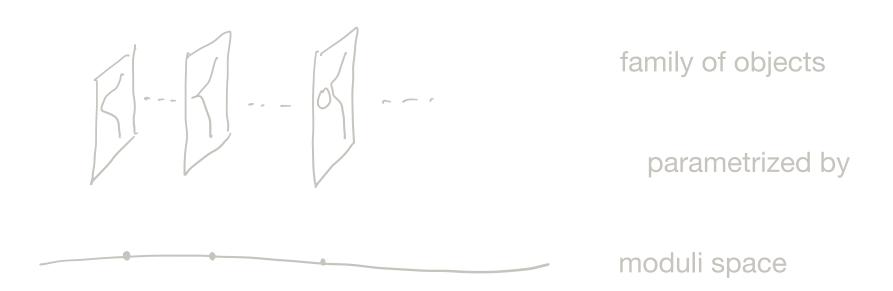
Southern University of Science and Technology

2023.7.25

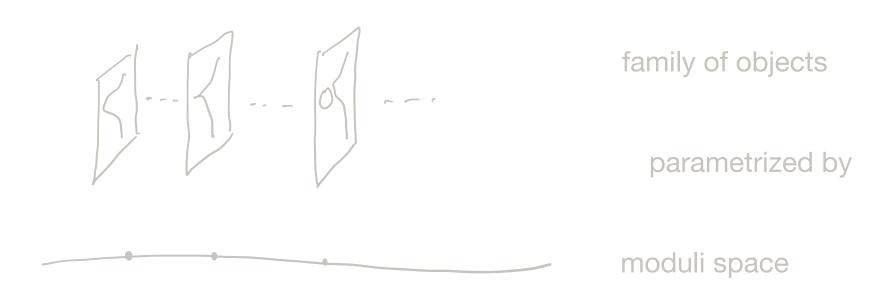
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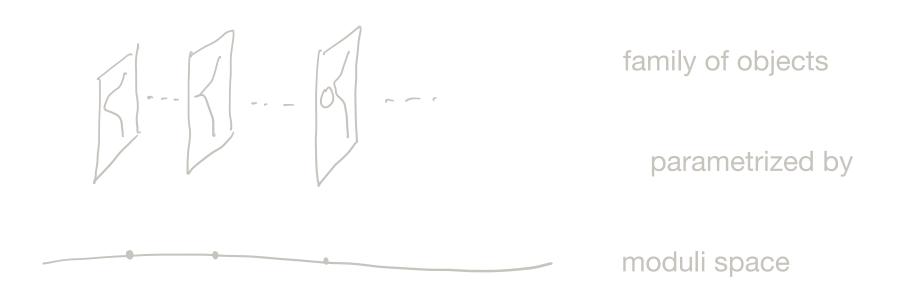
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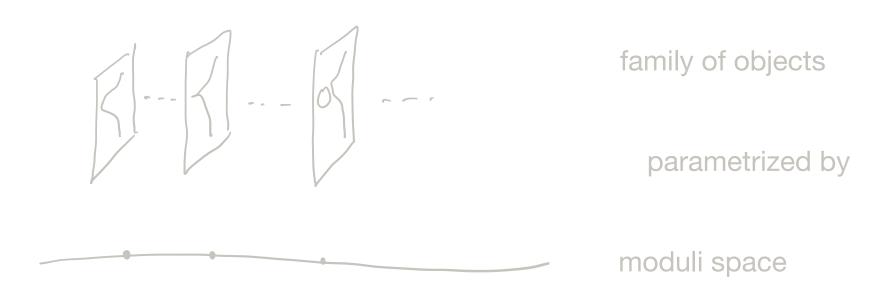
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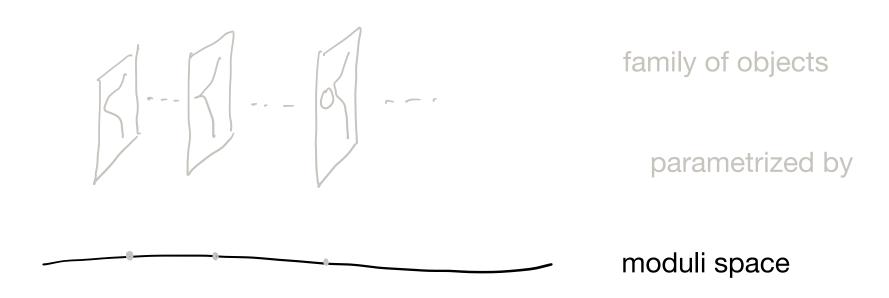
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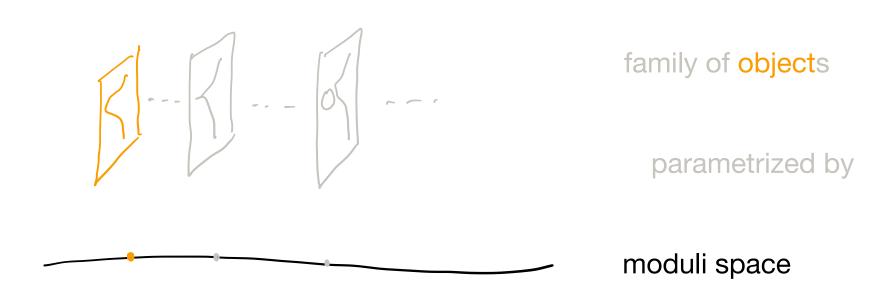
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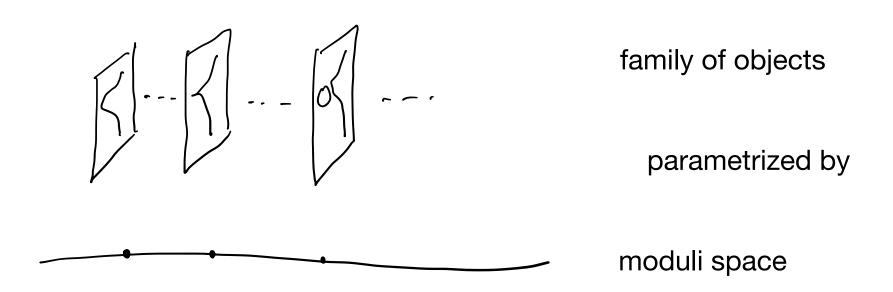
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- Instead of dealing with an isolated or static object, we would really like to understand a continuous family of objects, or how an object varies as the parameters on which it depends change.
- For these purposes, it is often fruitful to study this collection of parameters as a space on its own right.
- While the objects in question have structure, the parameter space also has its own structure, often rich and distinct from that of the objects.
- Understanding the moduli space offers in turn understanding of the objects individually and as a whole.
- In this sense, studying moduli spaces is of the second-order nature.

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First portraits: Algebro-geometric setup

• Objects to parametrize: Deformations of a formal group with level structure

G= formal group of height $h<\infty$ over a perfect field k of characteristic p>0

R = complete local ring with residue field k, nilpotent maximal ideal \mathfrak{m} , and natural projection $\pi \colon R \to R/\mathfrak{m}$

deformation of G/k to $R := (\mathbb{G}, i, \alpha)$ with \mathbb{G} a formal group over R, $i: k \hookrightarrow R/\mathfrak{m}$, and $\alpha: \pi^* \mathbb{G} \xrightarrow{\cong} i^* G$

deformation of G with a level- $\Gamma_0(p^n)$ structure := (G, H) with H a cyclic degree- p^n subgroup

 $= \psi \colon \mathbb{G} \to \mathbb{G}/\mathbb{H}$ over an extension of Rwhich lifts the relative Frobenius $\operatorname{Frob}^n \colon G \to G^{(p^n)}$

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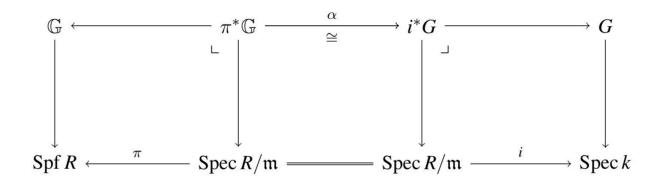
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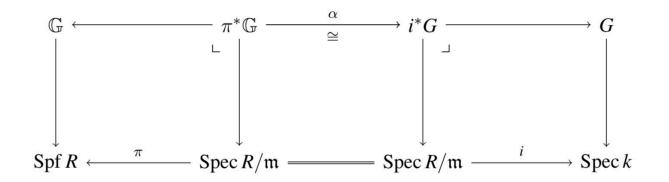


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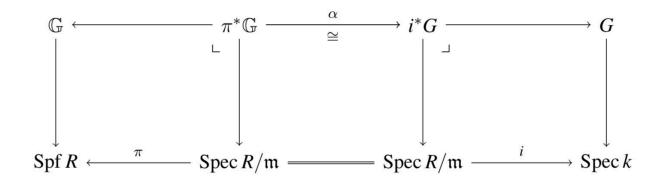


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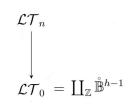
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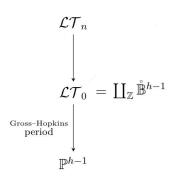
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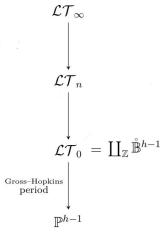
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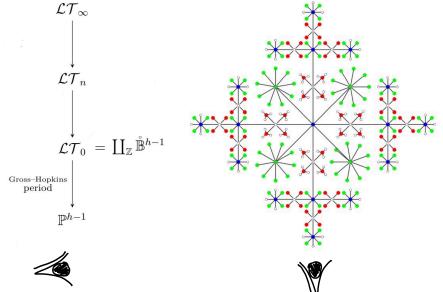
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Koszul complexes for computing unstable chromatic homotopy of spheres

[Devinatz–Hopkins '04, Wang '15] There is a homology-to-homotopy SS converging to the *vn*-periodic homotopy groups of the *q*-dimensional sphere

$$H^*_{\mathrm{c}}(\mathcal{G}_h; E^{\wedge}_* \Phi_h(S^q)) \implies \upsilon_n^{-1} \pi_* S^q$$

whose E_2 -page is the continuous group cohomology of the *h*'th Morava stabilizer group \mathcal{G}_h with coefficients in the completed *E*-homology of the Bousfield–Kuhn functor Φ_h applied to the *q*-sphere.

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Filtration (in the case of *h* = 2 for simplicity)

 Have a sequence of unstable spheres
 Applying *E*₀^ˆ Φ₂ (−), get a sequence of Koszul complexes

 $A_{0} = W(\overline{\mathbb{F}}_{p})\llbracket v_{1} \rrbracket \cong E^{0}(\mathrm{pt})$ $A_{1} = W(\overline{\mathbb{F}}_{p})[\alpha_{1}, \alpha_{1}']/(\alpha_{1}\alpha_{1}' - p)$ $\cong E^{0}(B\Sigma_{p})/I_{\mathrm{tr}}$ $E^{0} \xrightarrow{\psi} E^{0}(B\Sigma_{p})/I \xrightarrow{\psi} E^{0}(B(\Sigma_{p} \wr \Sigma_{p}))/I'$ $v_{1} \mapsto v_{1}' \qquad \alpha_{1} \mapsto \alpha_{1}'$

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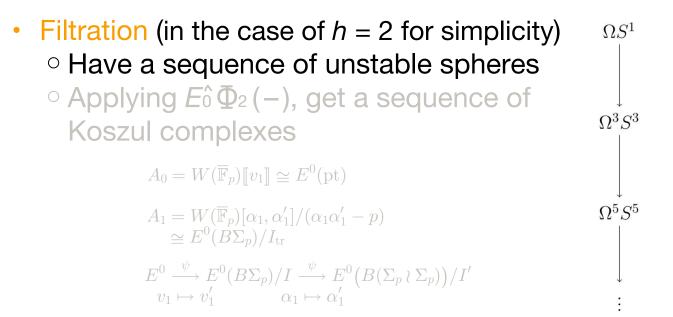
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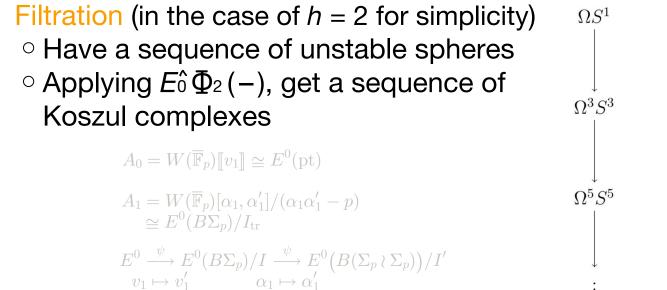
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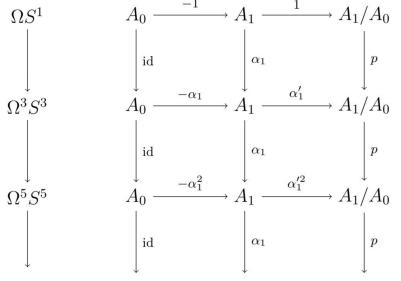
$$\cong E^{0}(B\Sigma_{p})/I_{tr}$$

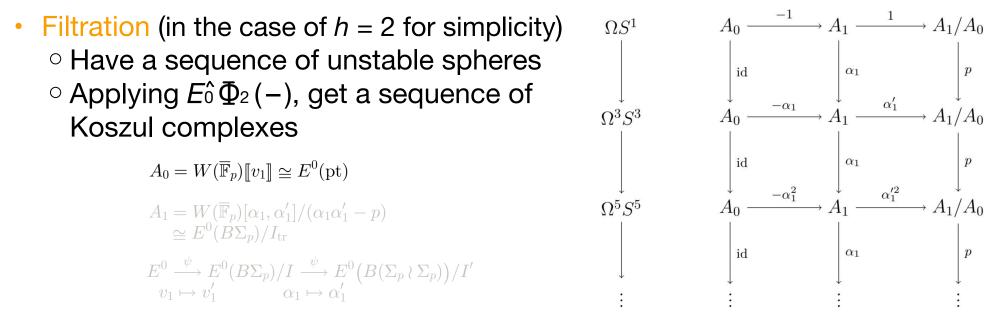
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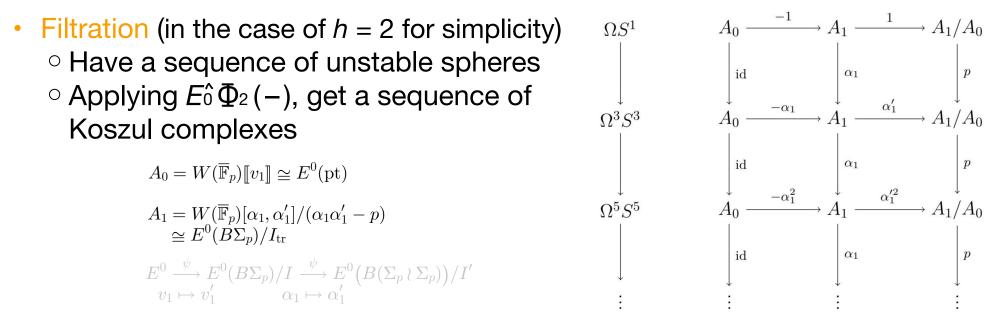
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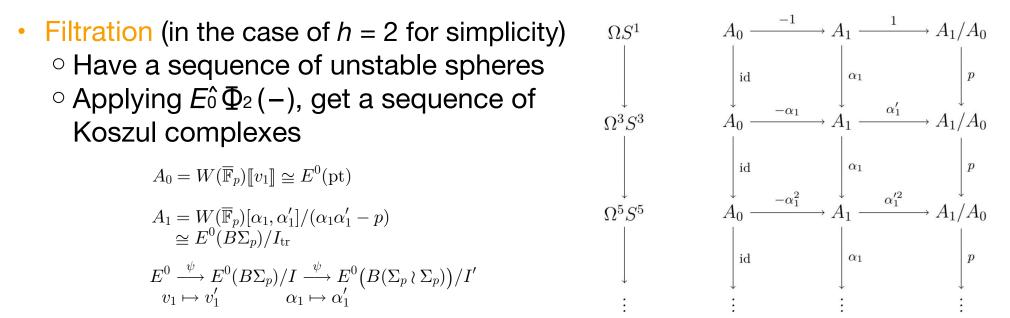


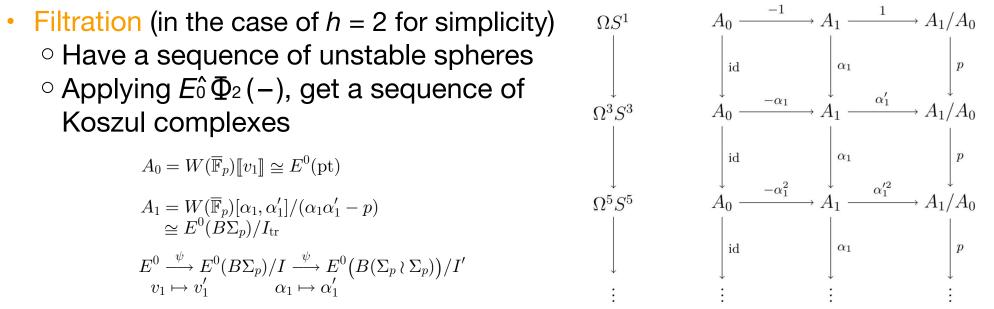


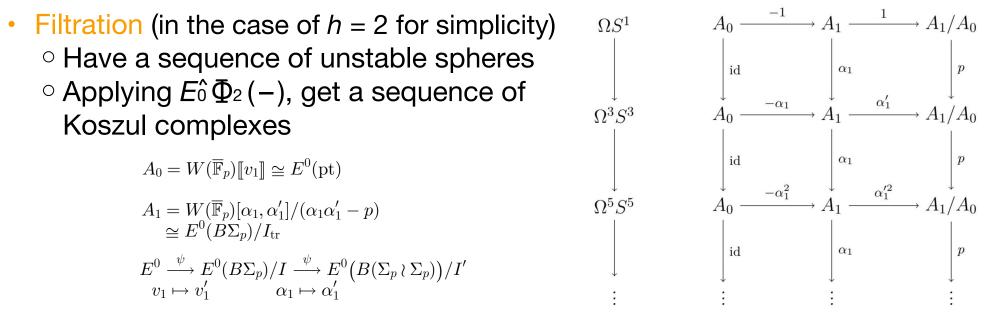


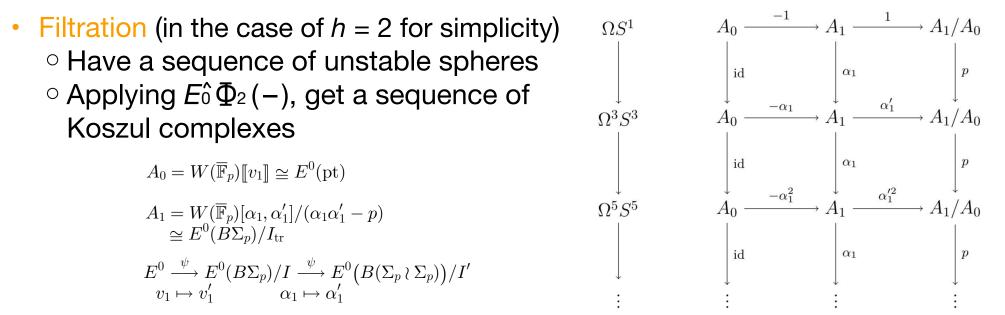


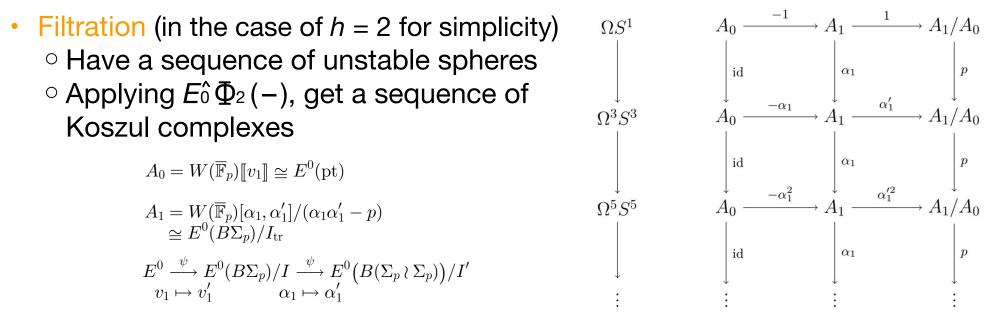












• More concretely, this forms a sheaf of functions over the moduli, filtered by orders *m* of pole at the irreducible component cut out by a_1 , which precisely correspond to dimensions q = 2m + 1 of the spheres.

<i>C</i> 4	<i>c</i> ₆	c_{4}^{2}	<i>C</i> ₄ <i>C</i> ₆	Δ , c_4^3	$c_{4}^{2}c_{6}$	$c_4\Delta$, c_4^4	$c_6\Delta$, $c_4^3c_6$
<i>a</i> ₁ <i>a</i> ₃	9 <i>a</i> ₃ ²	<i>a</i> ₁ <i>a</i> ₃ <i>c</i> ₄	$9a_3^2c_4$	$a_1 a_3 c_4^2$	$9a_{3}^{2}c_{4}^{2}$	$a_1 a_3 c_4^3$	$9a_3^2c_4^3$
x_{0}^{2}	3 <i>a</i> ² ₃	$a_1^2 a_3^2$	$3a_3^2c_4$	$a_1^2 a_3^2 c_4$	$3a_3^2c_4^2$	$a_1^2 a_3^2 c_4^2$	$3a_3^2c_4^3$
	a_3^2	$a_2 x_0^3 - 2a_4 x_0^2$	$a_{3}^{2}c_{4}$	$27a_3^4 \sim a_3^2c_6$	$a_3^2 c_4^2$	$a_3^2 c_6 c_4$	$a_{3}^{2}c_{4}^{3}$
		x ₀ ⁴	$a_1a_3^3(?)$	9a ₃ ⁴	$a_1 a_3^3 c_4$	$9a_{3}^{4}c_{4}$	$a_1 a_3^3 c_4^2$
			x_{0}^{5}	3a ₃ ⁴	$a_1^2 a_3^4$	$3a_3^4c_4$	$a_1^2 a_3^4 c_4$
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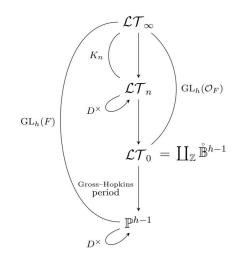
Part of calculations at p = 3

- The sequence of Koszul complexes is equivariant with respect to the action of the Morava stabilizer group $\mathcal{G}_{h} \cong \mathcal{O}_{D}^{\times}$ (with D / F = cent. div. alg. of inv. 1/h).
- Both $GL_h(F)$ (its congruence subgroups K_n) and D^* act on the Lubin–Tate tower, realizing the Jacquet–Langlands correspondence.
- [Faltings, Fargues '08] There is an equivariant isomorphism between the Lubin–Tate tower and another Drinfeld tower (parametrizing shtukas).

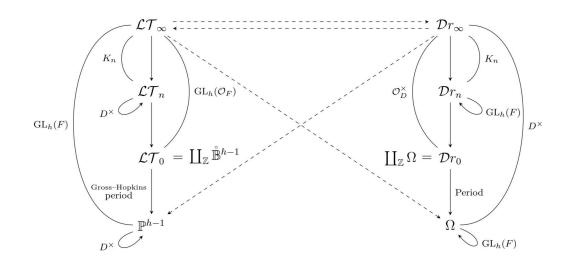
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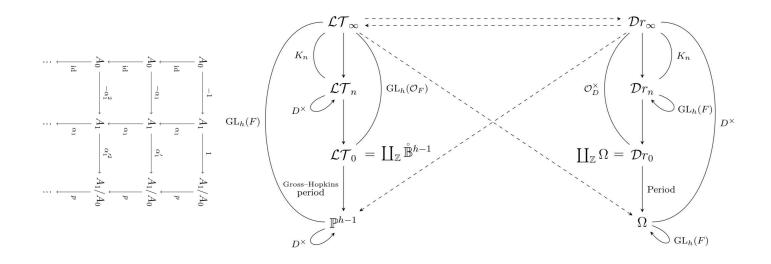
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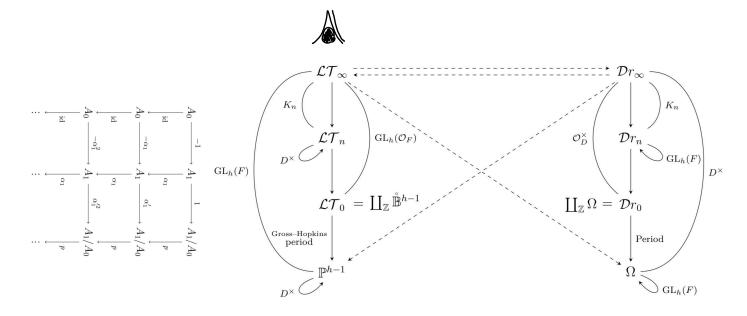
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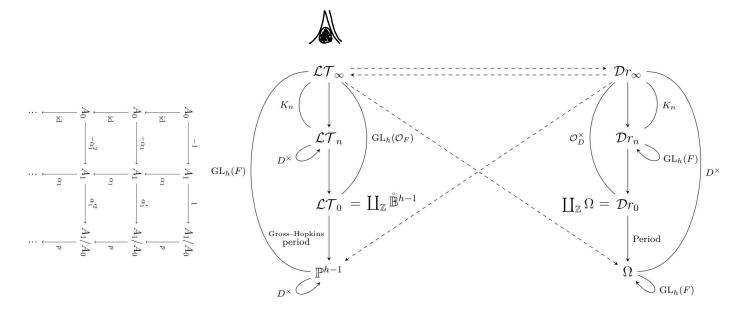
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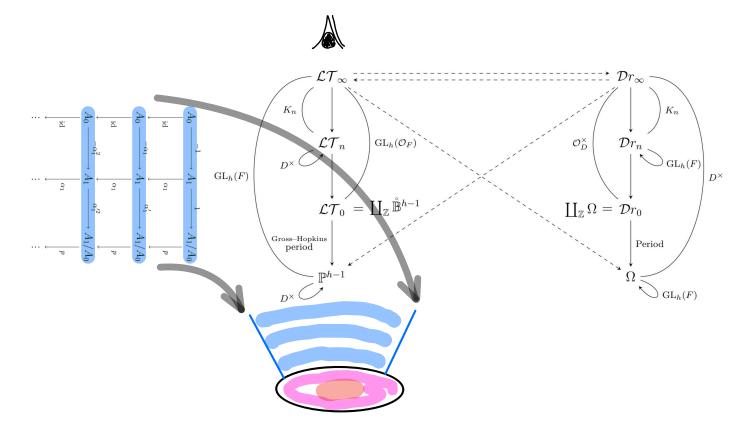


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Topological classifications for physical systems

Understanding continuous evolution of physical systems at the micro or quantum scale has a real impact on the larger-scale properties of materials. For example, holography is made possible via exceptional optical devices.

Moduli spaces of physical systems, especially their singular loci, play a pivotal role in designing such. Topological classifications enable physicists to fine-tune and create materials that can "do wonders" and cannot be found in nature, e.g., making invisible cloaks and other absorption devices.

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- Objects to parametrize: Hamiltonians with symmetries
 - Physical systems root in symmetries. For example, quantum mechanical systems can be described by their Hamiltonians, whose mathematical bearings are conventionally Hermitian matrices. Here, Hermiticity guarantees that the eigenvalues are real, corresponding to the fact that energies of the systems are observed to be real.
 - More recently, physicists have begun to model open systems by relaxing the Hermitian symmetry to allow eigenvalues with a nonzero imaginary part. This imaginary part measures energy exchange between the system and its surrounding environment. Still, some sorts of symmetry need to be imposed on the matrices to make them physically meaningful.
 - The size of the matrices corresponds to the number of energy band gaps.
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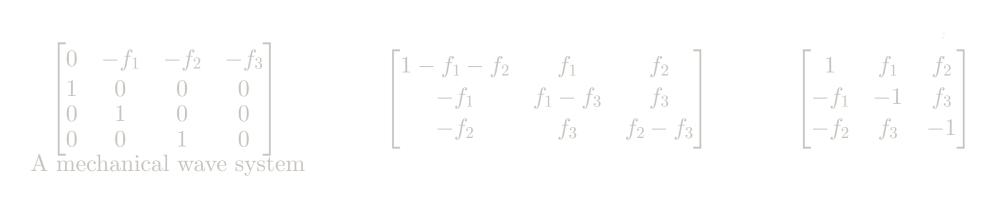
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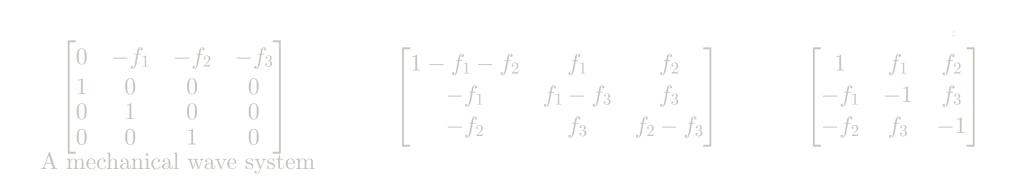
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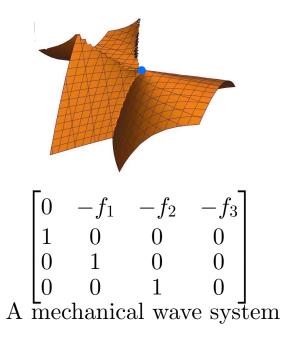
• Portraits of base spaces for gapless systems, the central figure within these configurations being the "swallowtail catastrophe" [Thom, Arnold]:



 $\begin{bmatrix} f_1 f_2 & f_1 & f_2 \\ -f_1 & f_1 & f_3 \\ -f_2 & f_3 & f_2 \end{bmatrix}$

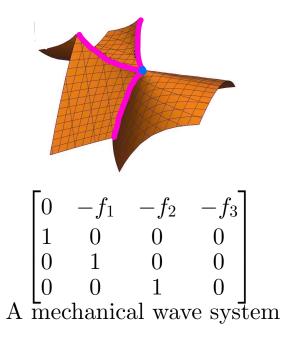


$$\begin{bmatrix} f_1 f_2 & f_1 & f_2 \\ -f_1 & f_1 & f_3 \\ -f_2 & f_3 & f_2 \end{bmatrix}$$



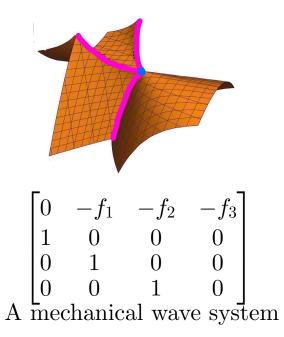
$$\begin{bmatrix} 1 - f_1 - f_2 & f_1 & f_2 \\ -f_1 & f_1 - f_3 & f_3 \\ -f_2 & f_3 & f_2 - f_3 \end{bmatrix} \begin{bmatrix} 1 & f_1 & f_2 \\ -f_1 & -1 & f_3 \\ -f_2 & f_3 & -1 \end{bmatrix}$$

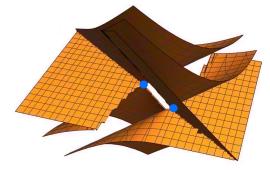
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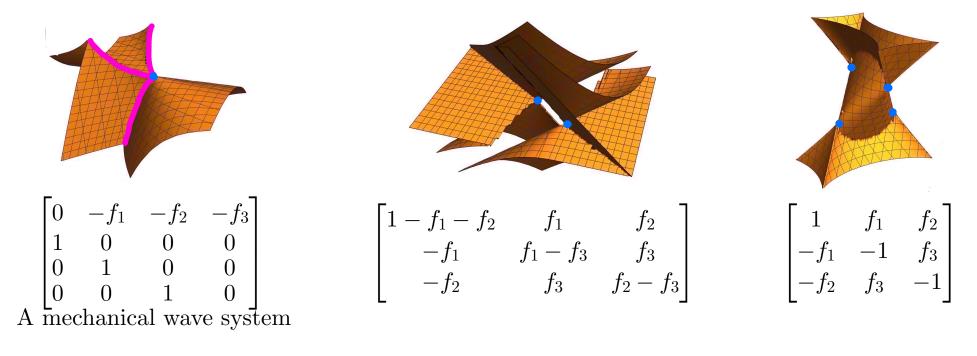
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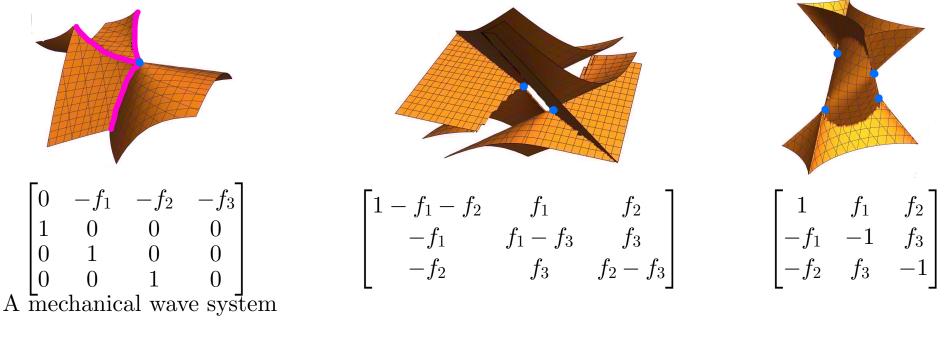
$[1 - f_1 - f_2]$	f_1	f_2	Γ	1	f_1	f_2
$-f_1$	$f_1 - f_3$	f_3		$-f_1$	-1	f_3
$-f_2$	f_3	$f_2 - f_3$	L	$-f_{2}$	f_3	$ \begin{bmatrix} f_2 \\ f_3 \\ -1 \end{bmatrix} $

$$\begin{bmatrix} f_1 f_2 & f_1 & f_2 \\ -f_1 & f_1 & f_3 \\ -f_2 & f_3 & f_2 \end{bmatrix}$$

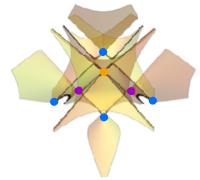


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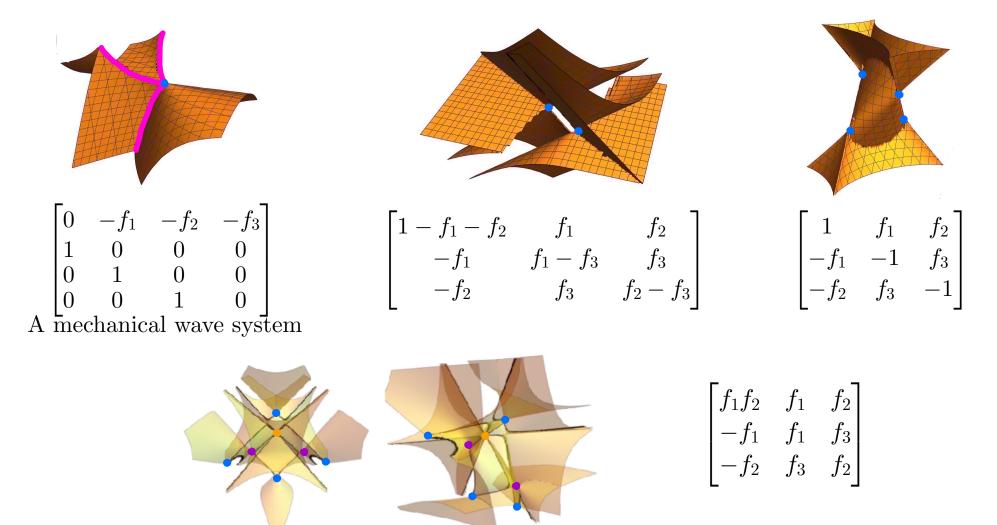
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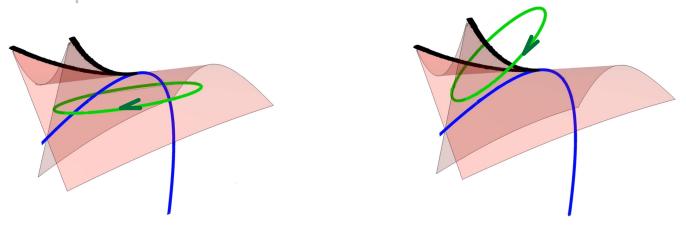
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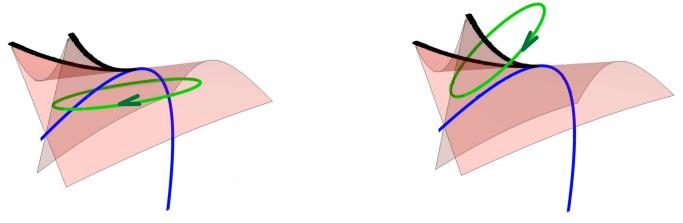
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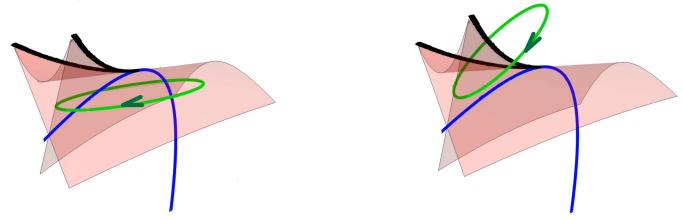
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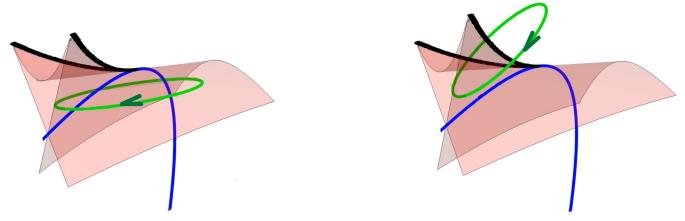
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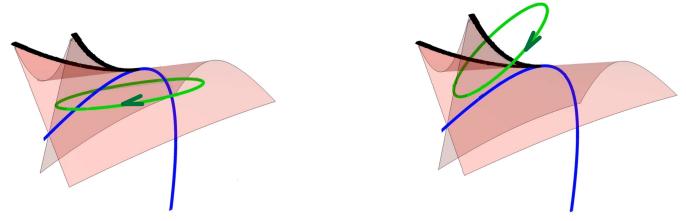
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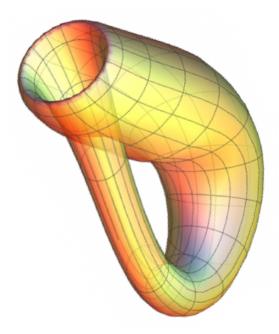
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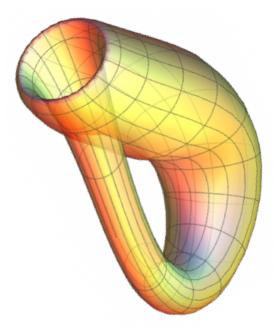
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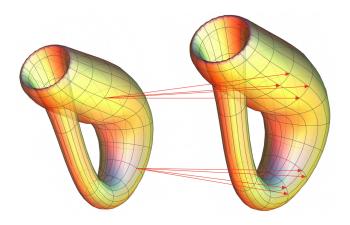
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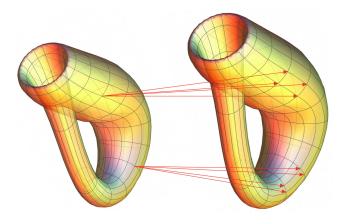
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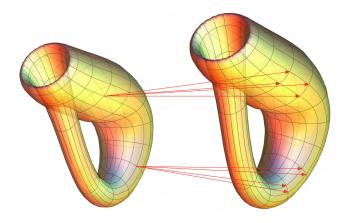
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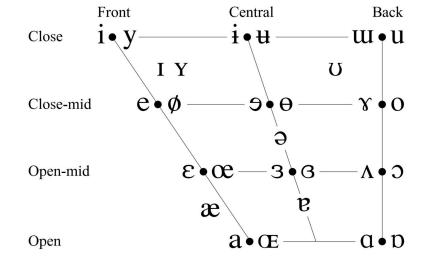
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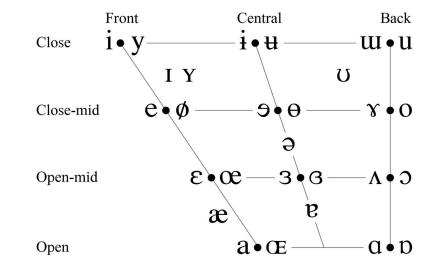
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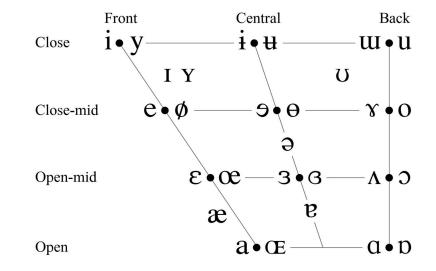


The vertical axis of the chart denotes vowel height. Vowels pronounced with the tongue lowered are at the bottom and raised are at the top. The horizontal axis of the chart denotes vowel backness. Vowels with the tongue moved towards the front of the mouth are in the left of the chart, while those with the tongue moved to the back are placed in right. The last parameter is whether the lips are rounded. At each given spot, vowels on the right and left are rounded and unnounded, respectively.

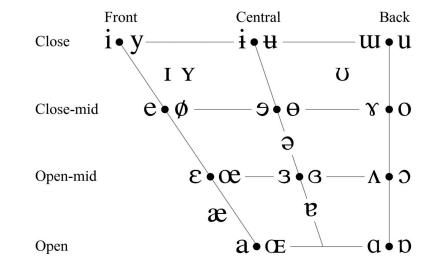
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Thank you.



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