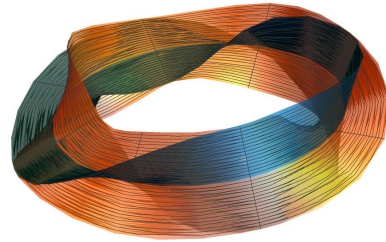


Explicit examples of Higgs bundles from physics and bulk–edge correspondence



Yifei Zhu

Southern University of Science and Technology

2024.7.2

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- With hindsight, how **explicit examples of rank-2 and rank-3 Higgs bundles** arise in topological classifications of (gapless) quantum mechanical systems

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*Holography, optical devices,
absorption devices, ...*

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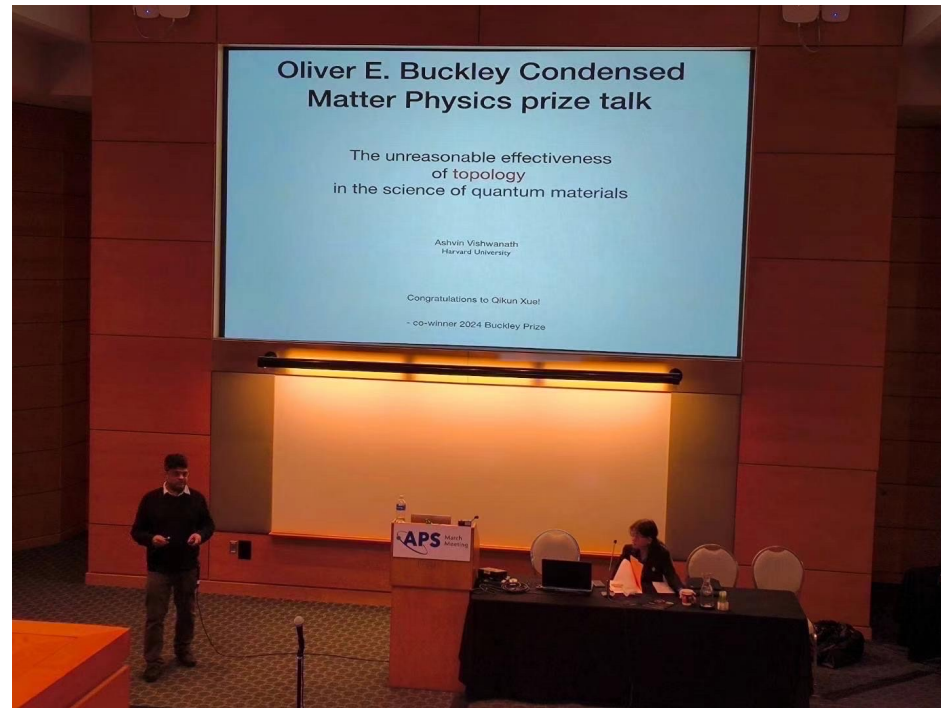
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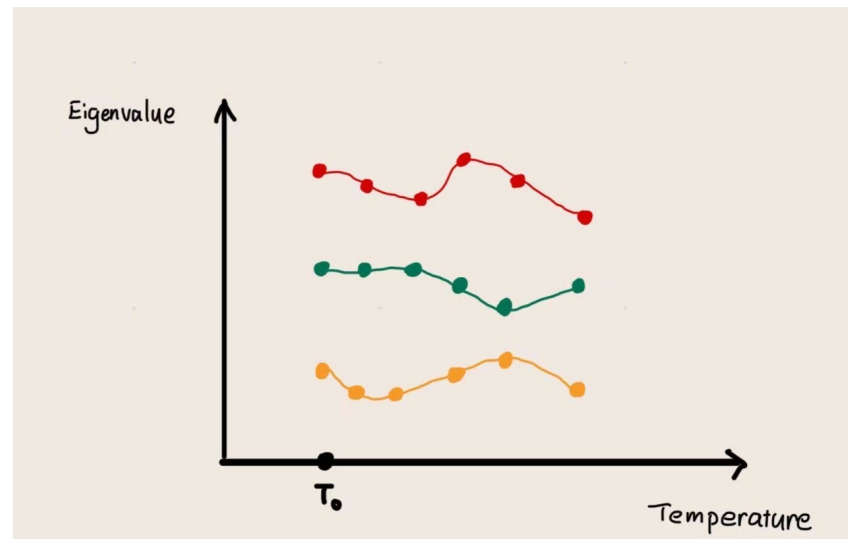
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- *The unreasonable effectiveness of topology in the science of quantum materials*, Ashvin Vishwanath of Harvard delivering the Buckley prize talk at this year's APS March Meeting in Minneapolis (薛其坤 of Tsinghua and SUSTech was the co-winner of the prize)
- U.S. Department of Energy, Office of Science. *Basic research needs for quantum materials: Research to discover, harness, and exploit exotic electronic properties* (brochure), 2016.
- 方忠 等, “**拓扑电子材料计算预测**”, 2023年度国家自然科学基金一等奖
- 第一届魅丽数学与交叉应用会议“数学与生物医药、**数学与先进材料**”, 2024年5月, 苏州

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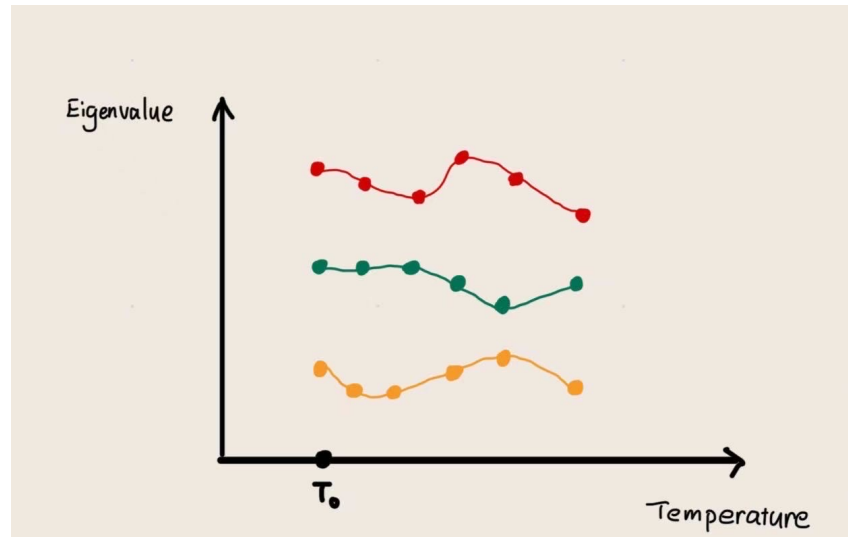
Mathematical modeling of electronic energy *band structures* therein



energy \longleftrightarrow eigenvalue

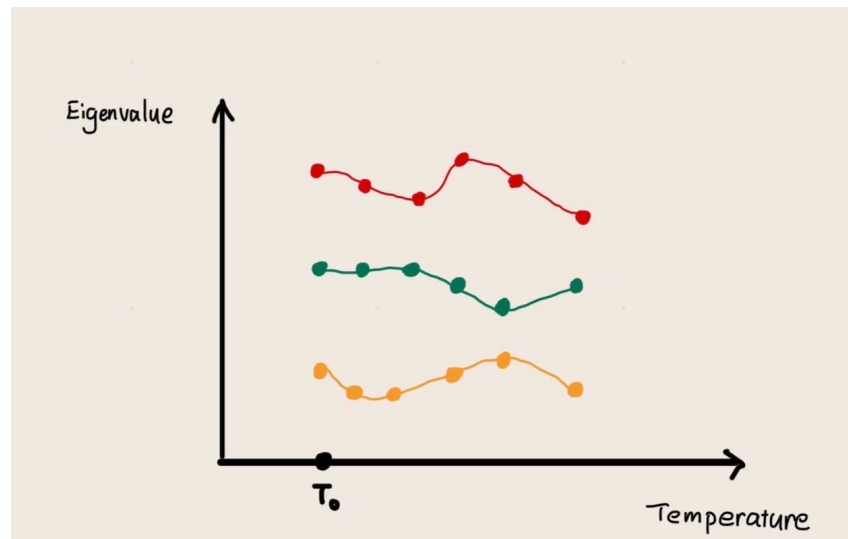
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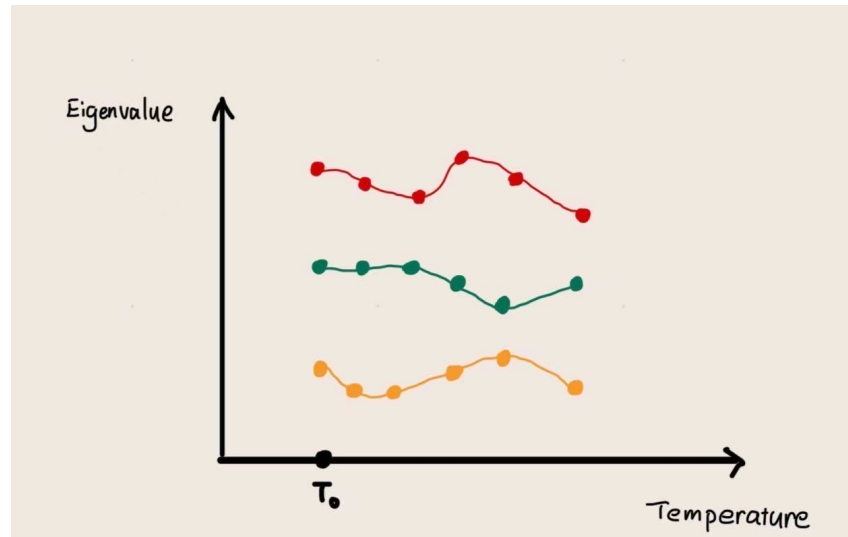
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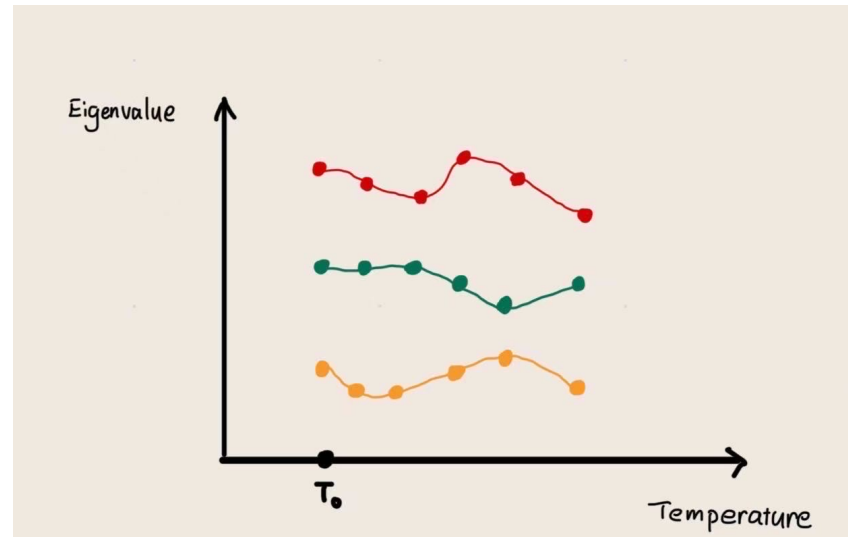
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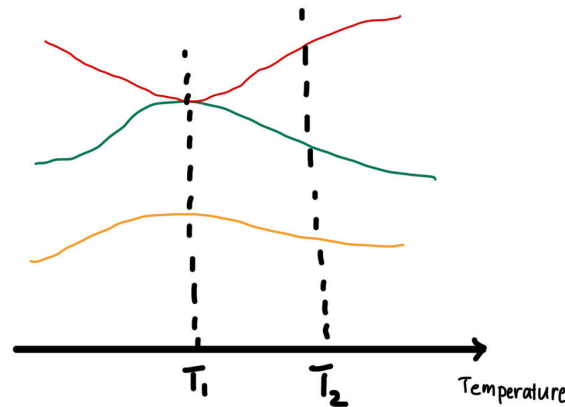


*Hermitian vs.
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*real eigenvalues
(observable
energies) vs.
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imaginary part
(counts for
energy exchange
with surrounding
environment or
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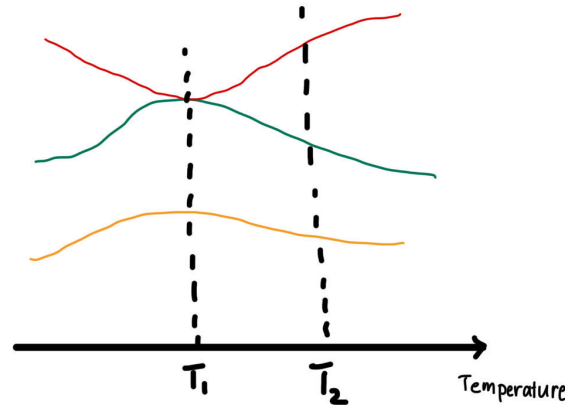


phases of matter
 H_2O : solid, liquid, gas

T_1 : singular points (points where eigenvalues degenerate)
 $H(T_1)$: gapless Hamiltonian
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Guzman et al., Model-free characterization of topological edge and corner states in mechanical networks, PNAS 2024.

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Thanks to Hopf bundles and Higgs bundles as *eigenbundles*, we now have a **conceptually more systematic**, visibly more intuitive understanding of the topic.

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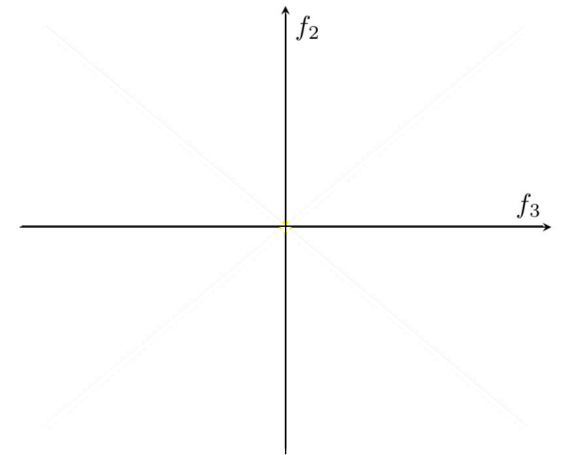
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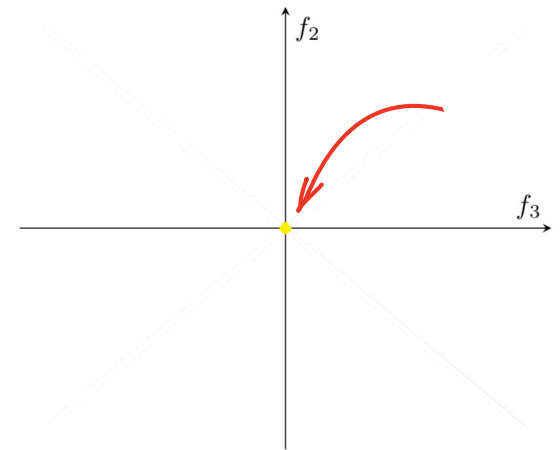
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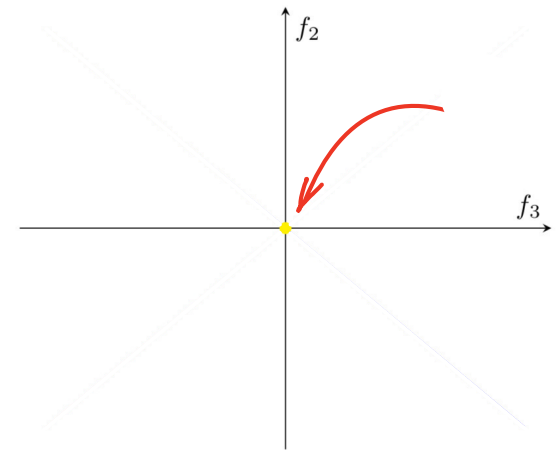
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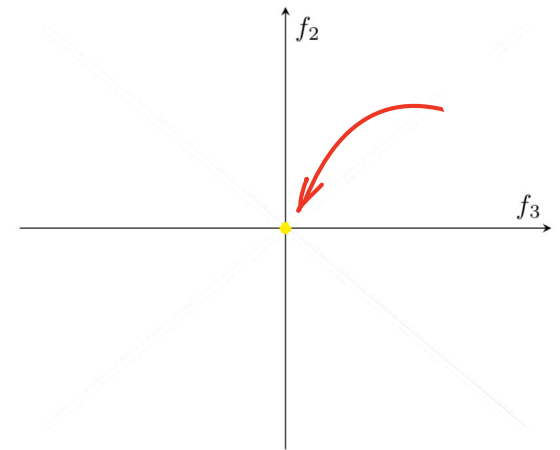
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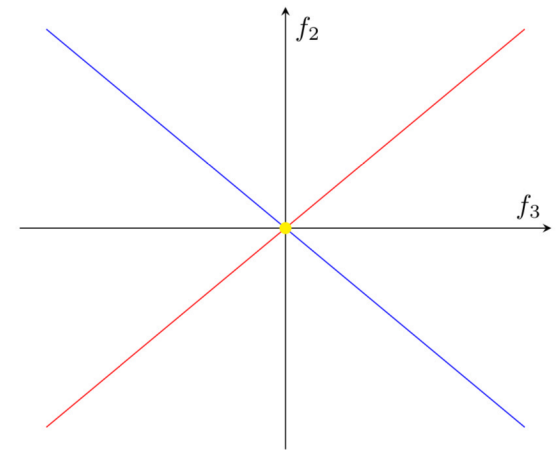
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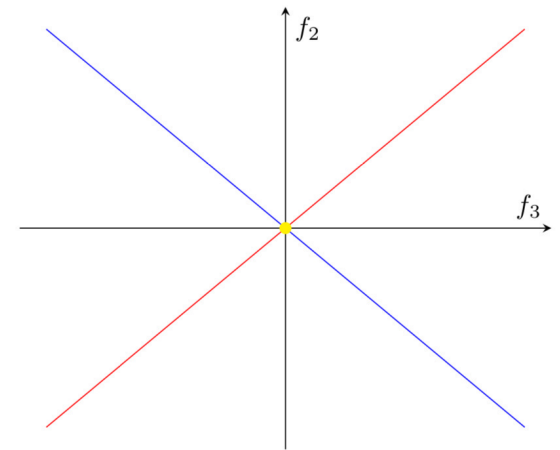
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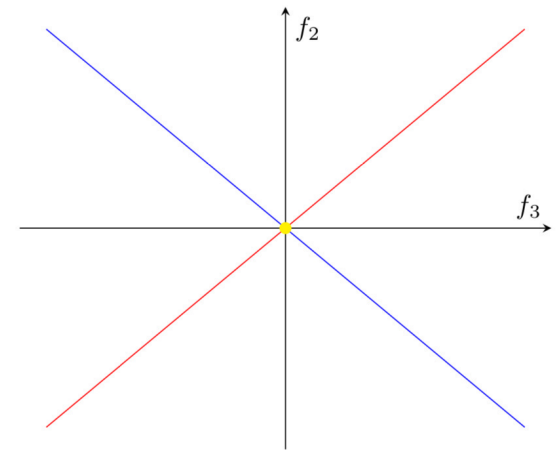
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1. Over $\{f_2 = \pm f_3\} - \{(0, 0)\}$, again H has a double eigenvalue, but its eigenspace is of **dimension 1**.



Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

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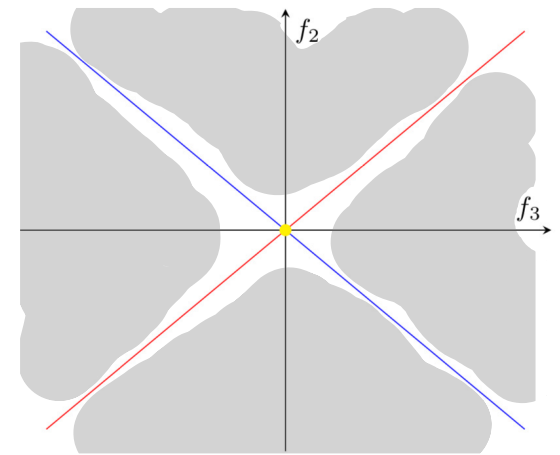
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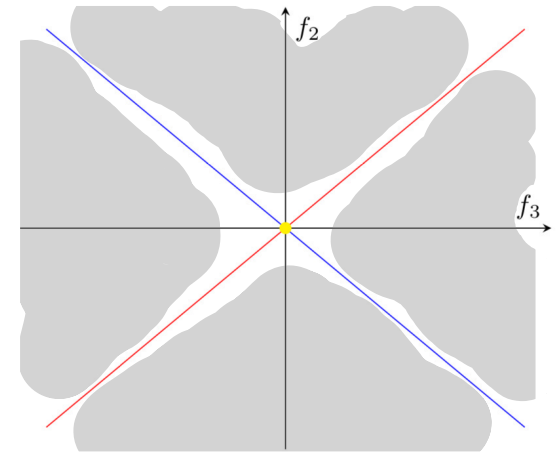
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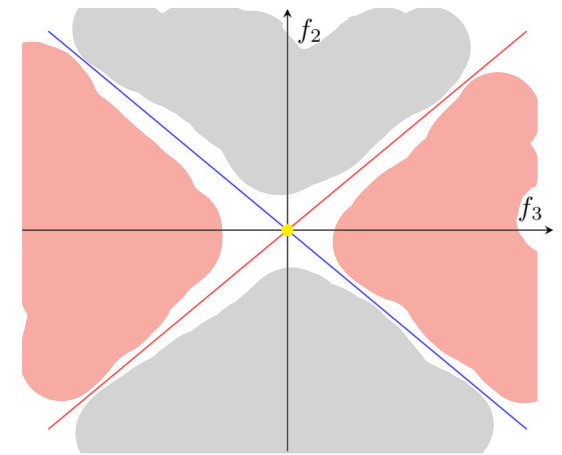
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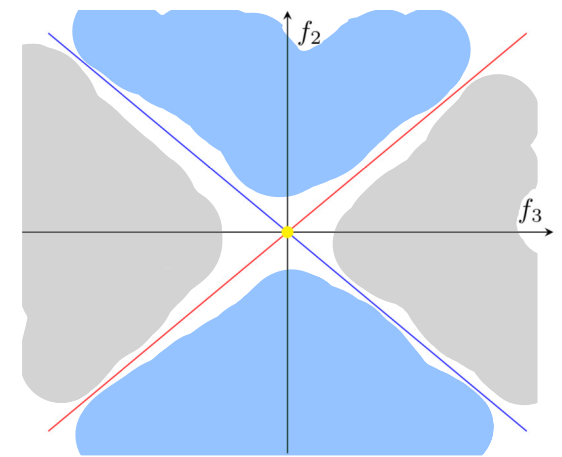
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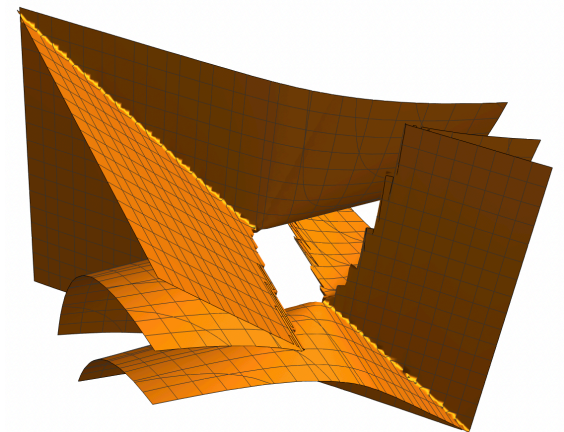
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The equation for this surface is a non-homogeneous real polynomial in f_1, f_2, f_3 of degree 6.



Swallowtail couple sw2

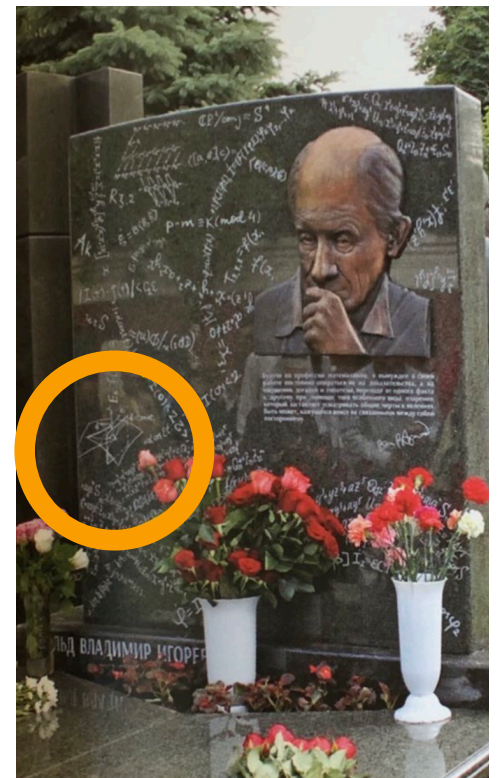
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V. I. Arnold's tombstone at the Novodevichy Cemetery in Moscow

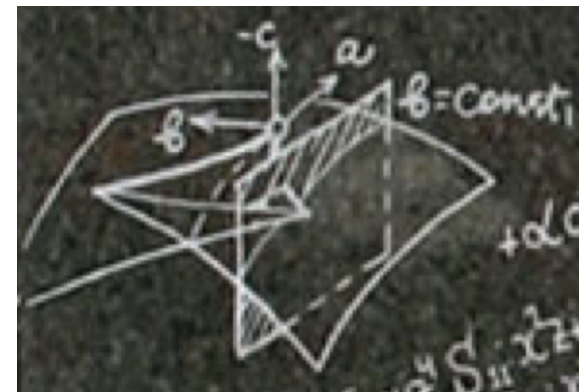
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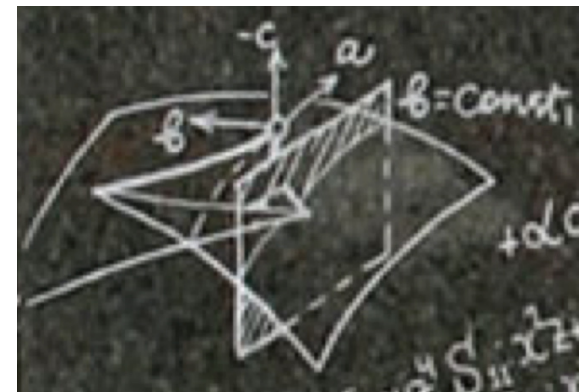
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A **local** model for moduli spaces of 3-band Hamiltonians

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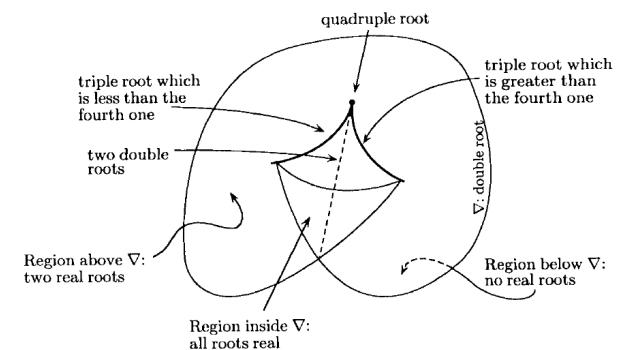
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Arnold, Braids of algebraic functions and the cohomology of swallowtails, 1968.

Homological stability of braid groups

Portrait from Gelfand, Kapranov, Zelevinsky, Discriminants, resultants, and multidimensional determinants.



The space of polynomials $x^4 + ax^2 + bx + c$

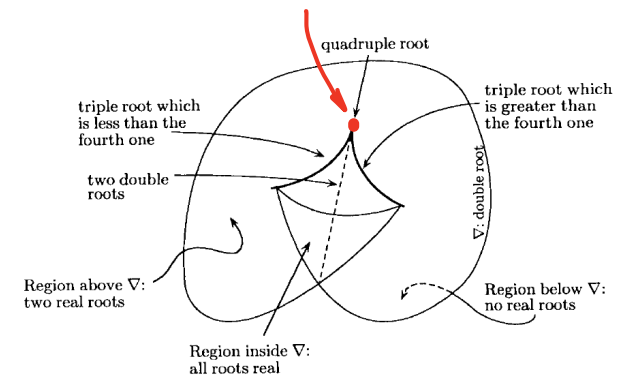
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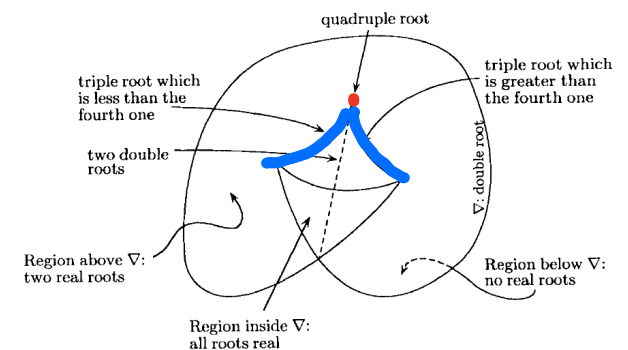
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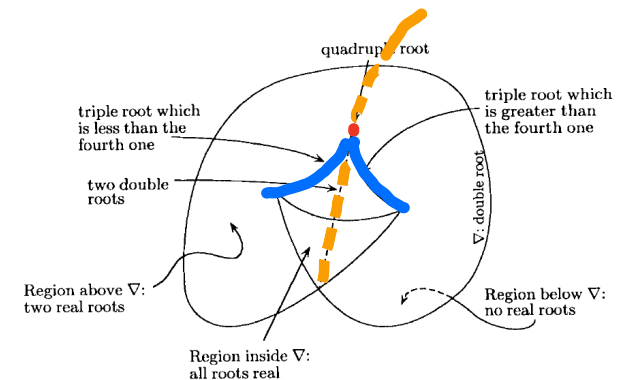
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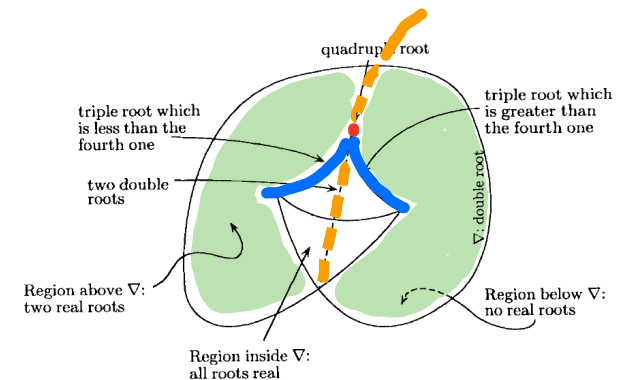
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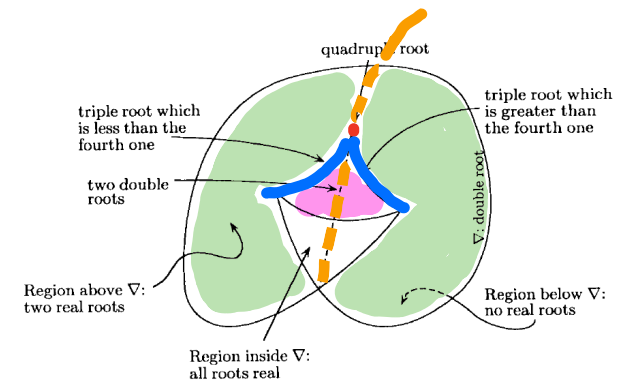
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Again, we aim to find computable algebraic invariants that systematically classify the evolutions of eigenvectors along loops in such stratified parameter spaces, including when they cross the discriminant surface resulting in degeneracies of various sorts.



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Eigenframe rotation as vector bundles: Revisiting the Hermitian case

The Hermitian case is simple, as the singularity is **isolated**, yet has profound physical implications already known to Arnold.

*Remarks on eigenvalues and eigenvectors of Hermitian matrices,
Berry phase, adiabatic connections and quantum Hall effect, 1995.*

Also: Polymathematics, 2000.

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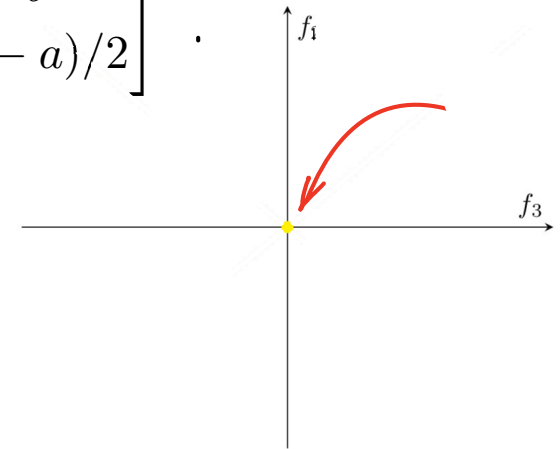
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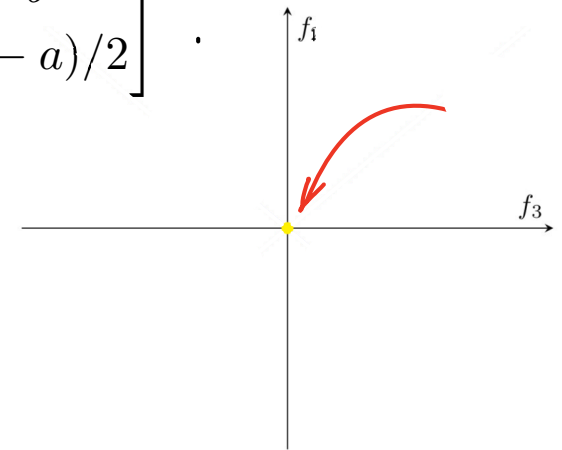
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Eigenframe rotation as vector bundles: Revisiting the Hermitian case

How does the eigenframe rotate over this stratified parameter plane?

As our starting point, previous work of Wu et al. [**Science**, 2019] classified the eigenframe rotation by

$$\pi_1(SO(2)/O(1)) \cong \pi_1(S^1) \cong \mathbb{Z}$$

Moreover, they obtained non-Abelian “topological charge” for n -band Hermitian systems when $n > 2$, such as

$$\pi_1(SO(3)/(O(1) \times O(1))) \cong \mathbb{H} = \{\pm 1, \pm i, \pm j, \pm k\}$$

However, their explanation for the appearance of the $O(1)^n$ -action in constructing the moduli spaces was rather ad hoc.

One of our key steps is a more conceptual understanding of the above moduli spaces in the case of $n=2$ through **bundle theory**. To see how they rotate, let us compute the unit eigenvectors explicitly.

$$\begin{vmatrix} f_3 - \omega & f_1 \\ f_1 & -f_3 - \omega \end{vmatrix} = \omega^2 - f_1^2 - f_3^2 = 0 \implies \omega_{\pm} = \pm \sqrt{f_1^2 + f_3^2}$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

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Eigenframe rotation as vector bundles: Revisiting the Hermitian case

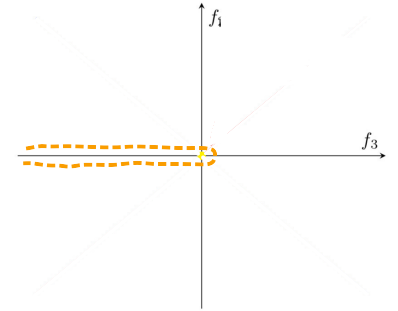
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Eigenframe rotation as vector bundles: Revisiting the Hermitian case

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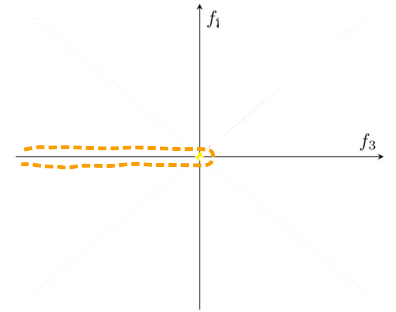
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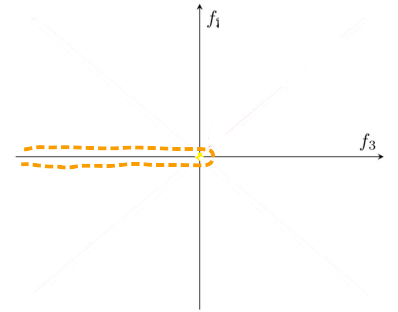
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Eigenframe rotation as vector bundles: Revisiting the Hermitian case

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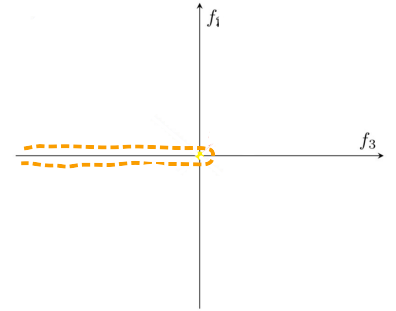
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Eigenframe rotation as vector bundles: Revisiting the Hermitian case

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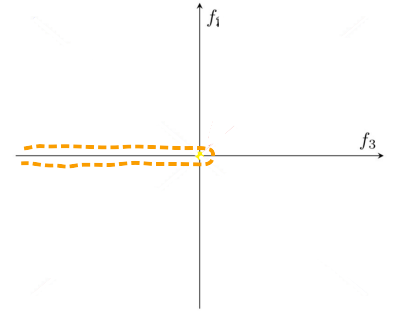
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Eigenframe rotation as vector bundles: Revisiting the Hermitian case

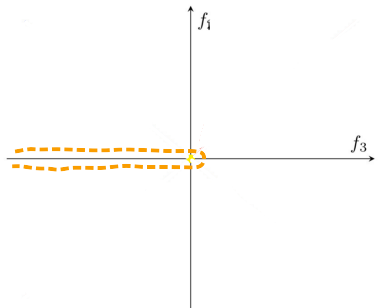
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Eigenframe rotation as vector bundles: Revisiting the Hermitian case

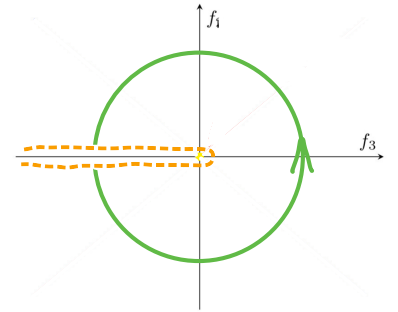
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Eigenframe rotation as vector bundles: Revisiting the Hermitian case

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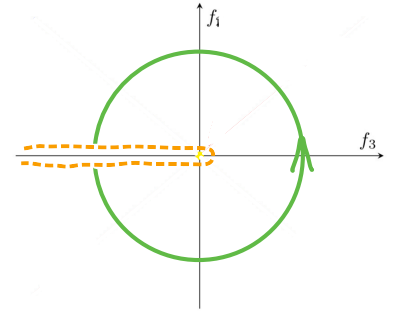
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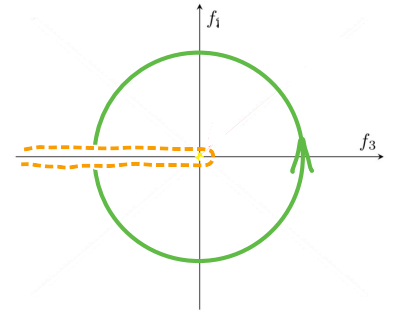


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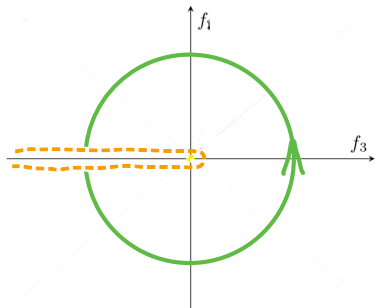
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Eigenframe rotation as vector bundles: Revisiting the Hermitian case

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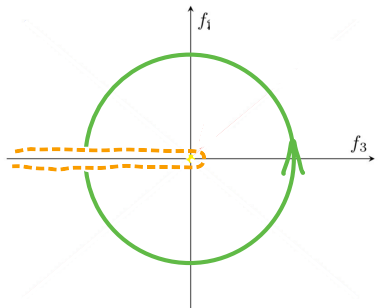
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Eigenframe rotation as vector bundles: Revisiting the Hermitian case

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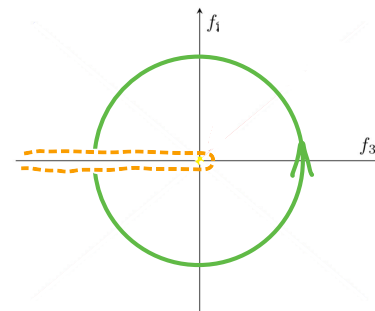
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Eigenframe rotation as vector bundles: Revisiting the Hermitian case

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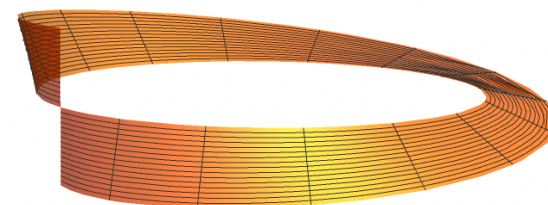
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Eigenframe rotation as vector bundles: Revisiting the Hermitian case

Lemma. The universal rank-1 eigenbundle for Hermitian 2-band systems is given by the **Hopf bundle**

$$\begin{array}{lll} S^0 \hookrightarrow S^1 \rightarrow S^1 & & \mathbb{R} \\ S^1 \hookrightarrow S^3 \rightarrow S^2 & \text{if the Hamiltonian is over} & \mathbb{C} \\ S^3 \hookrightarrow S^7 \rightarrow S^4 & & \mathbb{H} \end{array}$$

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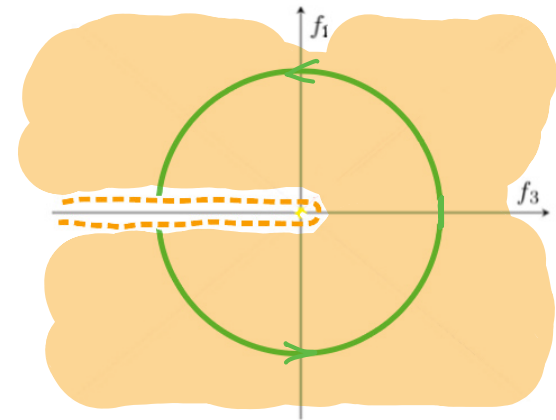
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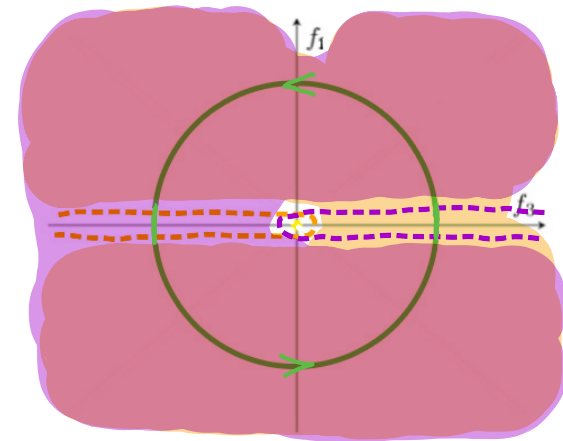
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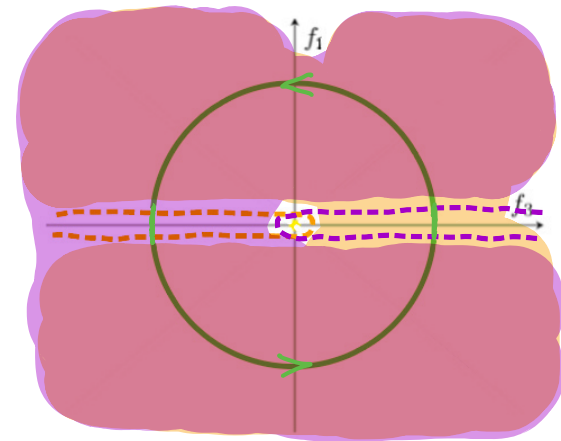
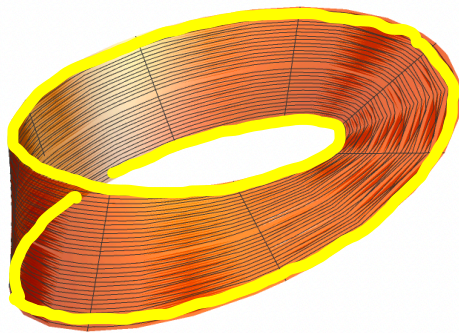
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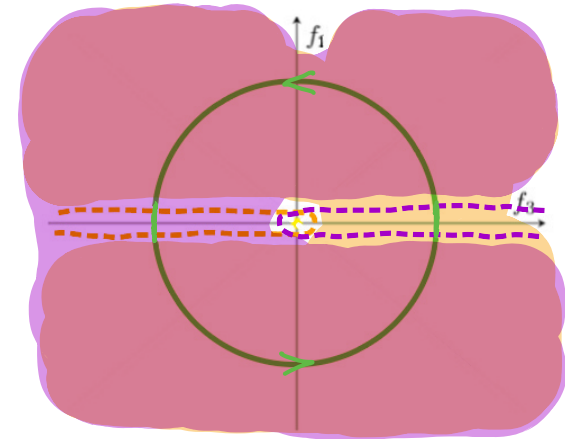
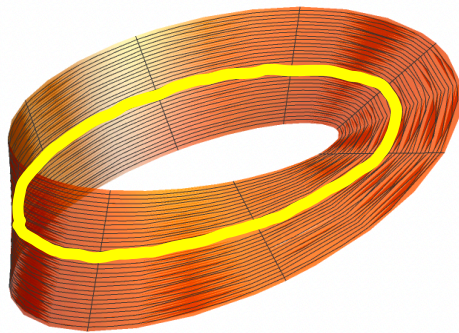
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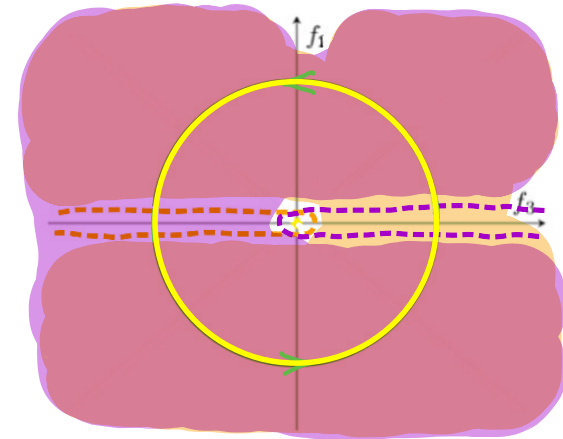
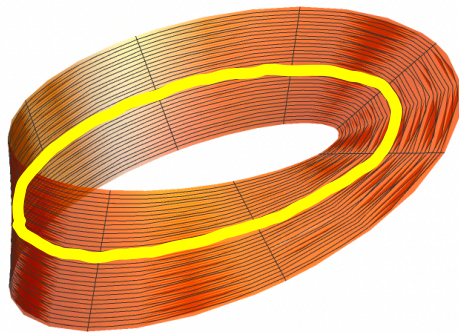
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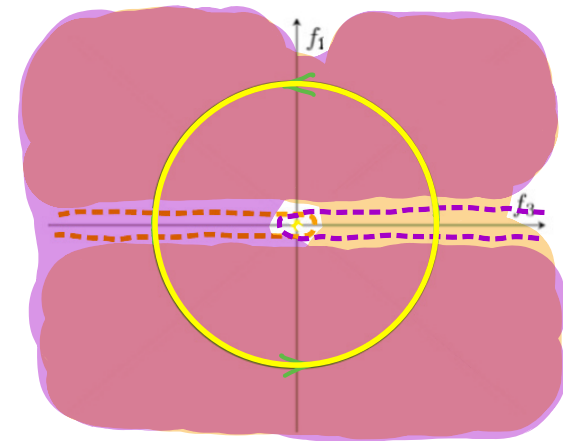
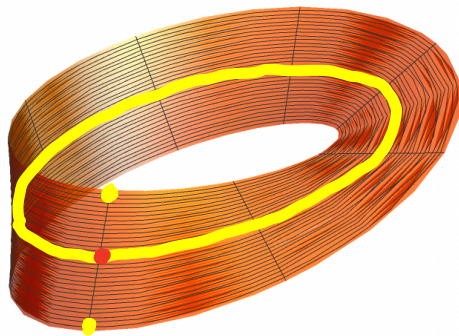
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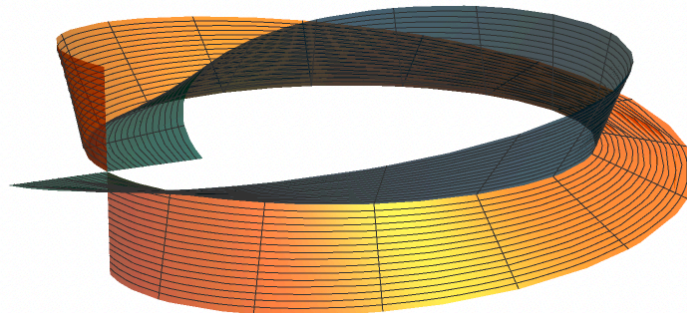
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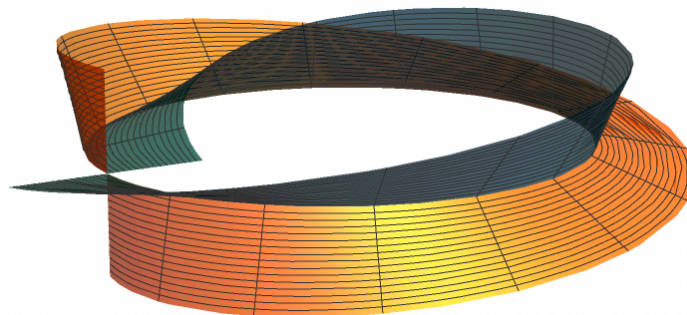


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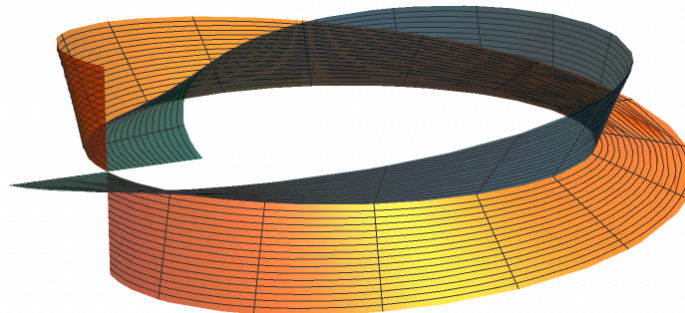


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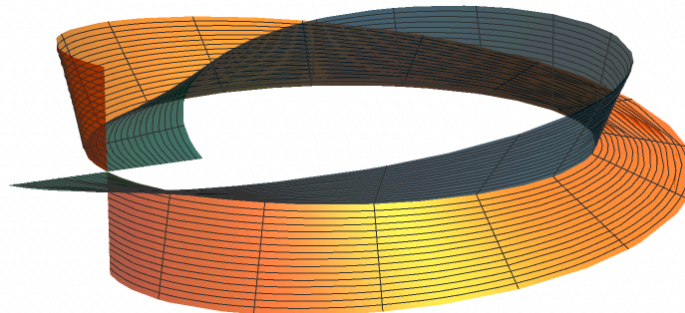


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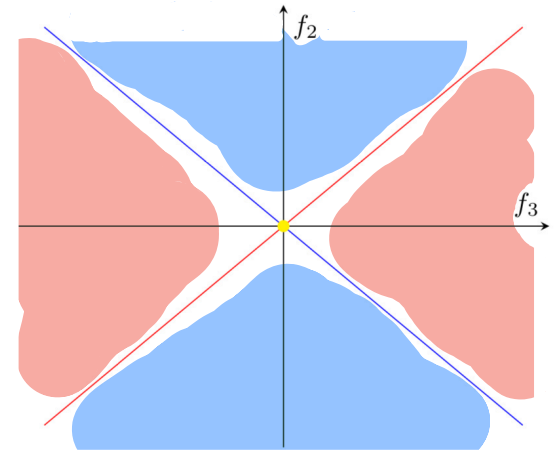
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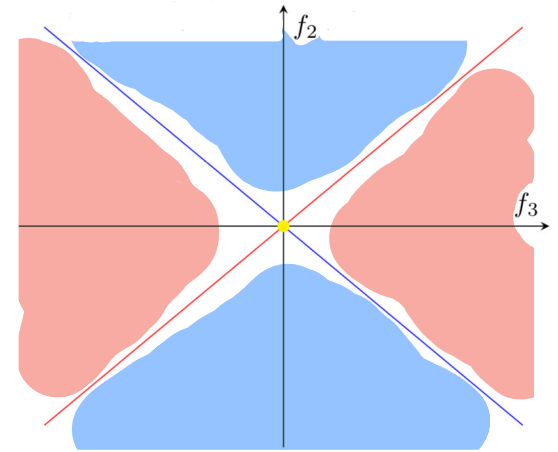


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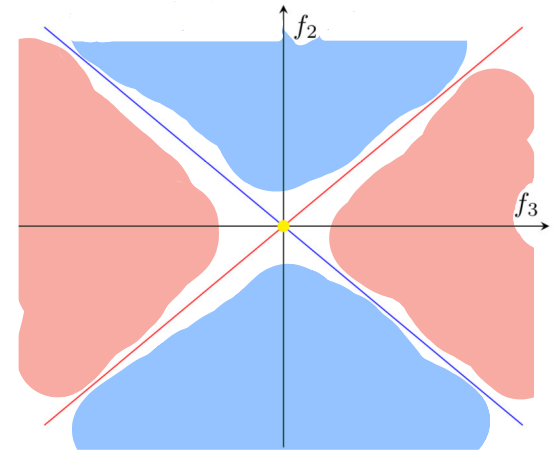


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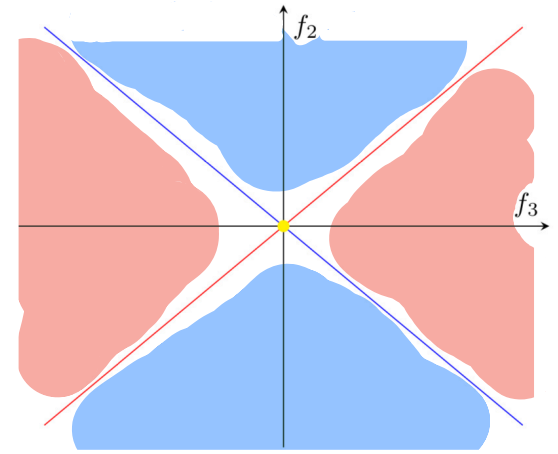


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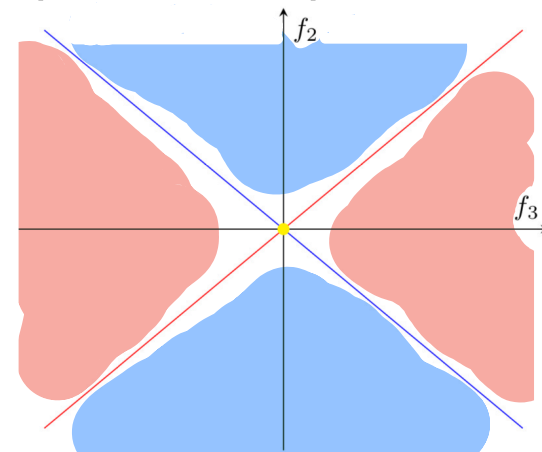


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Peter Higgs (bosons)

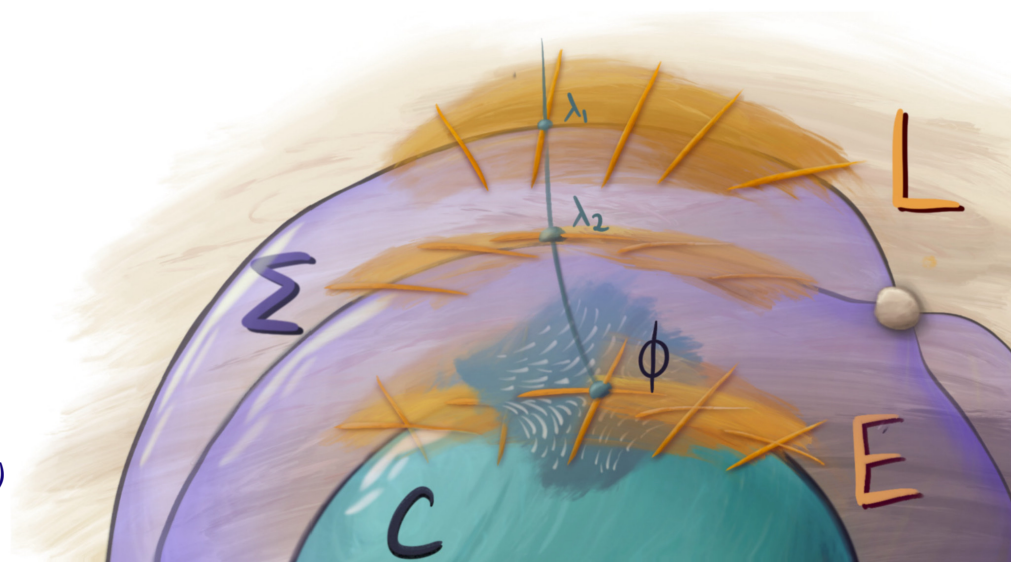
Nigel Hitchin 1987

Carlos Simpson

C compact Riemann surface (or more generally Kähler manifold)

E holomorphic vector bundle

ϕ Higgs field: a holomorphic 1-form taking values in the bundle of endomorphisms of E such that $\phi \wedge \phi = 0$

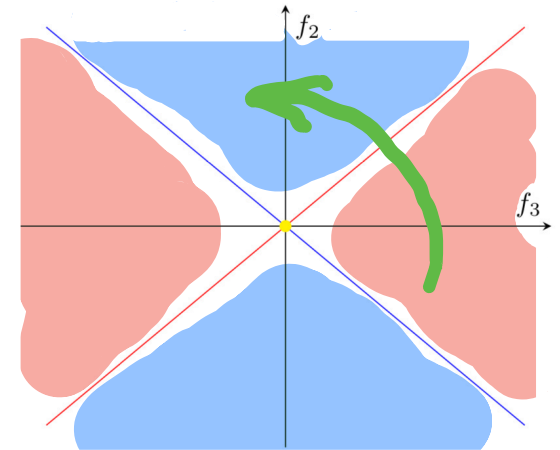


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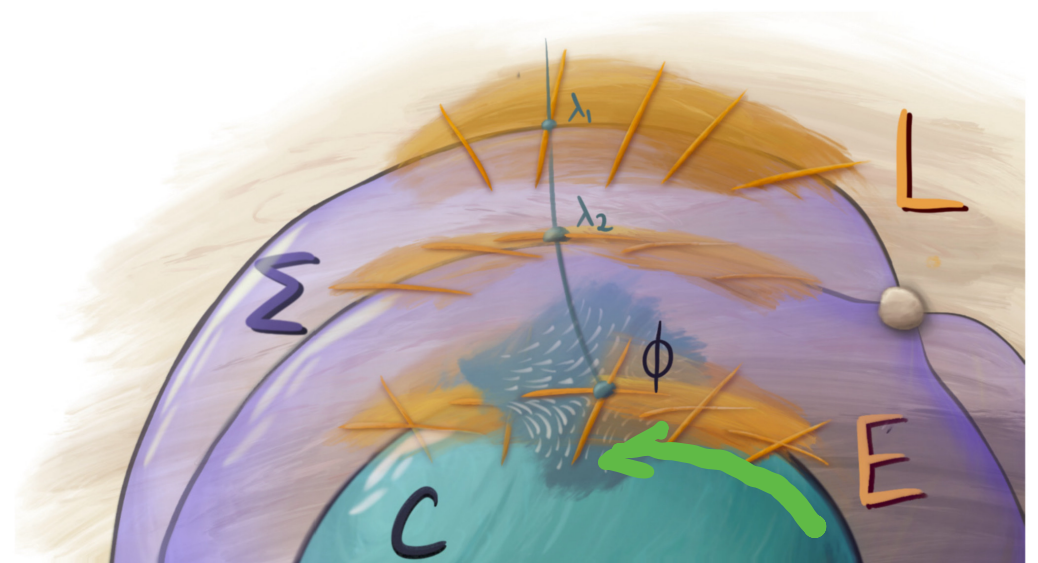
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$$\phi_x \in \text{End}(E_x), x \in C$$

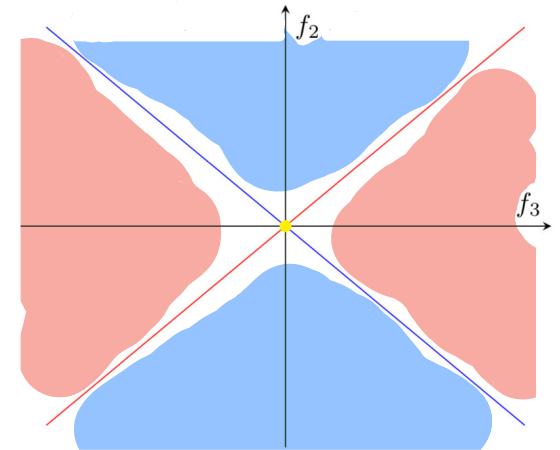


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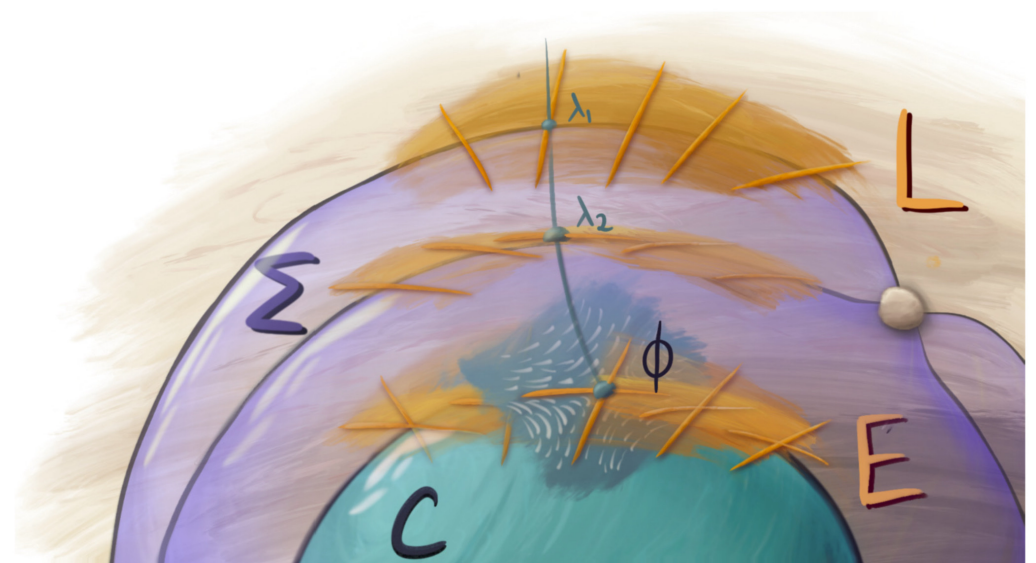
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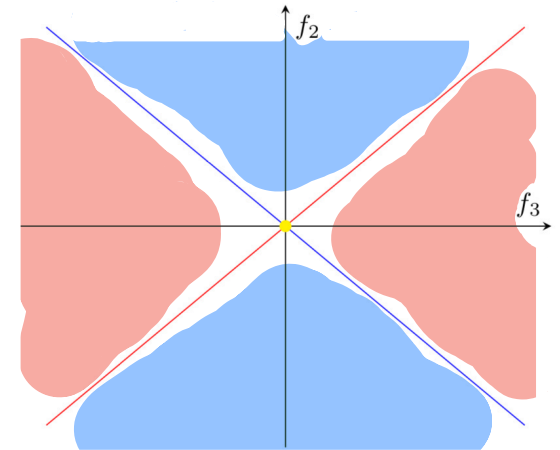


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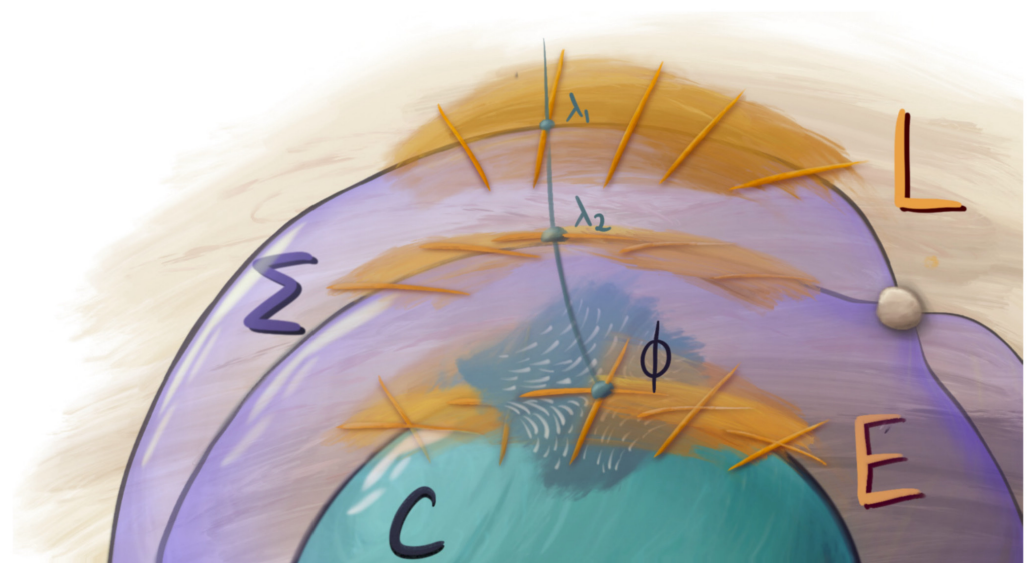
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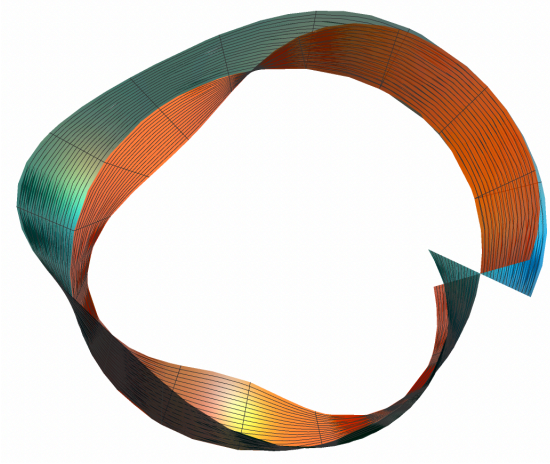
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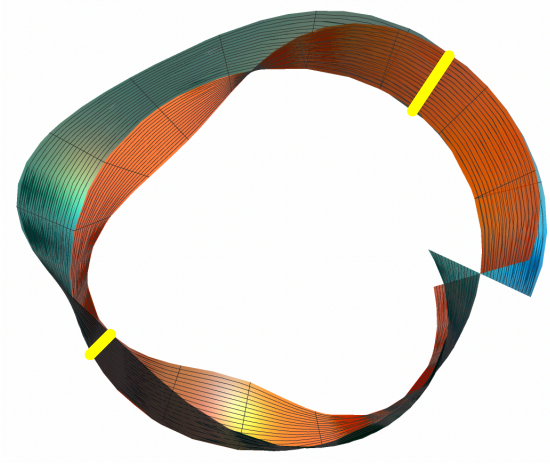
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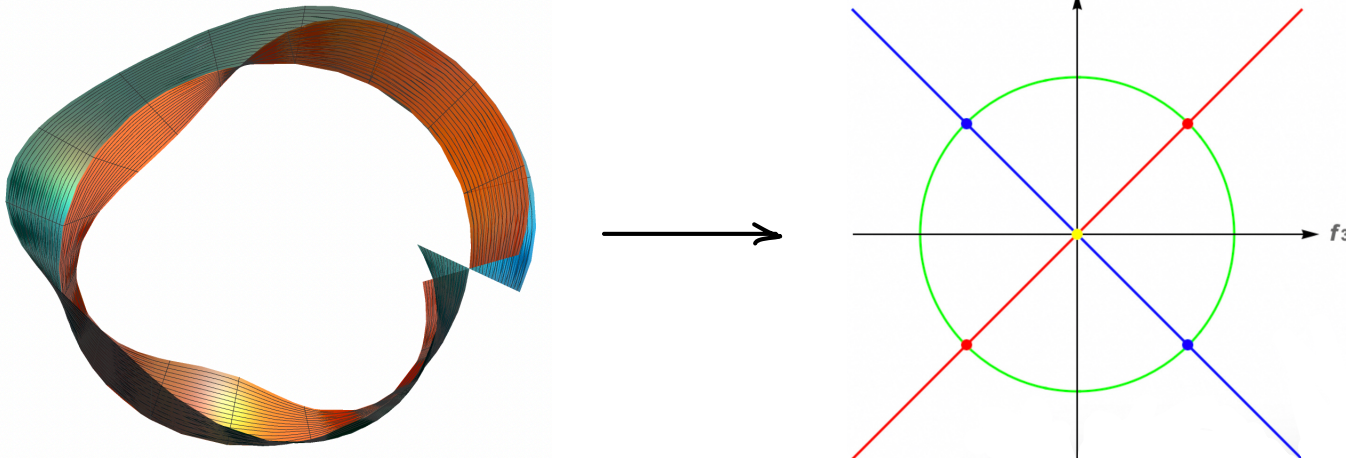
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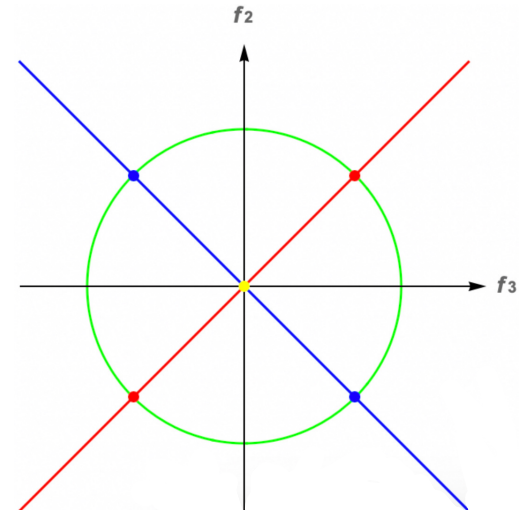
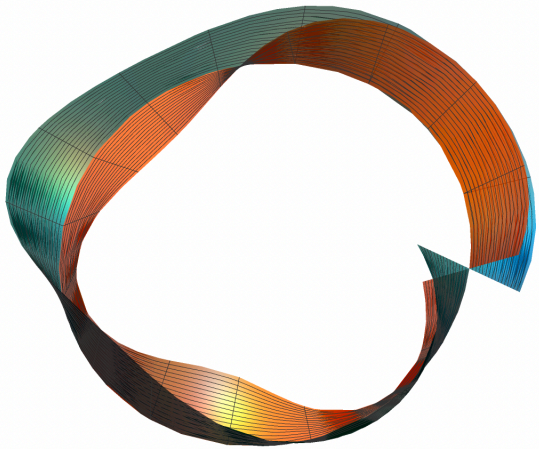
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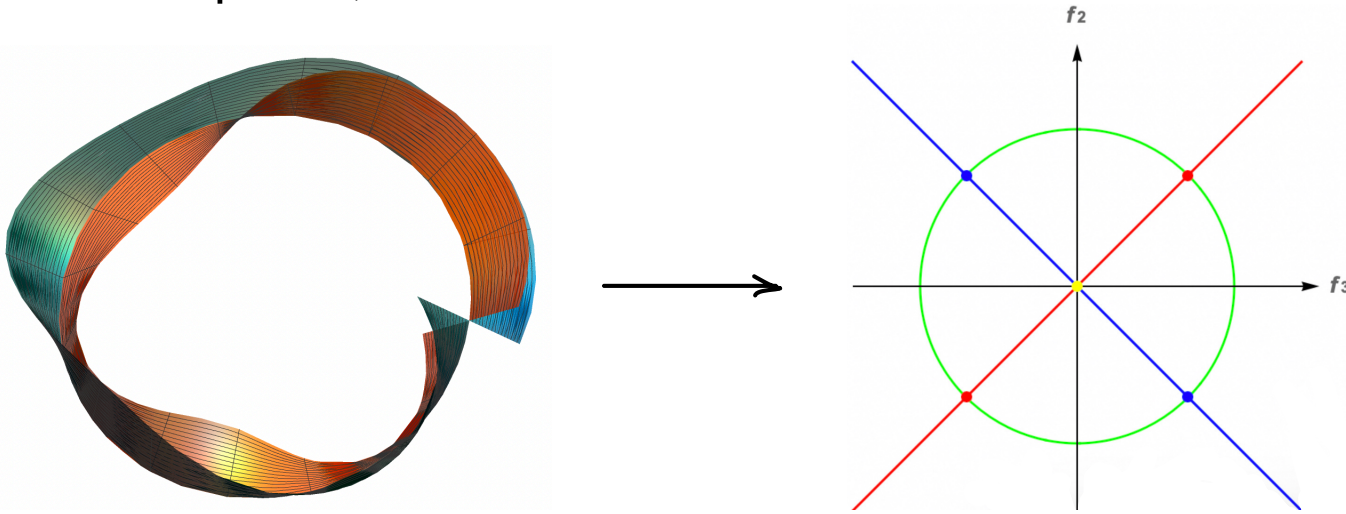
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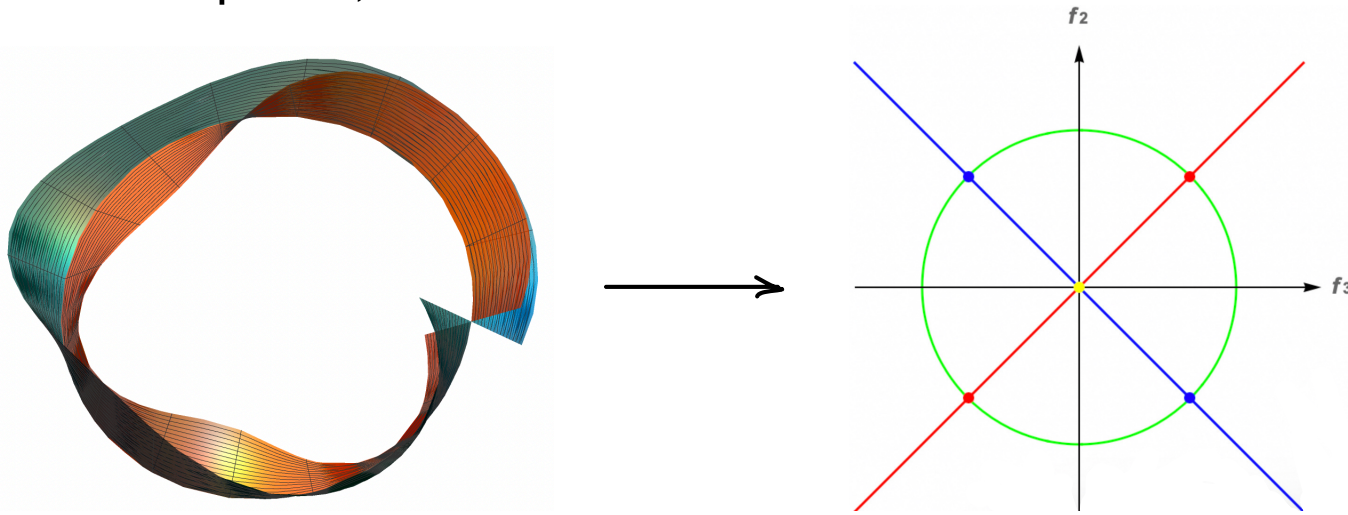
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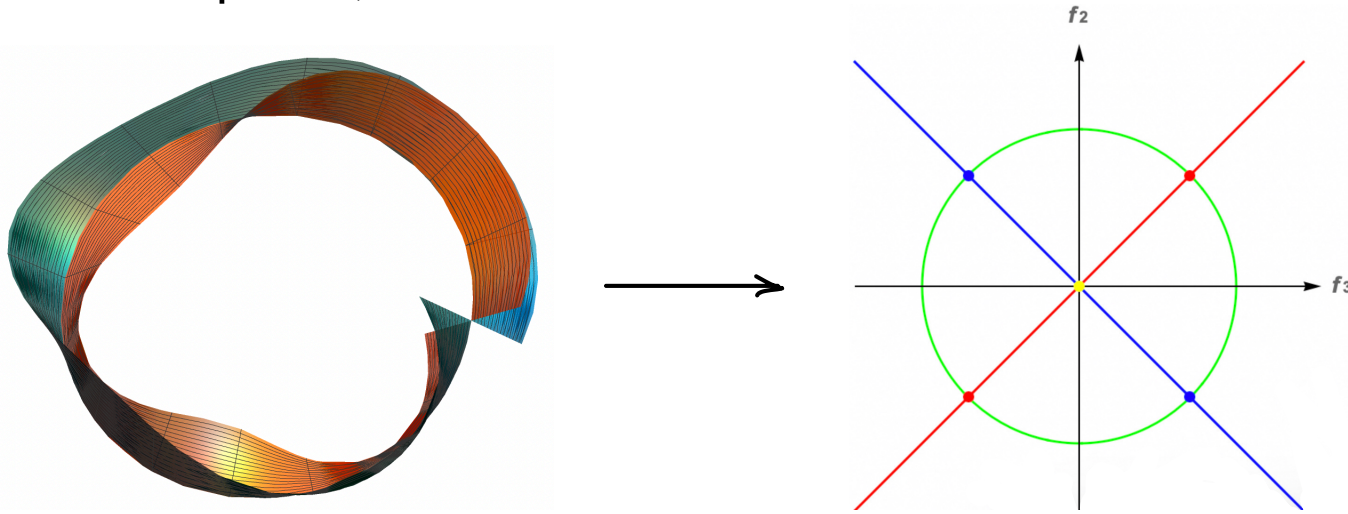


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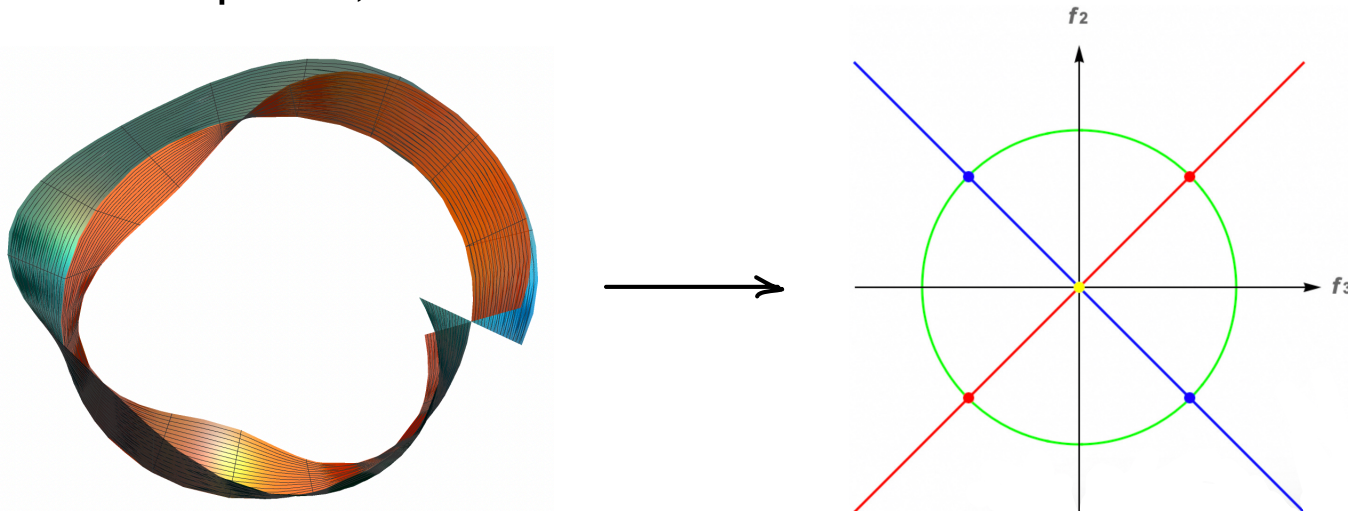


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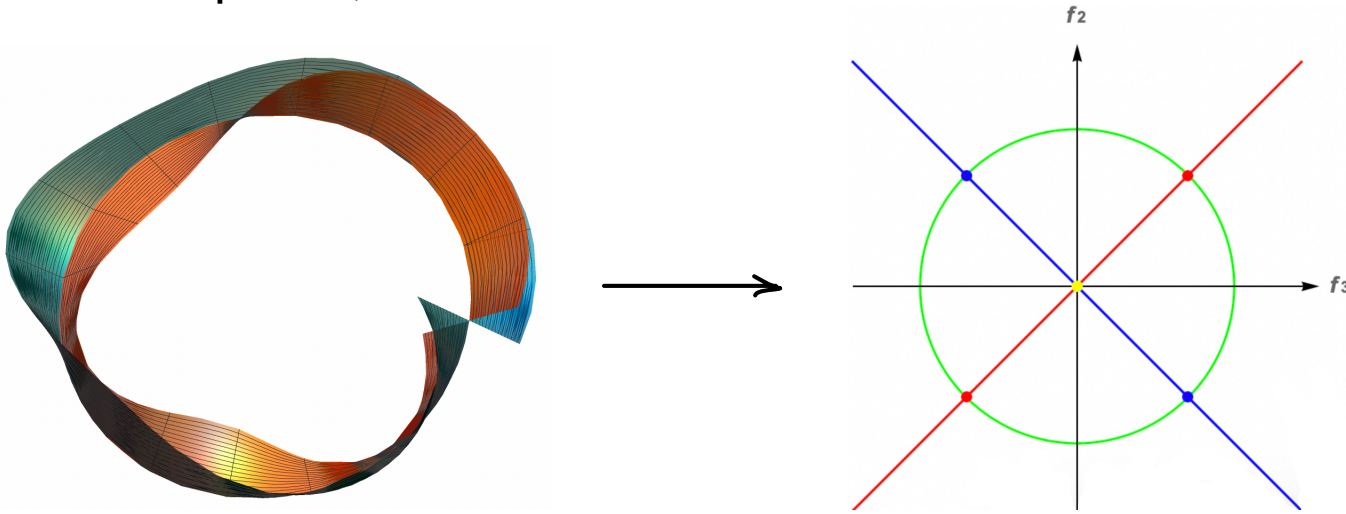
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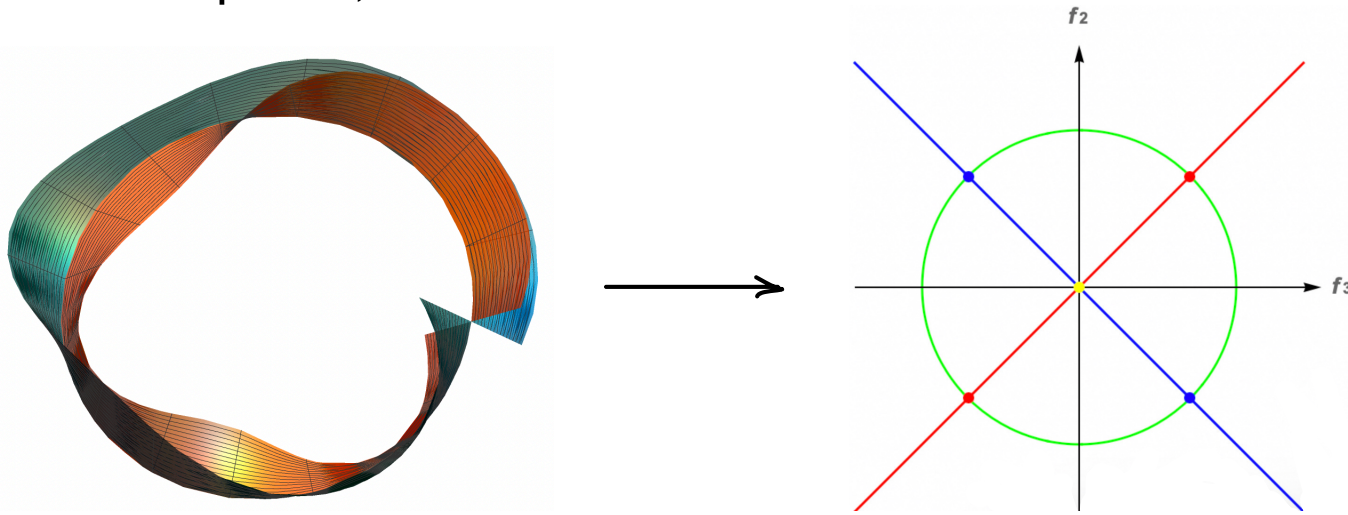
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Gajer, The intersection Dold–Thom theorem, Topology, 1996. (Ph.D. student of Blaine Lawson, 1993)

Goresky and MacPherson, 1974.

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Conjecture. It does so through a *deformation* (or homotopy) of Riemannian metrics

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$$\eta_t = \begin{bmatrix} e^{i\pi t} & 0 \\ 0 & 1 \end{bmatrix}$$

Eigenframe rotation as Higgs bundles: The non-Hermitian case

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Here is a video of the eigenbundle deformation: <https://yifeizhu.github.io/swallowtail/deform.mp4>

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Example (Swallowtail quadruple $sw4$).

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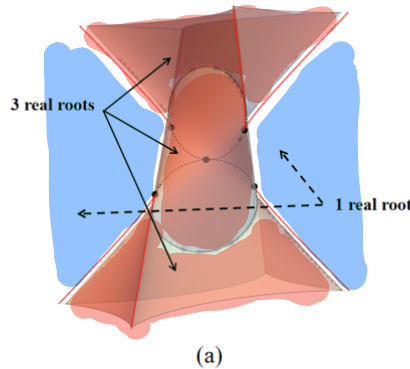
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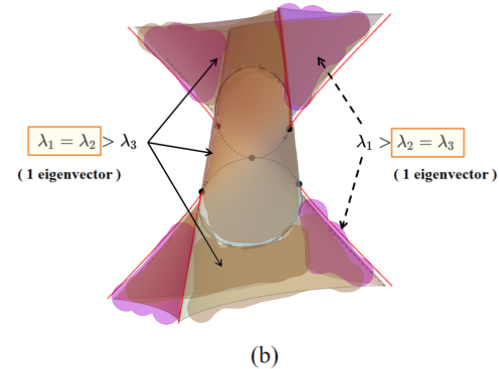
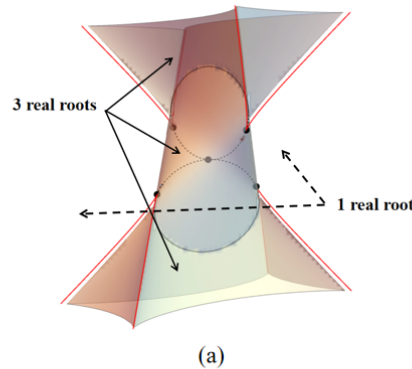
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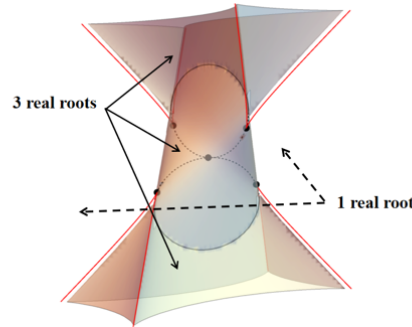
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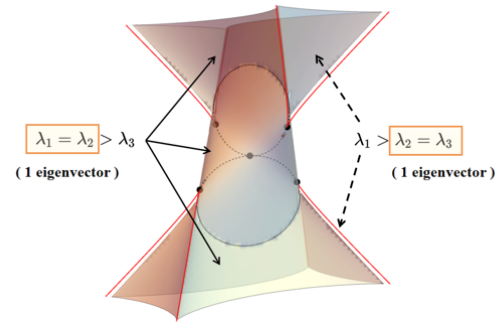
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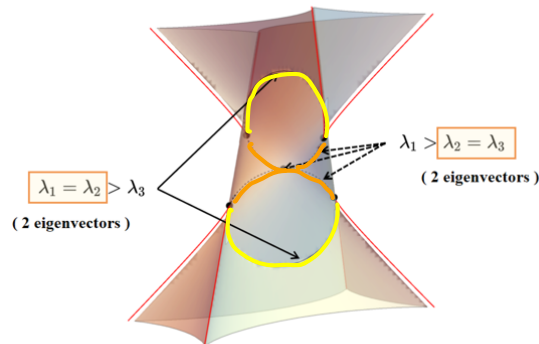
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(b)



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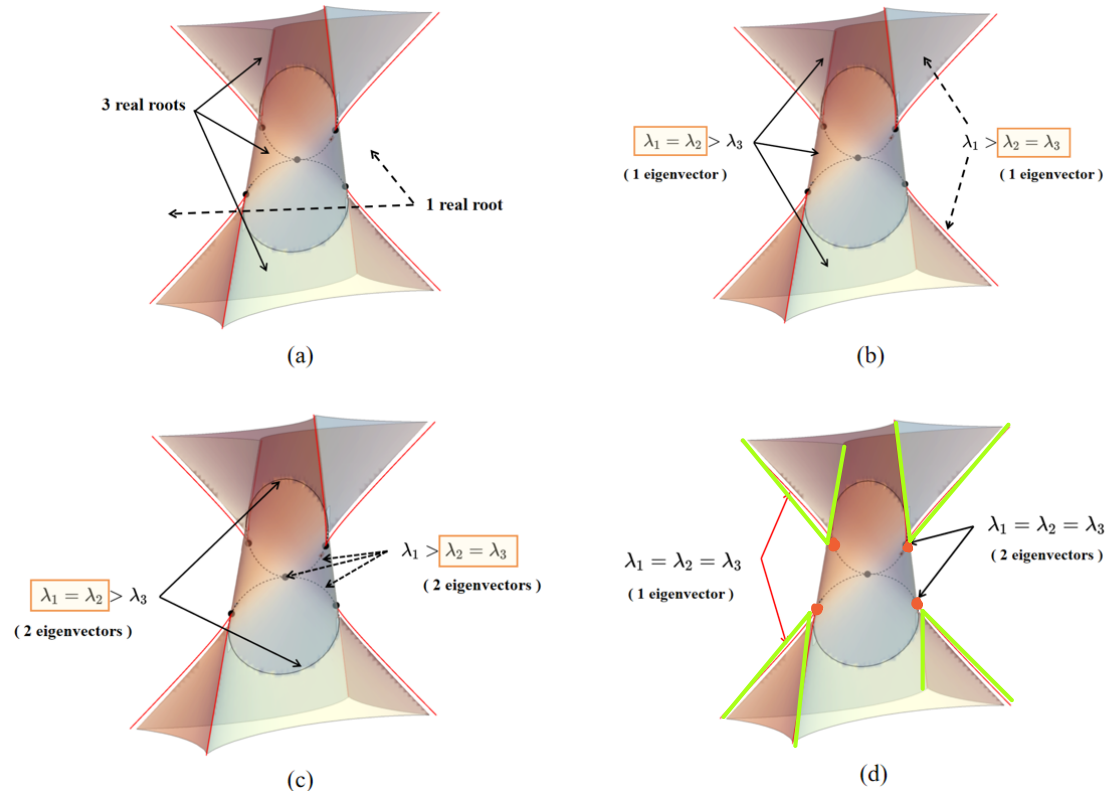
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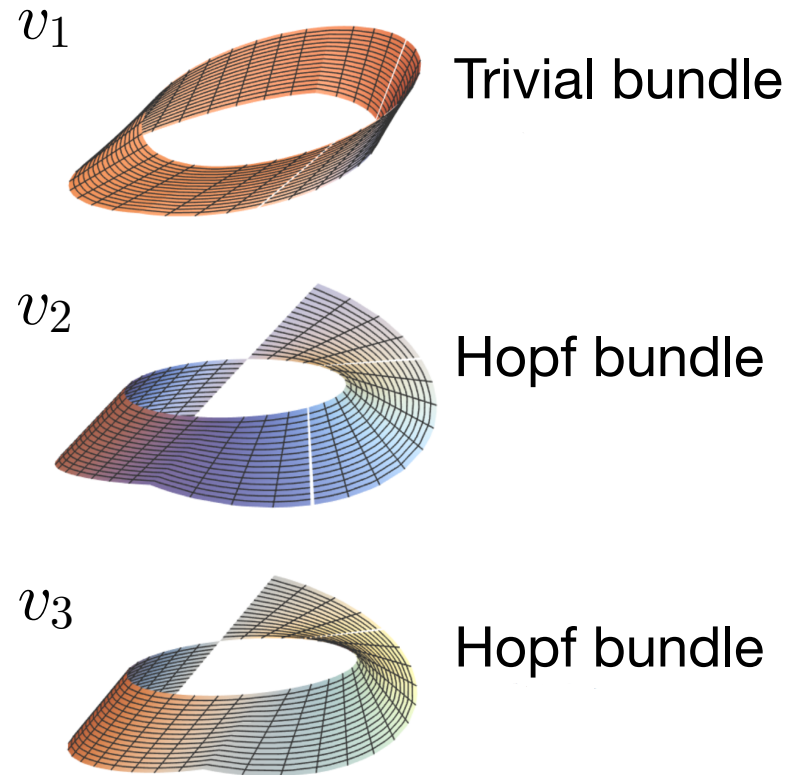
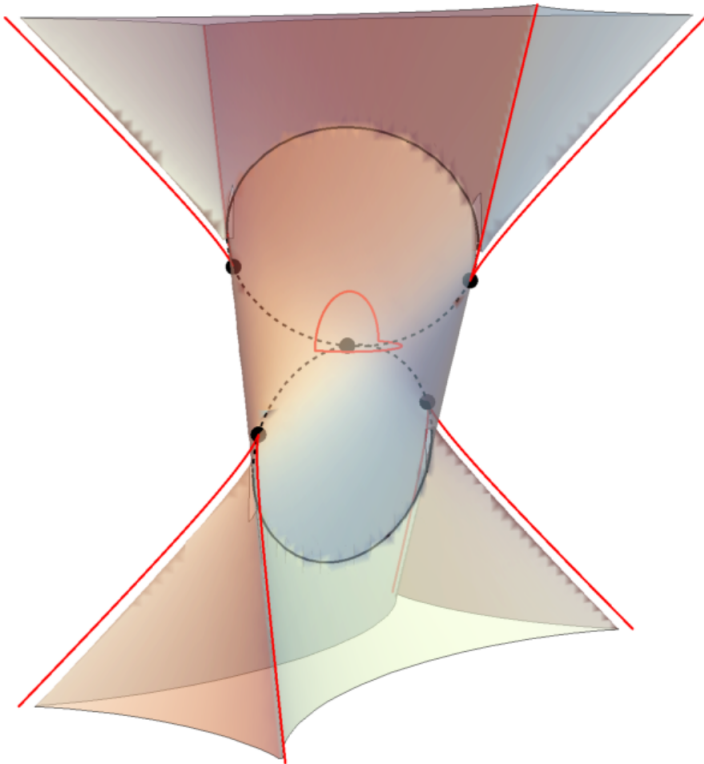


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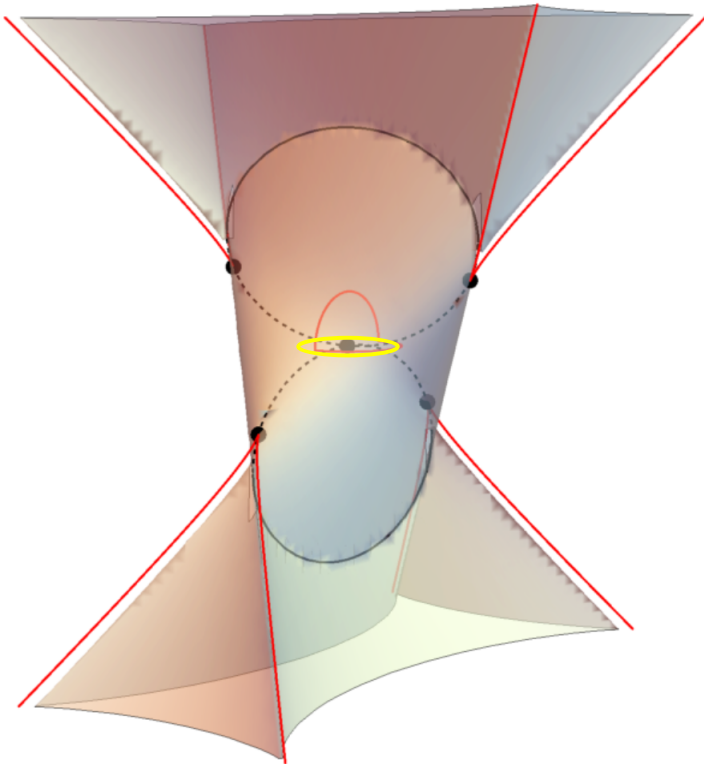


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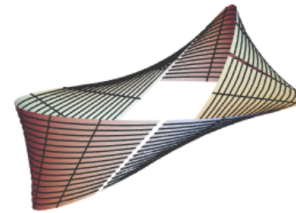
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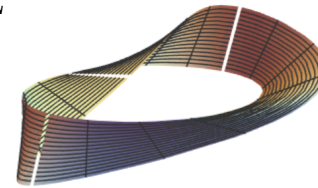


v_1



Trivial bundle

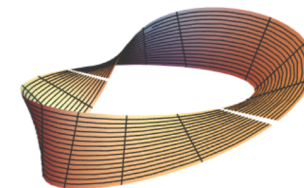
v_2



~~Hopf bundle~~

Trivial bundle

v_3



~~Hopf bundle~~

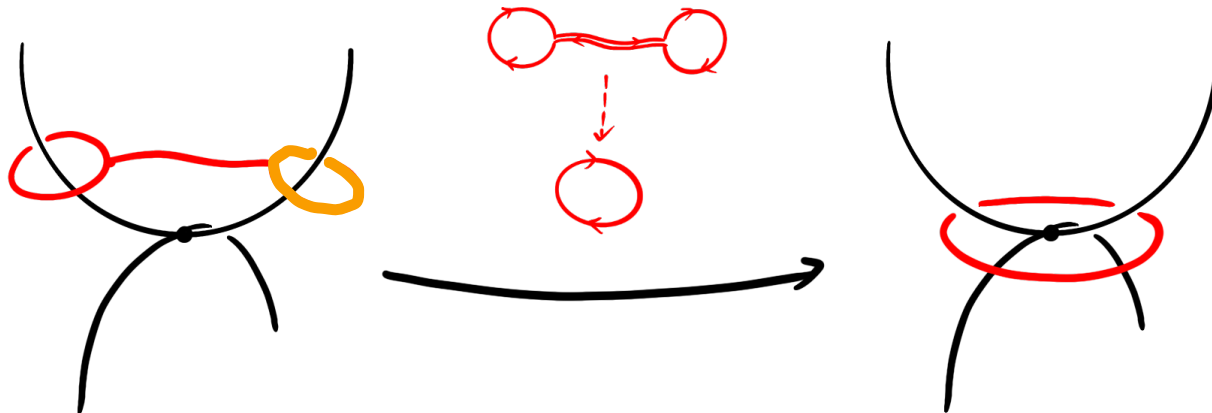
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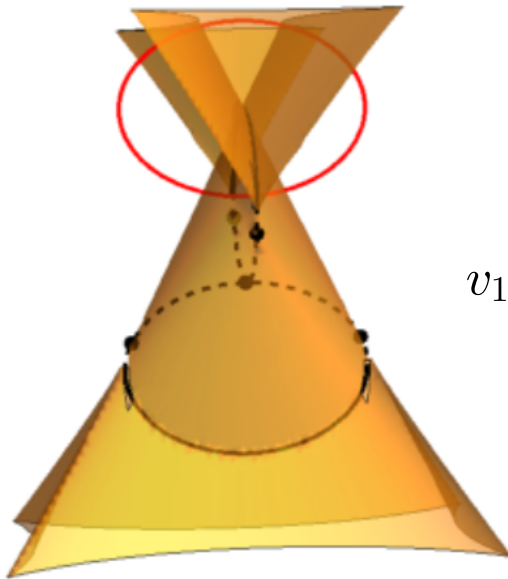


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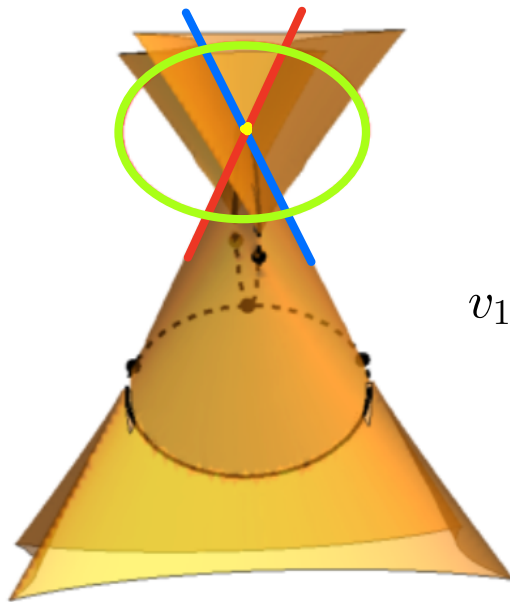
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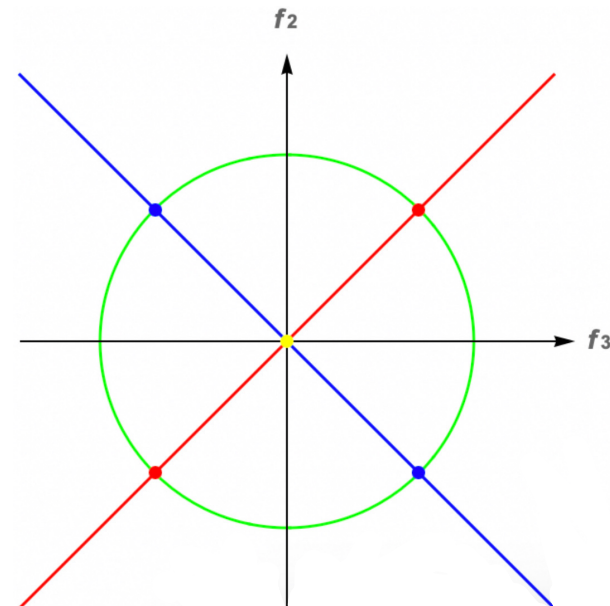
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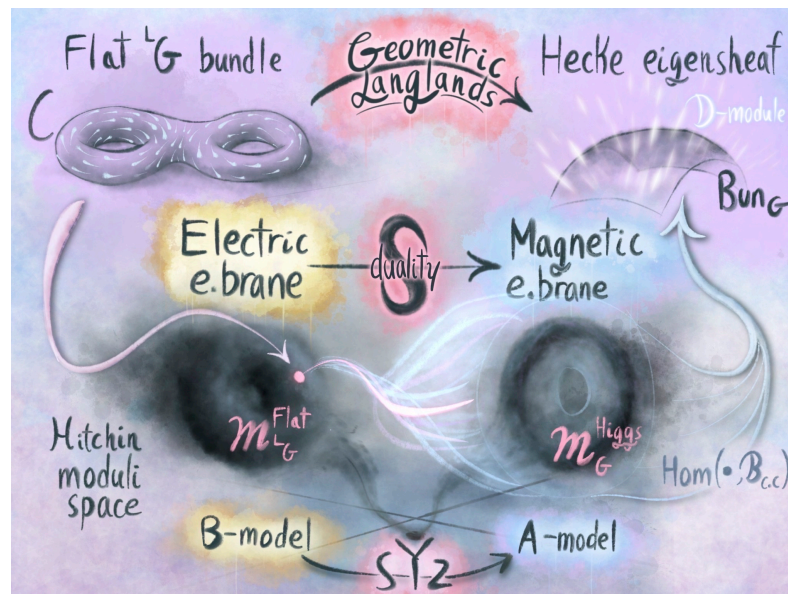
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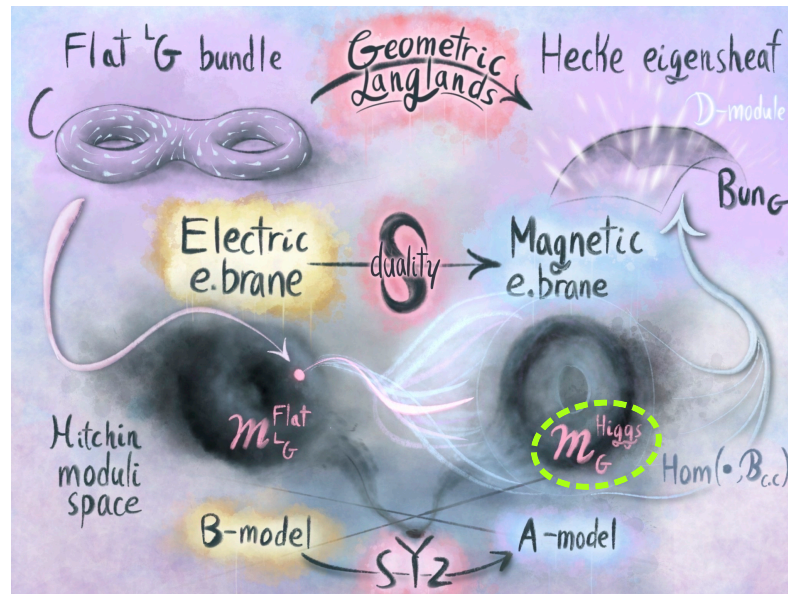
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There has not been a rigorous mathematical explanation for such a correspondence in general, but it is reminiscent of the **Langlands duality**. Indeed, **Higgs bundles** sit on one side of the geometric Langlands duality! We've at least found some **testing ground**.



Thank you.