

# Topology in magnetic systems

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# Introduction



## Electron = Charge + Spin

- There has been great achievements on the “**charge**” properties in semiconductor research since artificial structures were introduced in late 1960s to early 1970s.
- Similar pathway has been followed on the “**spin**” properties in magnetism research since artificial structures were introduced in 1980s.
- There is increasing interest on the combination of “**charge**” and “**spin**” to develop spintronics technology, especially since the discovery of Giant Magnetoresistance (2007 Nobel Prize in physics).

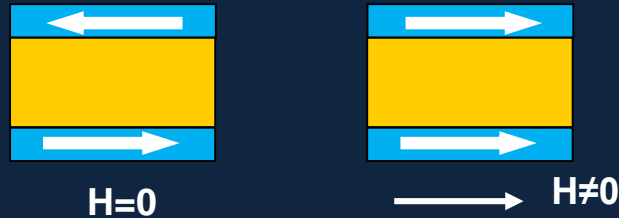


Albert Fert



Peter Grünberg

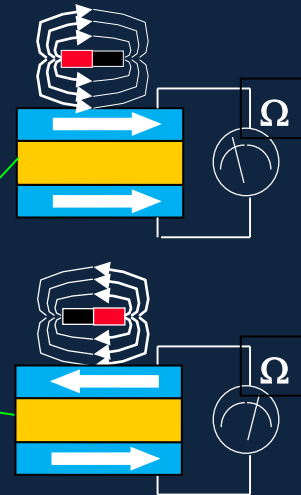
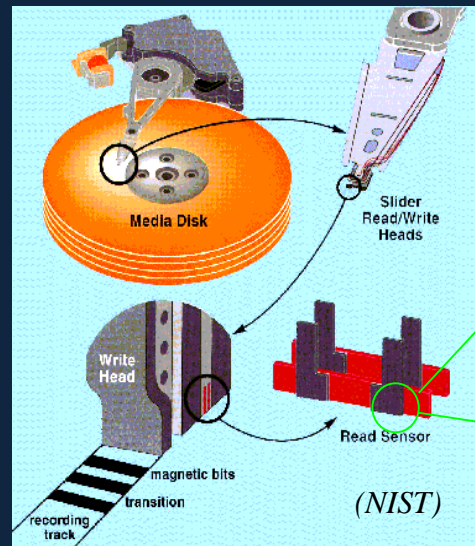
# Antiferromagnetic Coupling and Giant Magnetoresistance (GMR)



P. Grünberg, et. al., PRL **57**, 2442 (1986).

$$\text{GMR} = [R(0)-R(H)]/R(H) = 10-100\%$$

M. N. Baibich, et al., PRL **61**, 2472 (1988).  
G. Binasch, et al. PRB **39**, 4828 (1989).



# What will be the next breakthrough?

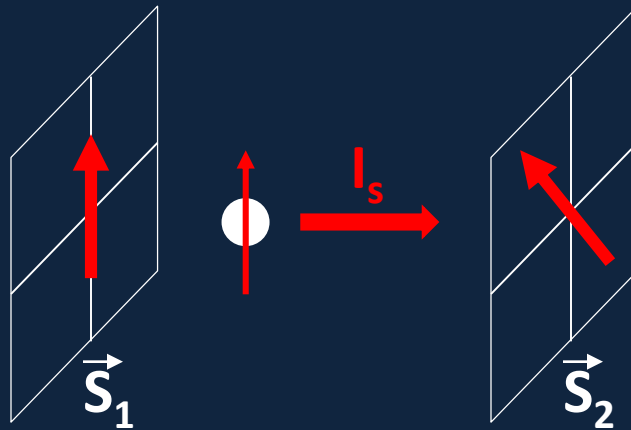
No one knows for sure, but two have emerged recently.

- Spin torque effect (spin current)

## 2013 Buckley Prize: Spin-Transfer Torque

J.C. Slonczewski, J. of Magn. Magn. Mat. 159, L1-L7 (1996).

L. Berger, Phys. Rev. B 54, 9353 (1996).



$$\dot{\vec{S}}_2 \sim \frac{I_e g}{e} \vec{S}_2 \times (\vec{S}_1 \times \vec{S}_2)$$

The flow of spins ( $I_s$ ) delivers a torque to  $S_2$ .

- Topological effect (vortices, skyrmions)

# Topology and Skyrmion

Topology



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# Topology and Skyrmion

Topology



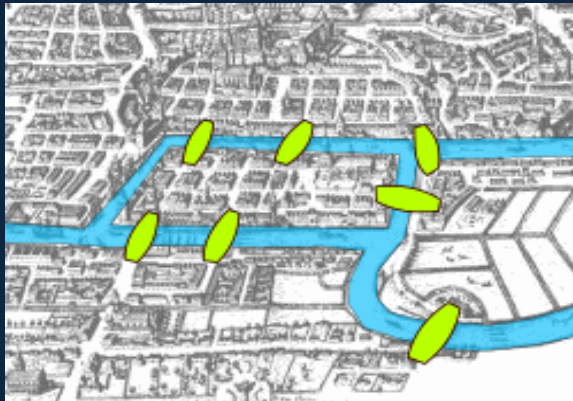
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## Königsberg Bridge Problem



Is it possible to find a route through the city that will cross each bridge once and only once?

# Topology and Skyrmion

Topology



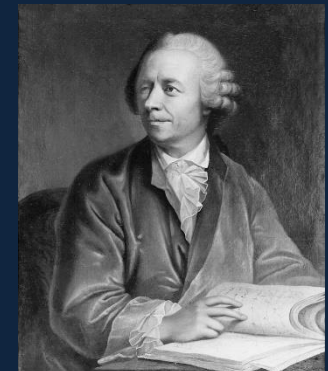
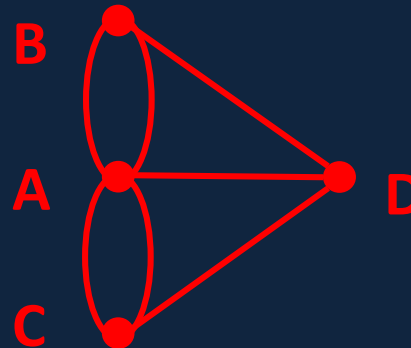
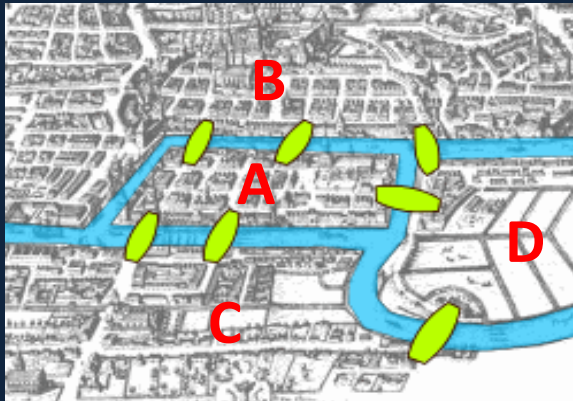
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## Königsberg Bridge Problem



Leonhard Euler

4/15/1707-9/18/1783

# Topology and Skyrmion

Topology



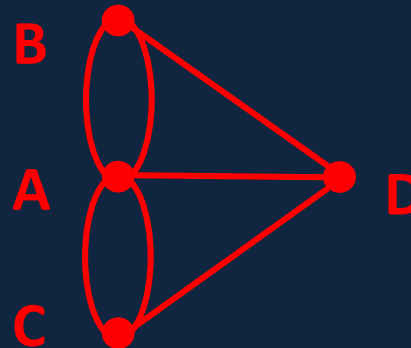
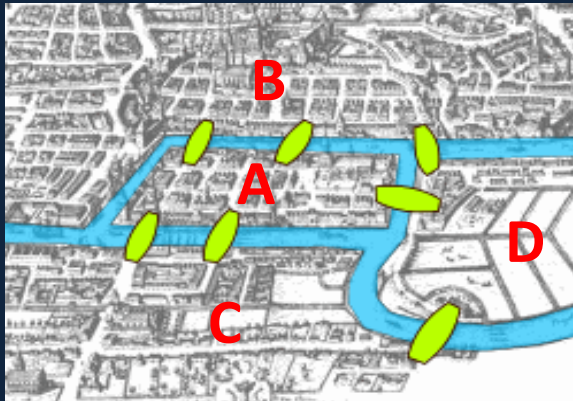
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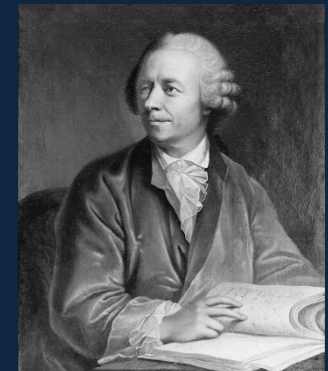
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## Königsberg Bridge Problem



*“This question is so banal, but seemed to me worthy of attention in that [neither] geometry, nor algebra, nor even the art of counting was sufficient to solve it.”*



**Leonhard Euler**  
4/15/1707-9/18/1783

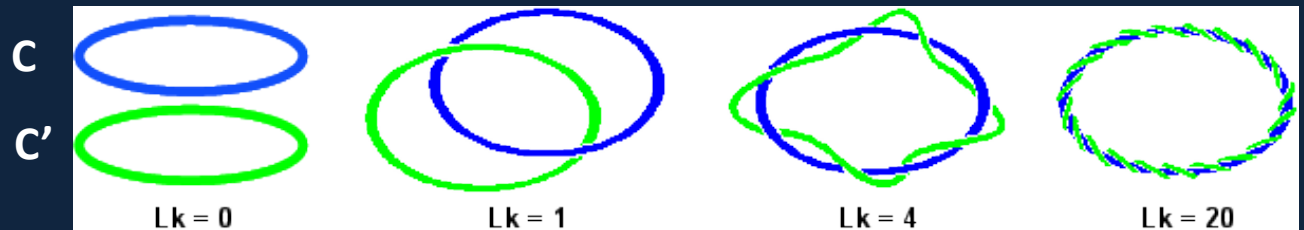


# Linking Number

What's the number of times that one closed loop winds around the other closed loop?



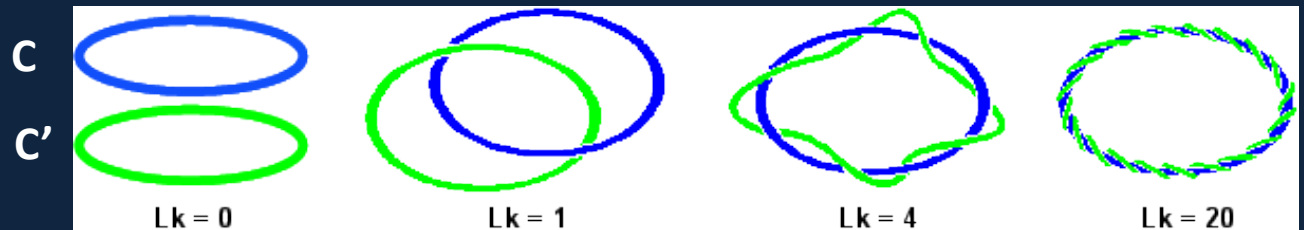
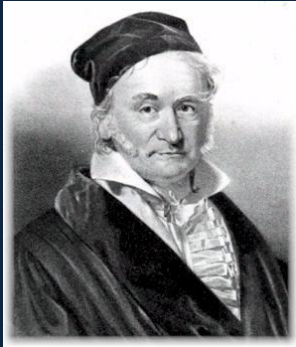
Johann Carl Friedrich Gauss  
4/30/1777-2/23/1855



$$\text{linking number} \quad n = \frac{1}{4\pi} \oint_C \oint_{C'} \frac{(d\vec{r} \times d\vec{r}') \cdot (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3}$$

# Linking Number

What's the number of times that one closed loop winds around the other closed loop?



Johann Carl Friedrich Gauss  
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$$\text{linking number} \quad n = \frac{1}{4\pi} \oint_C \oint_{C'} \frac{(d\vec{r} \times d\vec{r}') \cdot (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3}$$

# Electromagnetism

Biot–Savart law

$$\vec{B}(\vec{r}) = \frac{\mu_0 I_0}{4\pi} \oint_{C'} \frac{d\vec{r}' \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3}$$

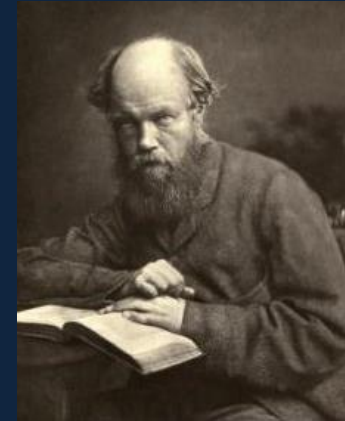
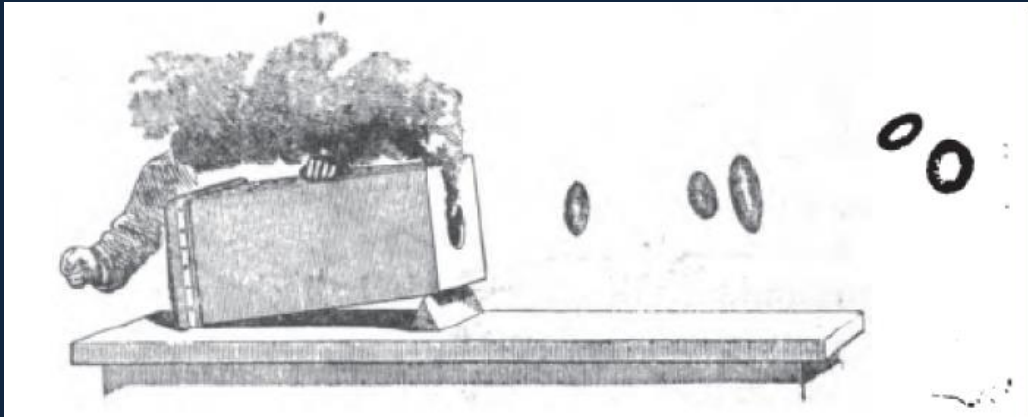
Ampère's law

$$\oint_C \vec{B}(\vec{r}) \cdot d\vec{r} = \mu_0 I_{enc} = \mu_0 n I_0$$



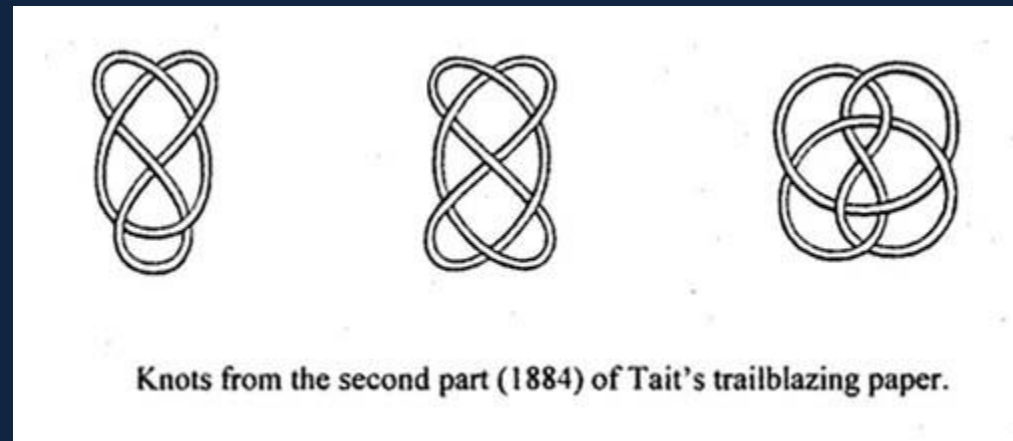
$$\frac{1}{4\pi} \oint_C \oint_{C'} \frac{[d\vec{r}' \times (\vec{r} - \vec{r}')] \cdot d\vec{r}}{|\vec{r} - \vec{r}'|^3} = n$$

## Experiment: Smoke Rings



Peter Guthrie Tait  
4/28/1831-7/4/1901

Elements = knots of swirling vortices in the ether





Tony Skyrme

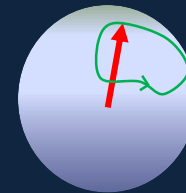
12/5/1922-6/25/1987

Skyrme, T. H. R. A Non-Linear Theory of Strong Interactions. *Proc. R. Soc. London Ser. A* **247**, 260-278 (1958).

Skyrme, T. H. R. A unified field theory of mesons and baryons. *Nuclear Physics* **31**, 556-569 (1962).

## Skyrmions in 2D Heisenberg system

$$H = -J \sum_{\langle i,j \rangle} \vec{S}_i \cdot \vec{S}_j = J \int d^2\vec{r} [\nabla \vec{n}(\vec{r})]^2 \quad \vec{n} \equiv \frac{\vec{S}}{S}$$



Berry Phase = S x Solid Angle

**Question:** How many static solutions?

**Answer:** Many solutions which are “quantized” by an integer index N.

### Metastable states of two-dimensional isotropic ferromagnets

A. A. Belavin and A. M. Polyakov

Gor'kii State University

(Submitted October 4, 1975)

Pis'ma Zh. Eksp. Teor. Fiz. **22**, No. 10, 503-506 (20 November 1975)

$$q = \frac{1}{8\pi} \int \epsilon_{\alpha\beta\gamma} \epsilon_{\mu\nu} n^\alpha \frac{\partial n^\beta}{\partial x_\mu} \frac{\partial n^\gamma}{\partial x_\nu} d^2x \quad (5)$$

### Skyrmion Number

$$N = \int \frac{dxdy}{4\pi} \vec{n} \cdot (\partial_x \vec{n} \times \partial_y \vec{n}) = \text{integer}$$

# Mapping of 2D spins from the xy plane to a sphere

Skyrmion solution

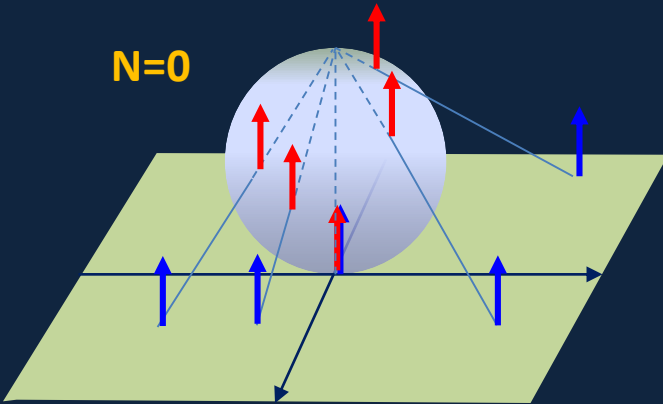
$$\vec{S} = (\sin \theta \cos \phi, \sin \theta \sin \phi, \cos \theta) \quad \Leftrightarrow \quad w = (x, y)$$

$$w = x + iy = \frac{S_x + iS_y}{1 - S_z} = \text{ctg}(\theta/2)e^{i\phi}$$

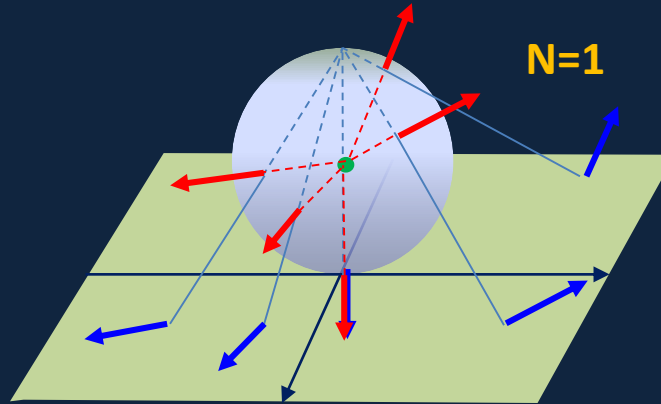
**Skyrmion Number**

$$N = \int \frac{dxdy}{4\pi} \vec{n} \cdot (\partial_x \vec{n} \times \partial_y \vec{n})$$

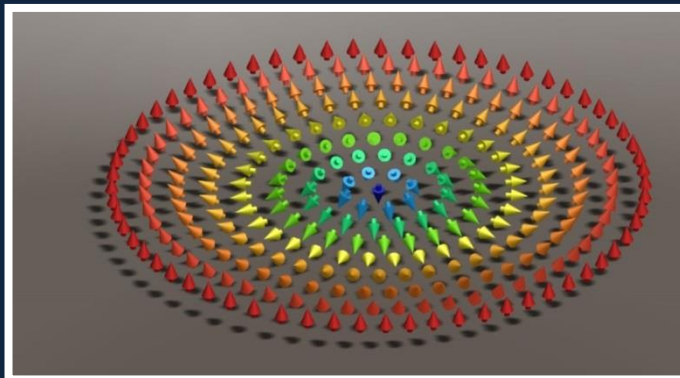
N=0



N=1

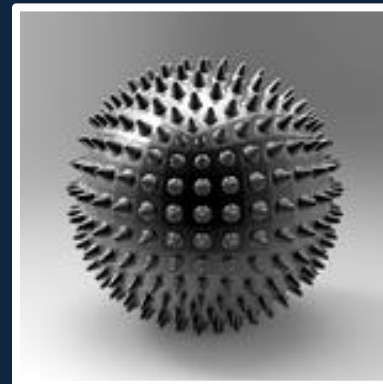


$$\begin{cases} S_x = \frac{2Rr}{R^2 + r^2} \cos(\theta + \phi) \\ S_y = \frac{2Rr}{R^2 + r^2} \sin(\theta + \phi) \\ S_z = \frac{R^2 - r^2}{R^2 + r^2} \end{cases}$$



**Skyrmion of N=1**

=



**Spiky Ball**



## Grreat Choice® Spiky Ball Set Dog Toy



**\$ 6.99**  
**Free Shipping**

# Creation of magnetic Skyrmions in experiment

- **Dzyaloshinskii-Moriya (DM) interaction**

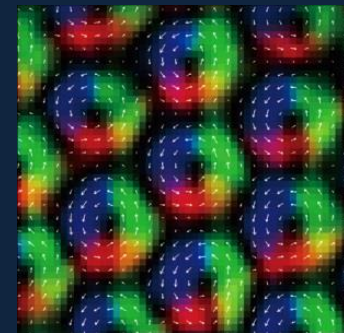
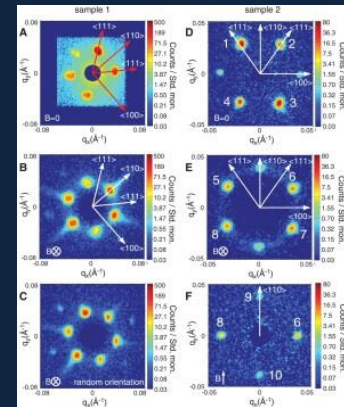
$$H_{DM} = -\vec{D}_{DM} \cdot (\vec{S}_i \times \vec{S}_j)$$

U. K. Röbler, A. N. Bogdanov, & C. Pfleiderer, “Spontaneous skyrmion ground states in magnetic metals”. *Nature* **442**, 797-801 (2006).

S. Mühlbauer et al. “Skyrmion Lattice in a Chiral Magnet”. *Science* **323**, 915-919 (2009).

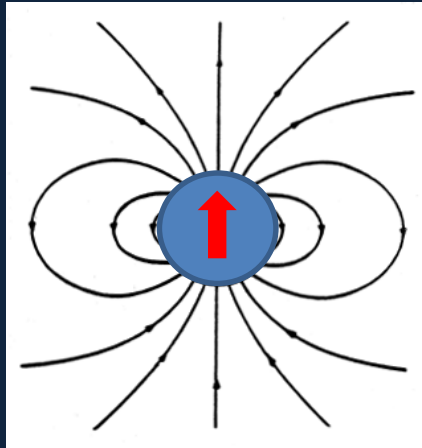
X. Z. Yu et al. “Real-space observation of a two-dimensional skyrmion crystal” *Nature* **465**, 901–904 (17 June 2010).

Naoto Nagaosa & Yoshinori Tokura, “Topological properties and dynamics of magnetic skyrmions”. *Nature Nanotechnology* **8**, 899-911 (2013).



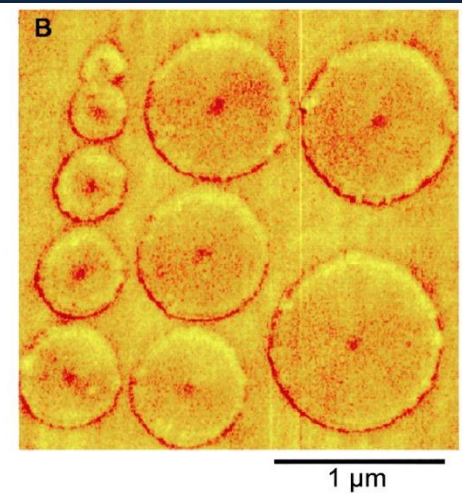
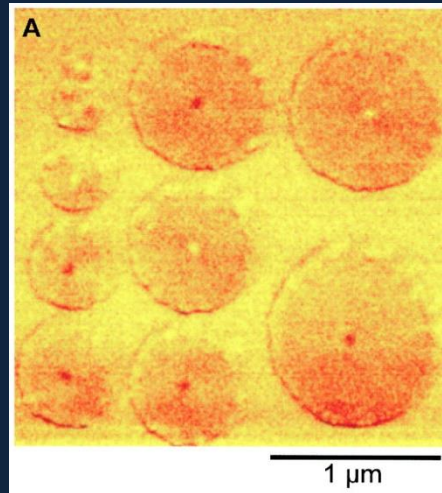
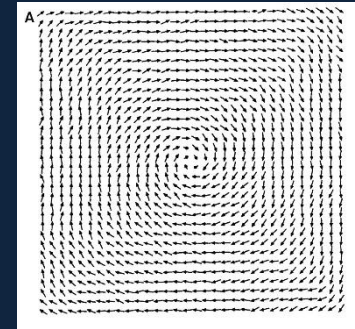
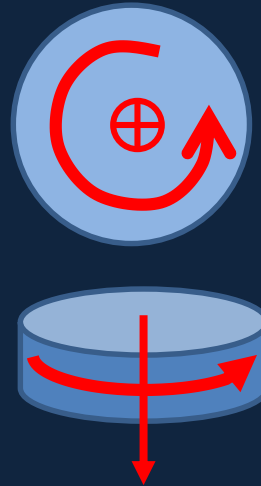


# Magnetic Vortex vs skyrmion



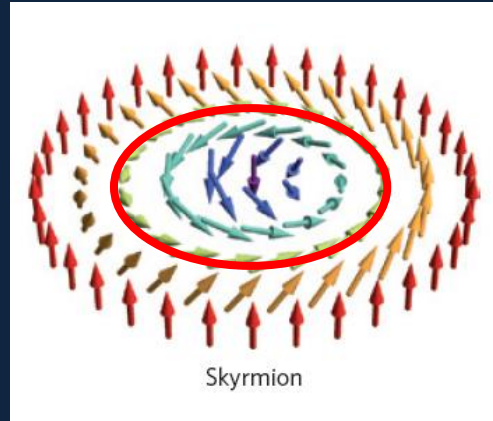
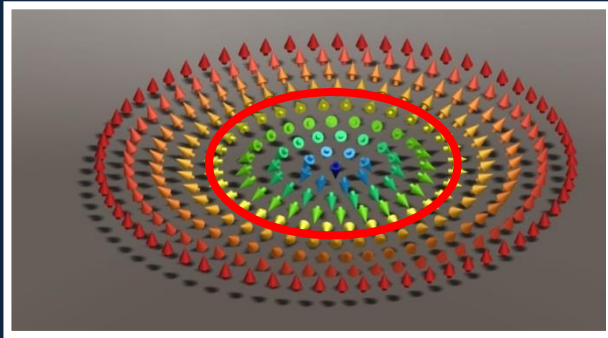
$$\sigma_M = \vec{M} \cdot \hat{n}$$

$$H = -J \sum_{\langle i,j \rangle} \vec{S}_i \cdot \vec{S}_j$$



T. Shinjo et. al. *Science* **289**, 930-932 (2000).





**One magnetic vortex is  $\frac{1}{2}$  skyrmion**

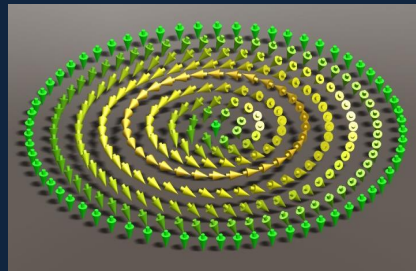
O. A. Tretiakov and O. Tchernyshyov, *Phys. Rev. B* **75**, 012408 (2007).

**One skyrmion = vortex + surrounding out-of-plane spins**

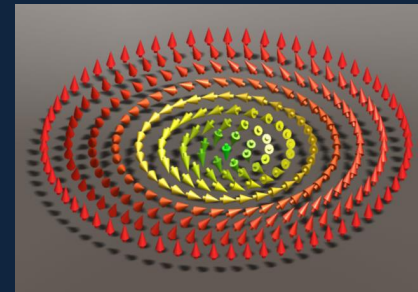
$$N = \int \frac{dxdy}{4\pi} \vec{n} \cdot (\partial_x \vec{n} \times \partial_y \vec{n})$$

**N=0** vortex core //  $\mathbf{S}_{\text{surrounding}}$

**N=1** vortex core //  $-\mathbf{S}_{\text{surrounding}}$



**N=0**

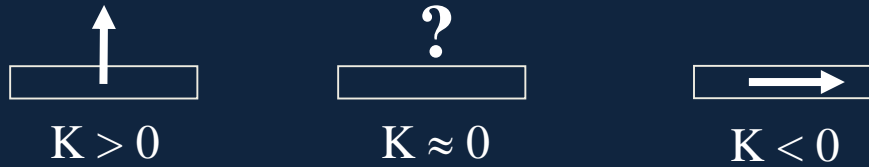


**N=1**

# Stripe, bubble, and artificial skyrmions in magnetic thin films

## Spin Reorientation Transition (SRT)

$$H = -J \sum S_i \cdot S_j - K \sum S_{iz}^2 \quad \text{What will happen at } K \approx 0 ?$$



- Spin-orbit interaction gives  $-KS_z^2$  if  $z$  symmetry is broken.
- Dipolar interaction gives  $2\pi M^2 S_z^2$ .

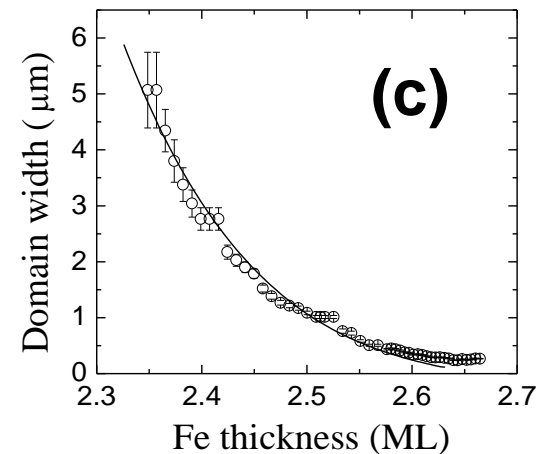
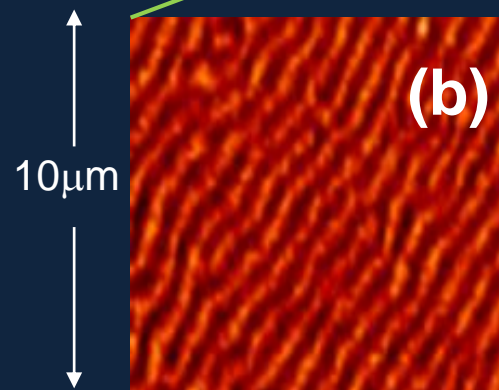
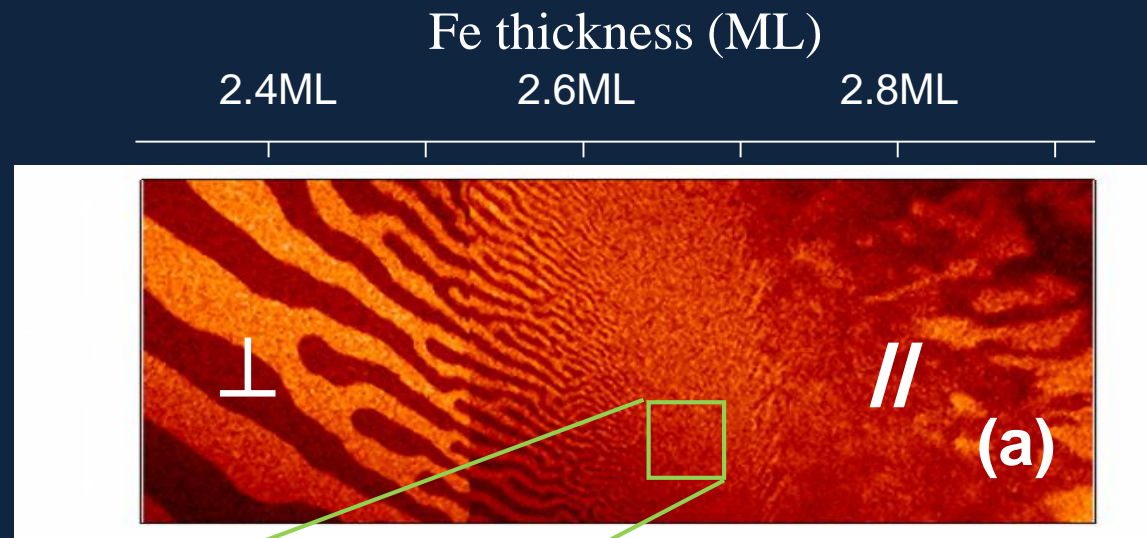
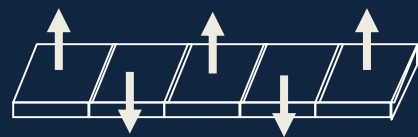
For example,  $E = -Ku_z^2 = -(K_S/d - 2\pi M^2)S_z^2$

Changing  $d$  to tune  $K$  around zero.



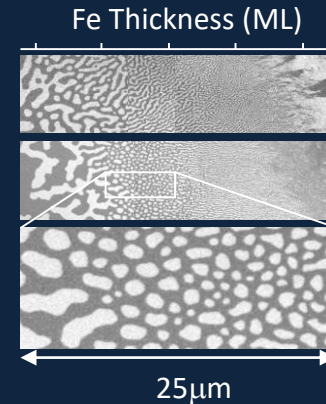
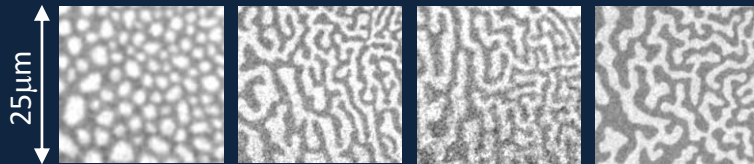
# Stripe Phase

- Y. Z. Wu, C. Won, A. Scholl, A. Doran, H. W. Zhao, X. F. Jin, and Z. Q. Qiu, *Phys. Rev. Lett.* **93**, 117205 (2004).

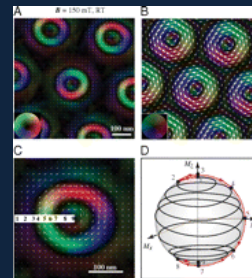
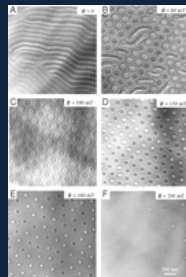


# Bubbles and Skyrmions

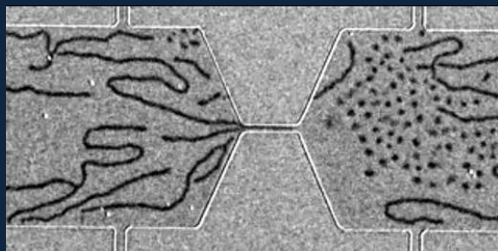
- J. Choi, J. Wu, C. Won, Y. Z. Wu, A. Scholl, A. Doran, T. Owens, and Z. Q. Qiu, “Magnetic bubble domain phase at the spin reorientation transition of ultrathin Fe/Ni/Cu(001) film”, *Phys. Rev. Lett.* **98**, 207205 (2007).



- Xiuzhen Yu *et al.* “Magnetic stripes and skyrmions with helicity reversals.” *Proceedings of the National Academy of Sciences* **109**, 8856-8860 (2012).

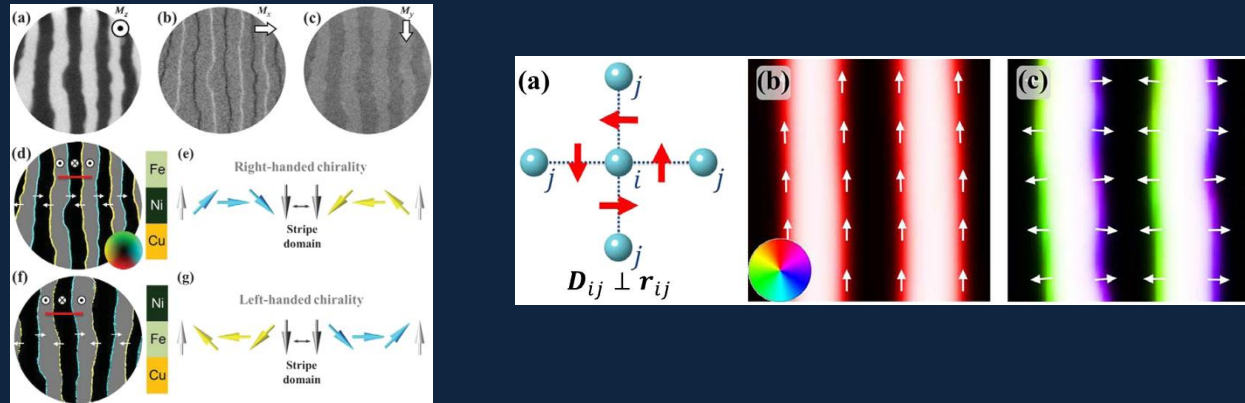


- Wanjun Jiang *et al.* “Blowing magnetic skyrmion bubbles.” *Sciences* **349**, 283-286 (2015).

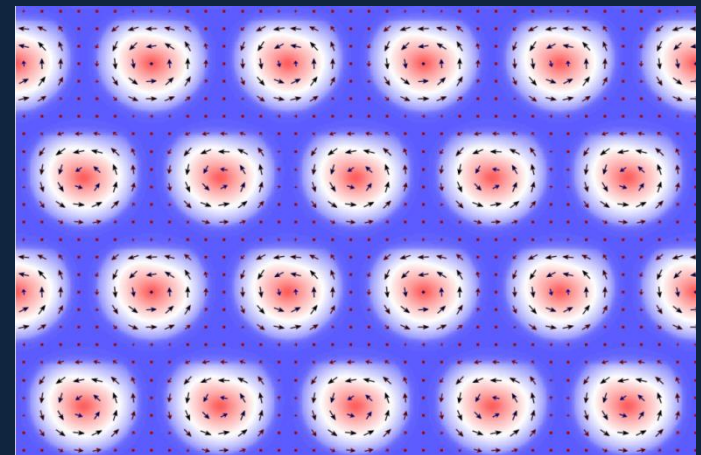
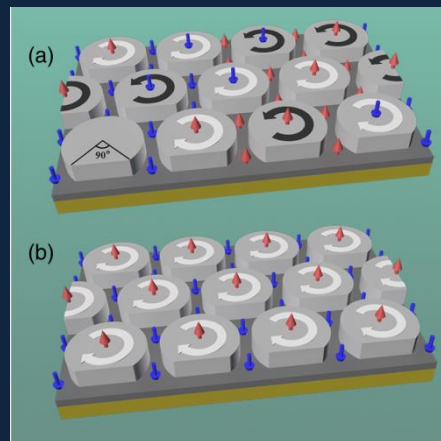
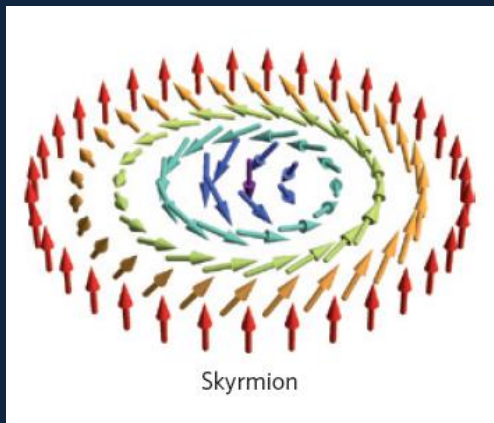




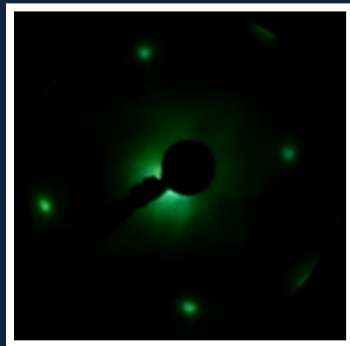
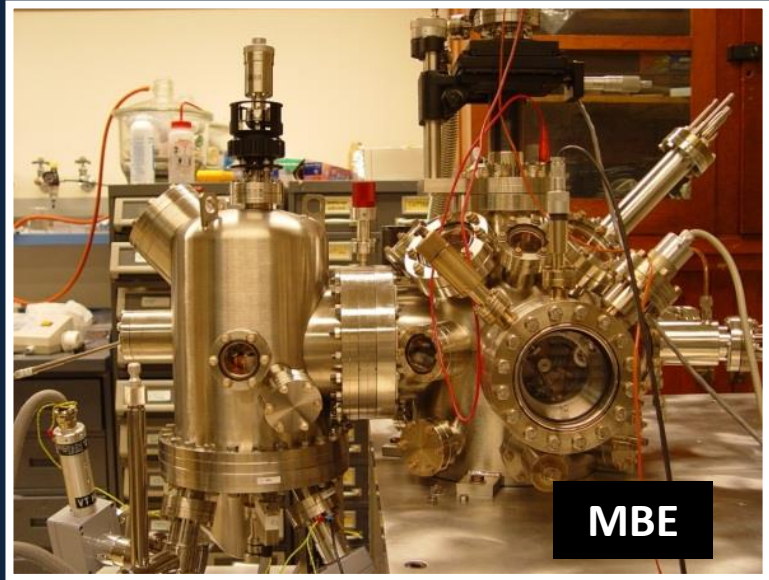
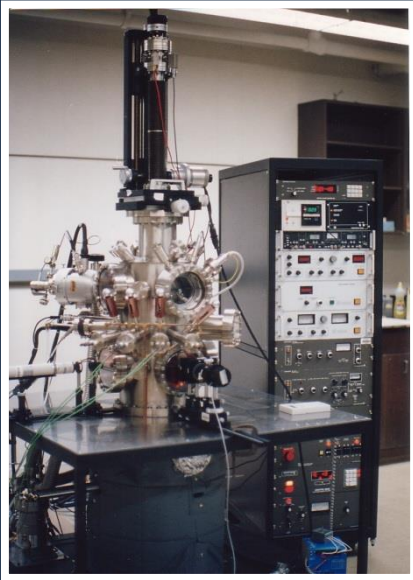
- G. Chen *et al.* “Novel chiral magnetic domain wall structure in Fe/Ni/Cu(001) films”, *Phys. Rev. Lett.* **110**, 177204 (2013).



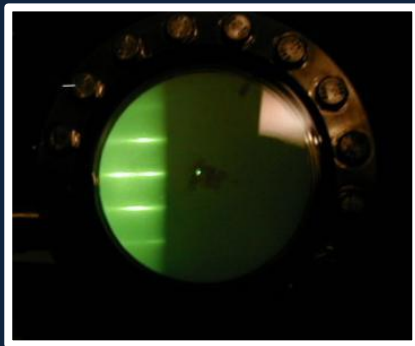
- L. Sun, *et al.* “Creating an Artificial Two-Dimensional Skyrmion Crystal by Nanopatterning.” *Phys. Rev. Lett.* **110**, 167201 (2013).



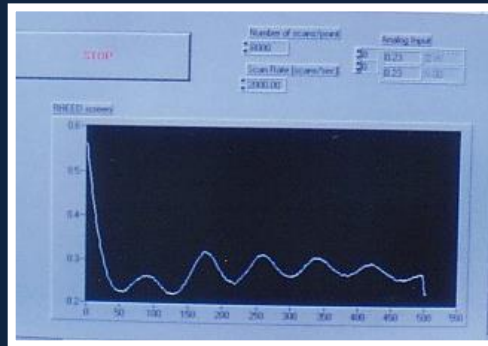
# MBE grown magnetic thin films



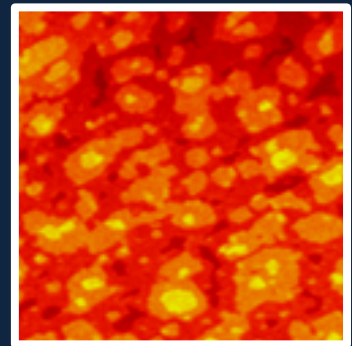
LEED



RHEED

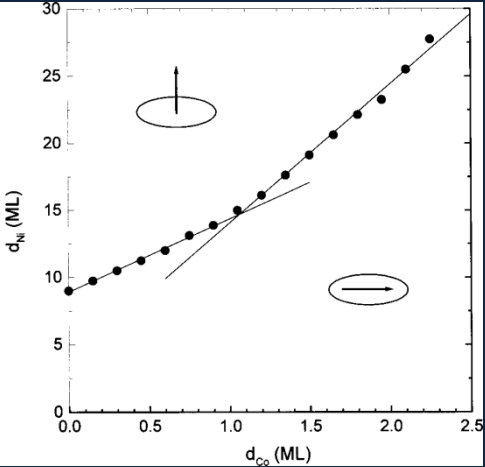
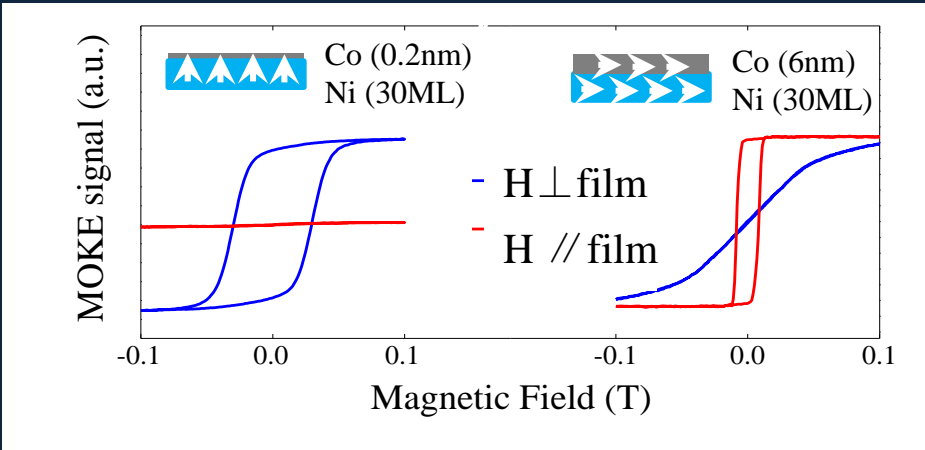
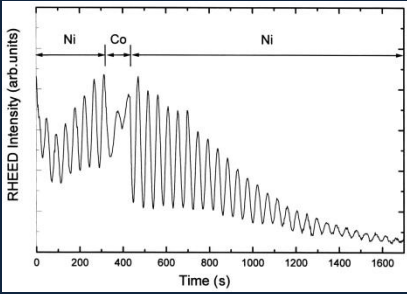
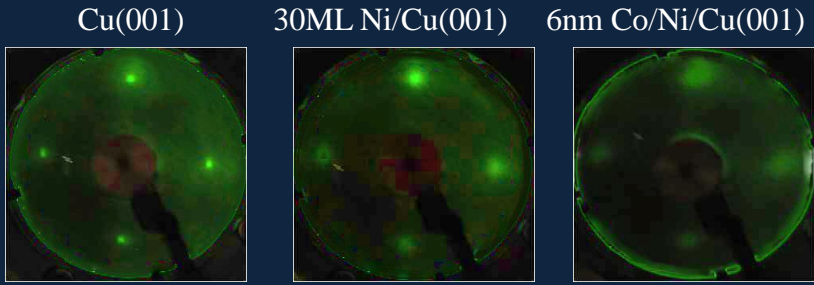


RHEED Oscillations

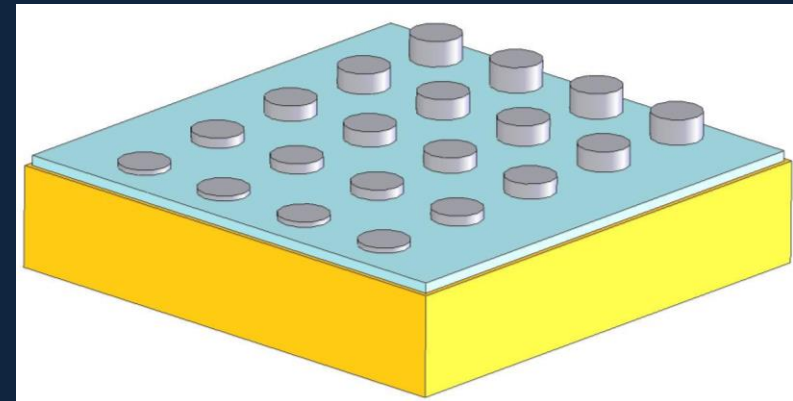
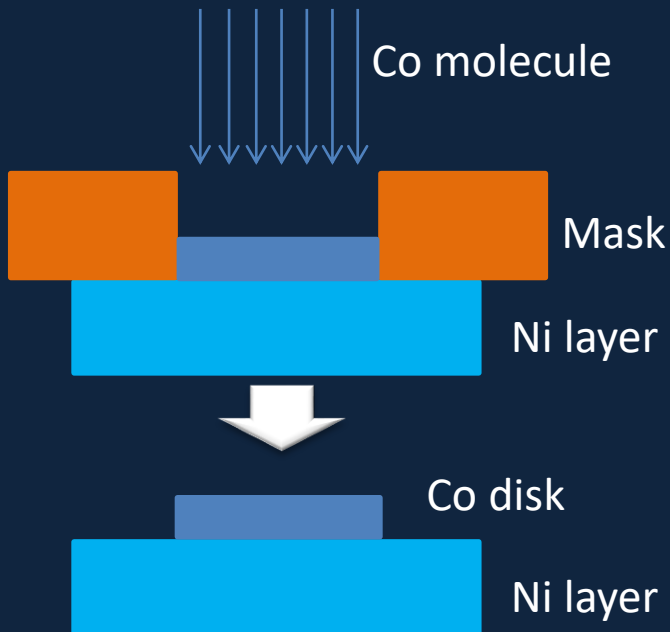
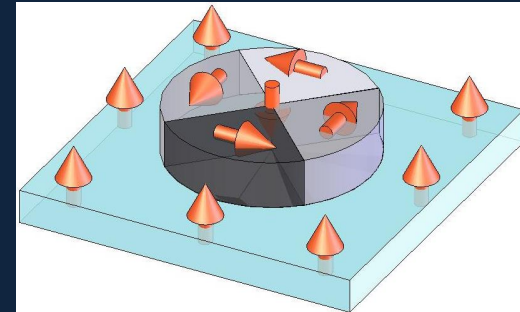
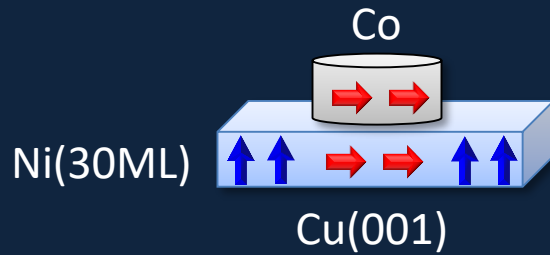


STM

# Co/Ni/Cu(001)



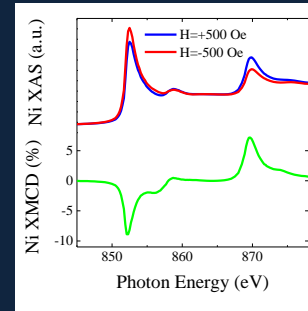
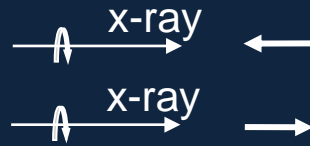
# Co(disk)/Ni(30ML)/Cu(001)



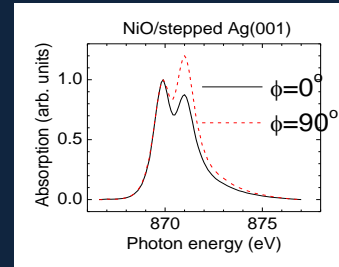
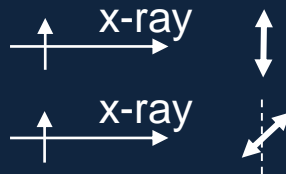


# X-ray Magnetic Dichroism — Element-resolved measurement

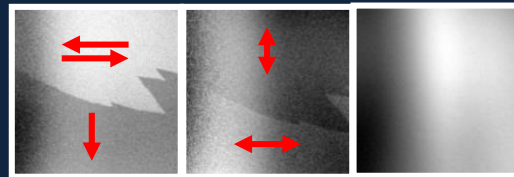
**XMCD**



**XMLD**



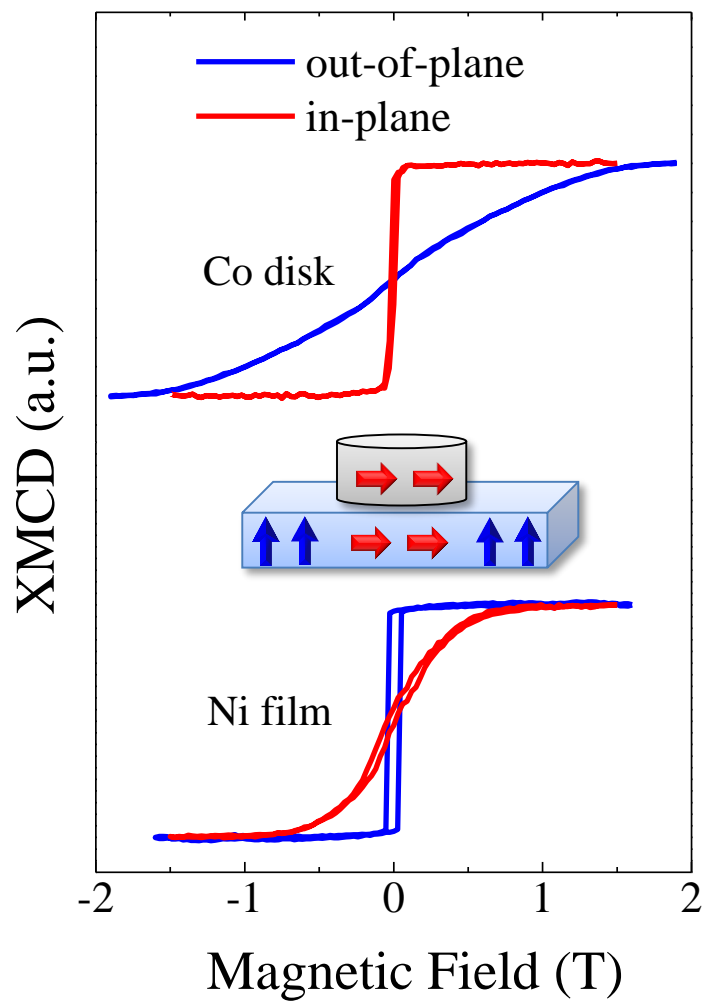
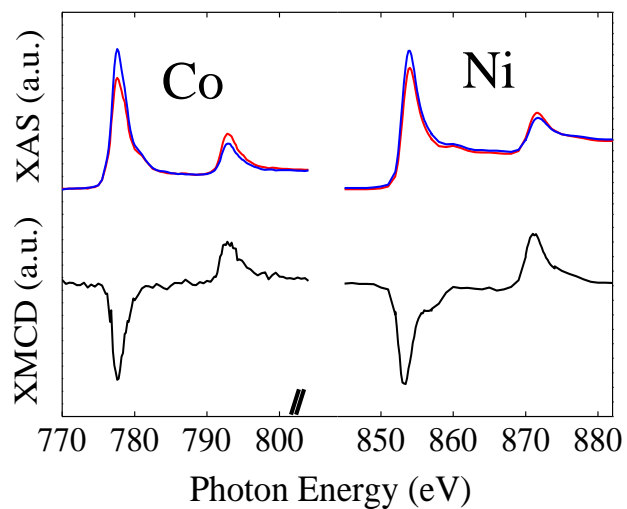
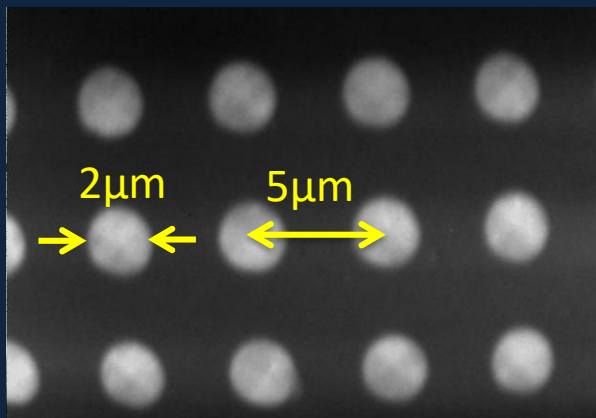
**PEEM**



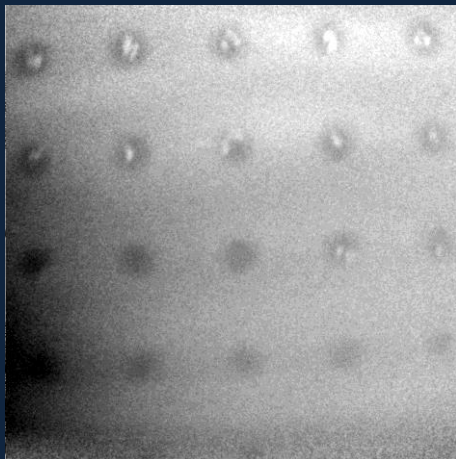
NiO(14ML)/Fe(15ML)/Ag(001)

Fe XMCD Ni XMLD Ni XMCD

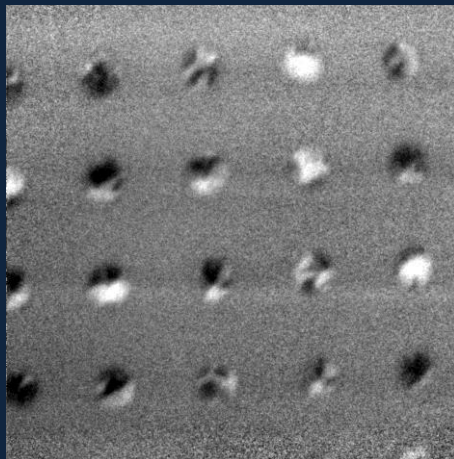




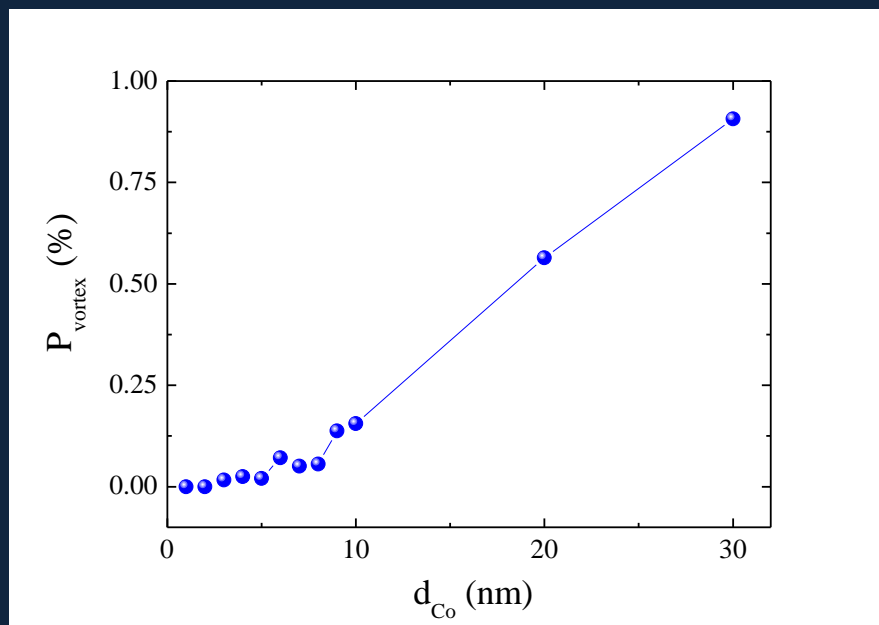
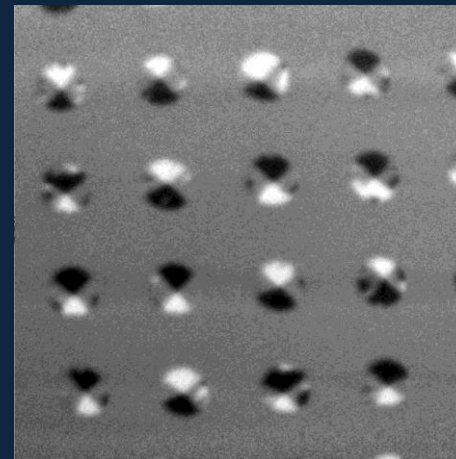
$d_{Co} = 1\text{ nm}$



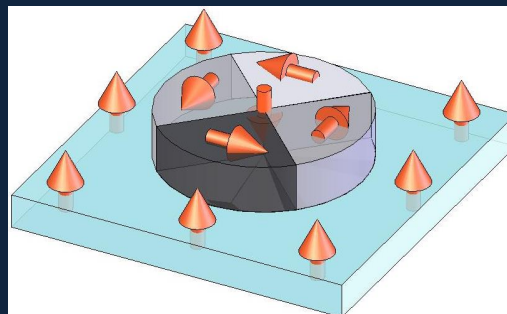
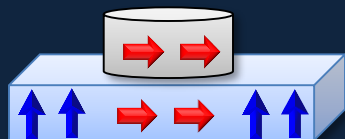
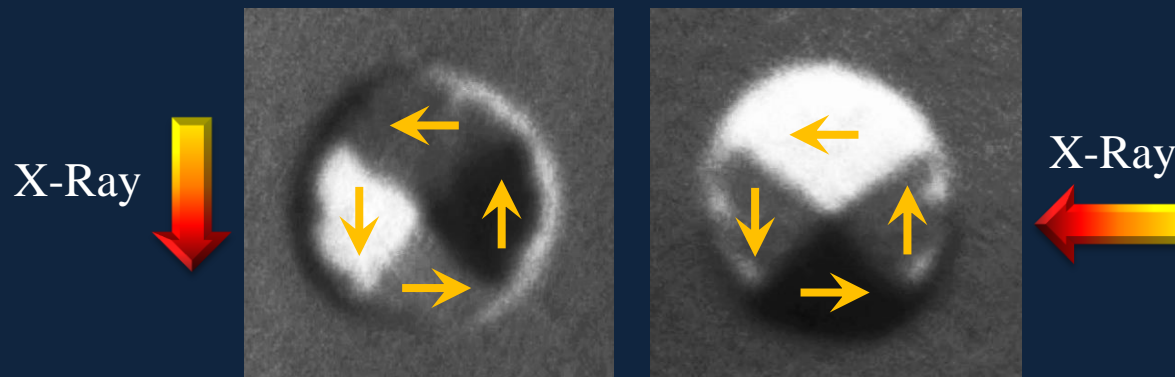
$d_{Co} = 3\text{ nm}$



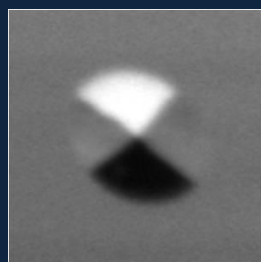
$d_{Co} = 30\text{ nm}$



# Co(disk)/Ni(30ML)/Cu(001)



## Ni(5ML)/Co(disk)/Cu(001)

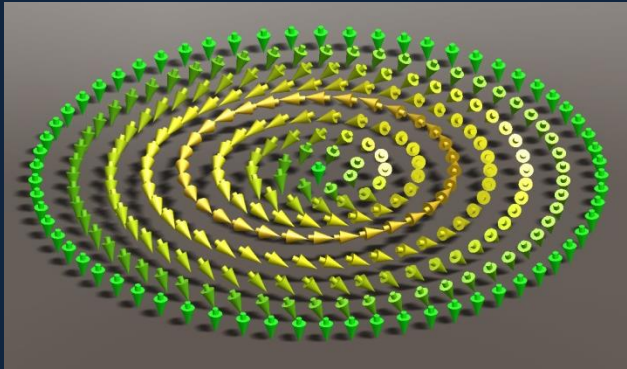


Co domain

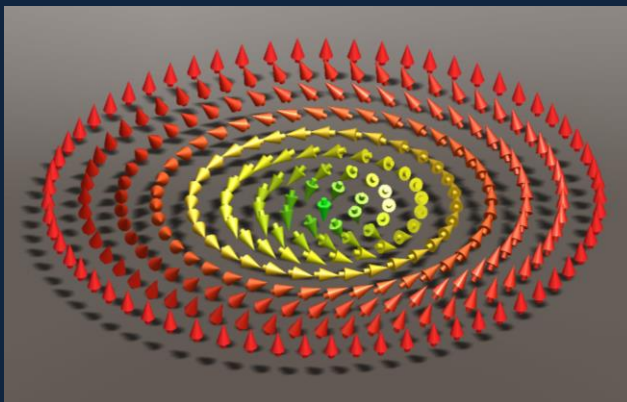
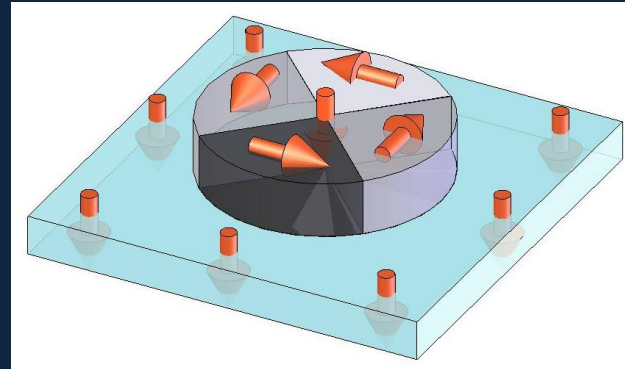


Ni domain

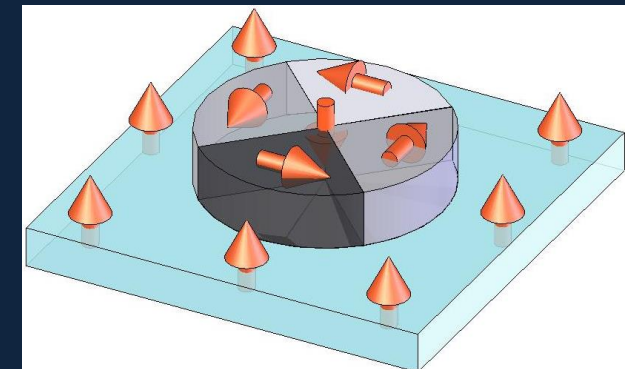
# Preparing N=0 and N=1 Skyrmions



N=0



N=1



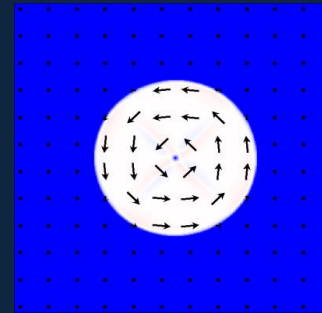
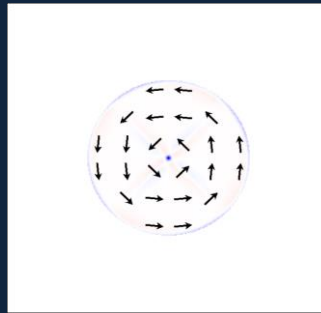
N=0 State:  $H=20,000\text{Oe}$  and then  $H=0$

N=1 State:  $H=-800\text{Oe}$  to N=0 state and then  $H=0$

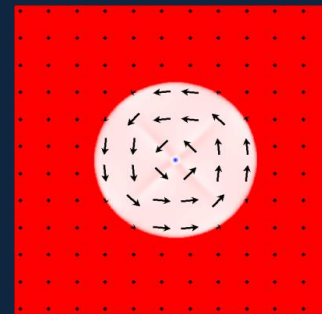
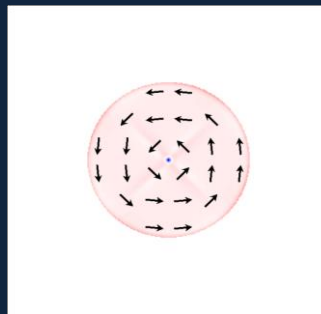
Co

Ni

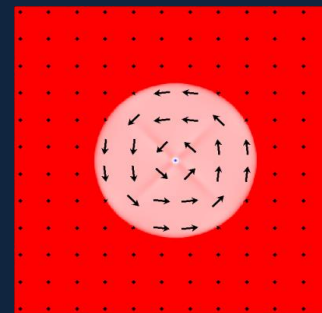
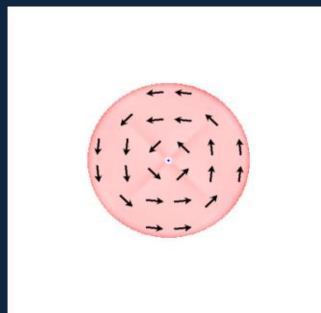
H= 0 Oe



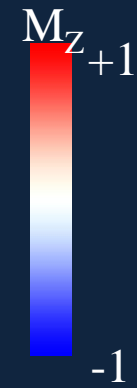
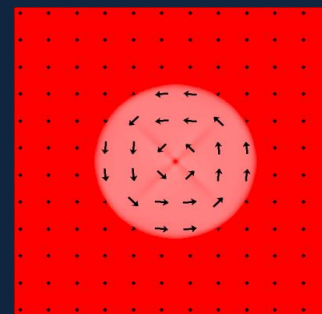
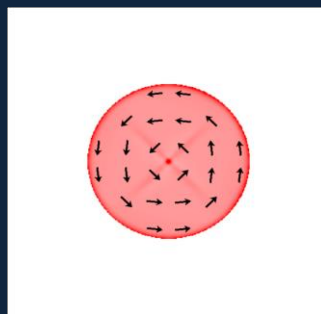
1000 Oe



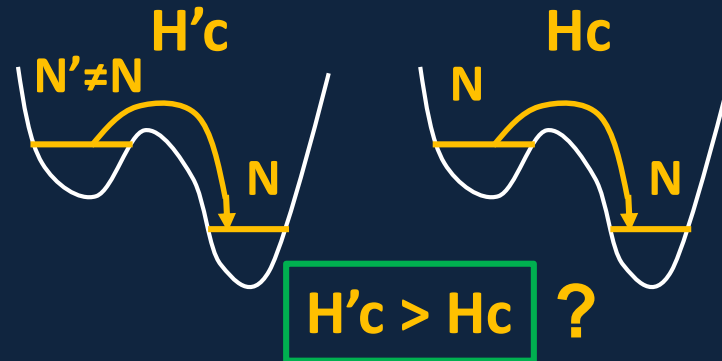
3000 Oe



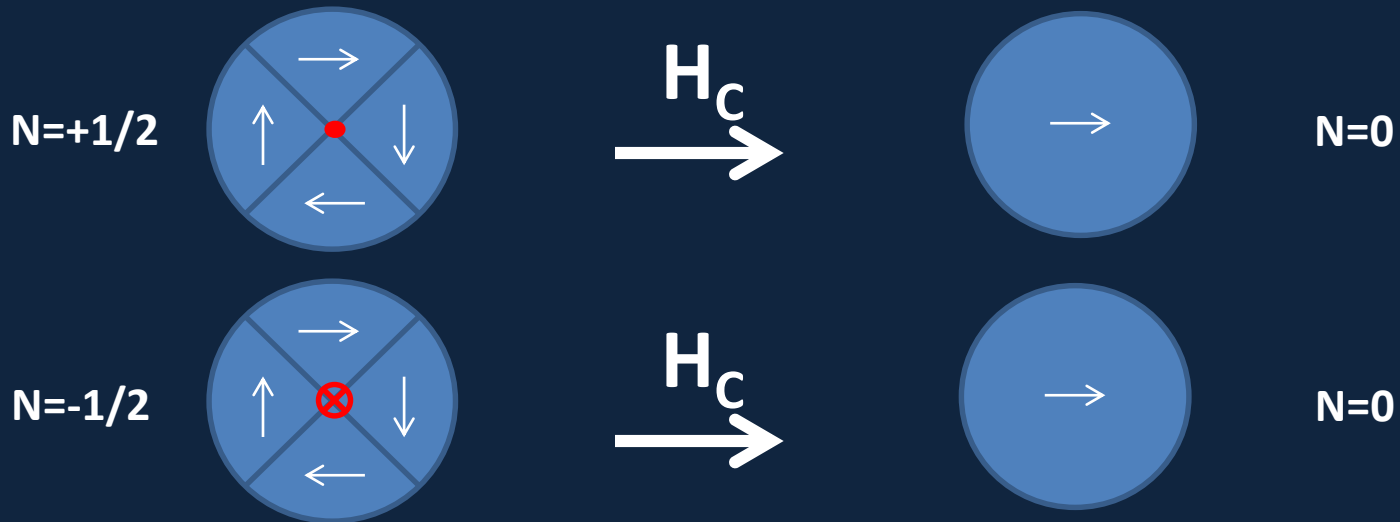
6000 Oe



# Searching for the topological effect of Skyrmions



## Vortex core annihilation



Apply an in-plane field to saturate the vortex state into a single domain state.

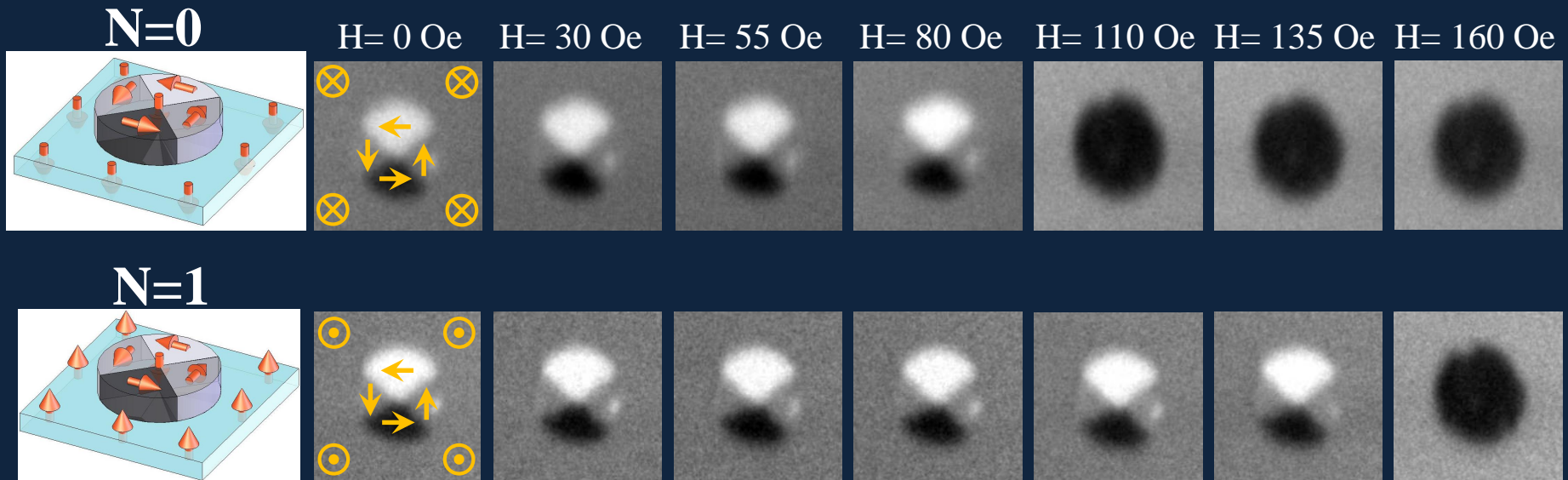


## Skyrmion core annihilation



Apply an in-plane field to saturate the middle vortex into a single domain state.



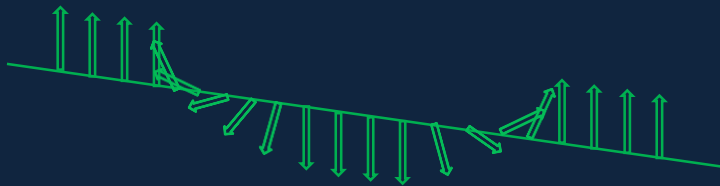
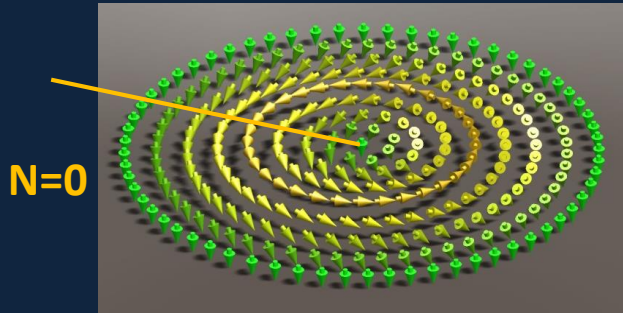
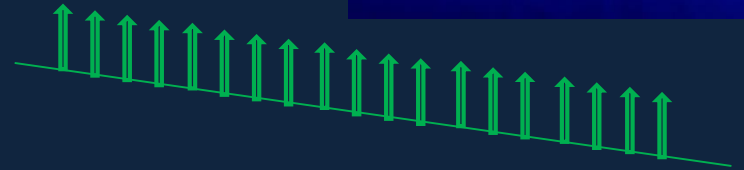


The critical field for N=1 is greater than N=0 state, suggesting a topological effect in the Skyrmion core annihilation process.

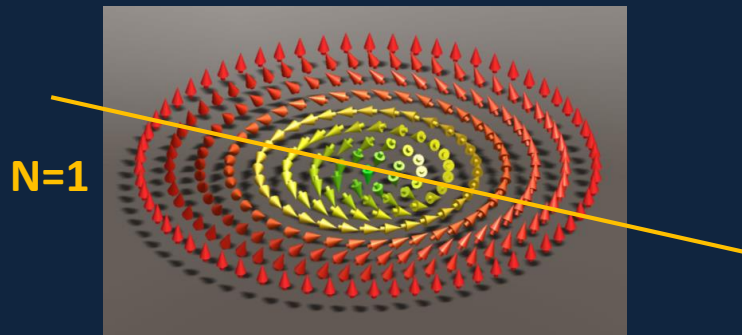
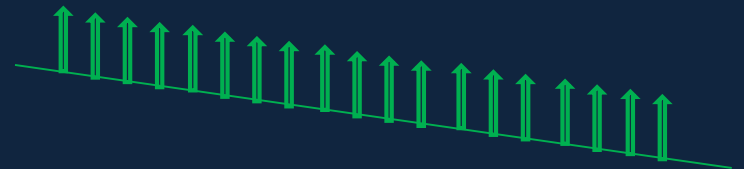
# Topology of a domain wall



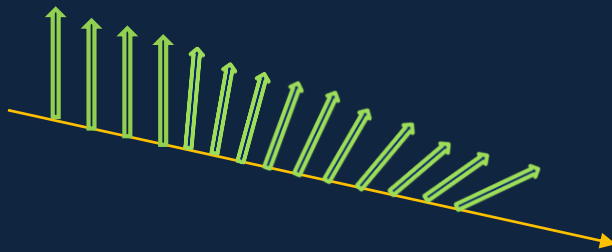
=



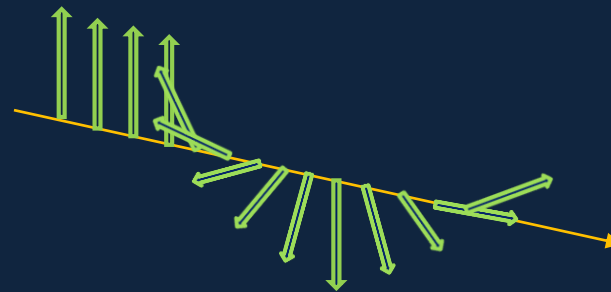
≠



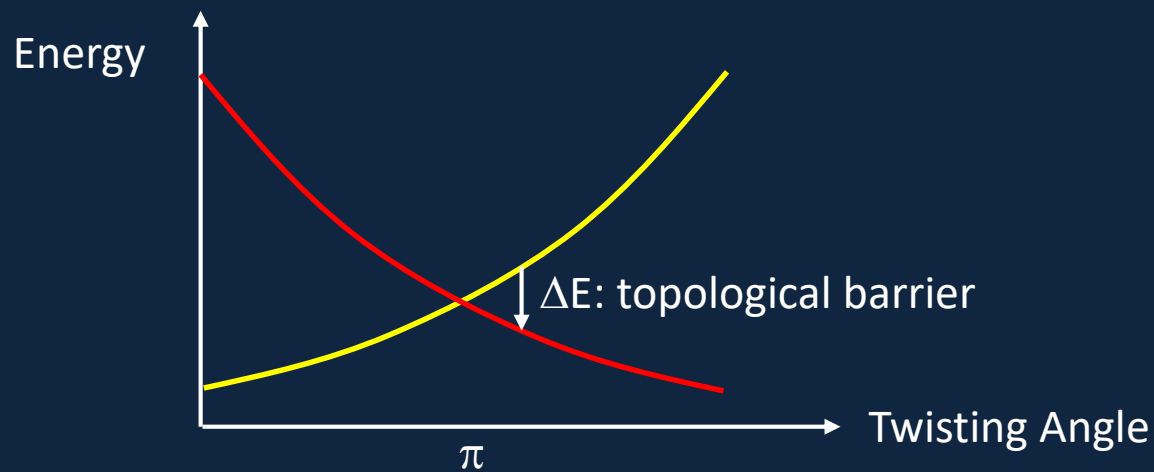
# Breaking the topology of the domain wall



Left-handed



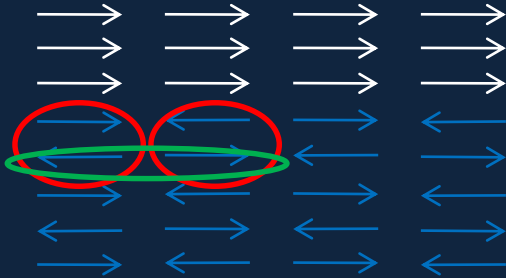
Right-handed



# Exchange bias

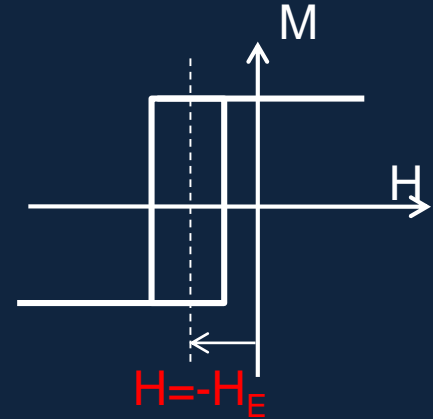
$$\sum_i \vec{S}_{F,i} \neq 0$$

**FM**



$$\sum_i \vec{S}_{AF,i} = 0$$

**AFM**



W. H. Meiklejohn and C. P. Bean, Phys. Rev. **102**, 1413 (1956).

## Mauri's Model

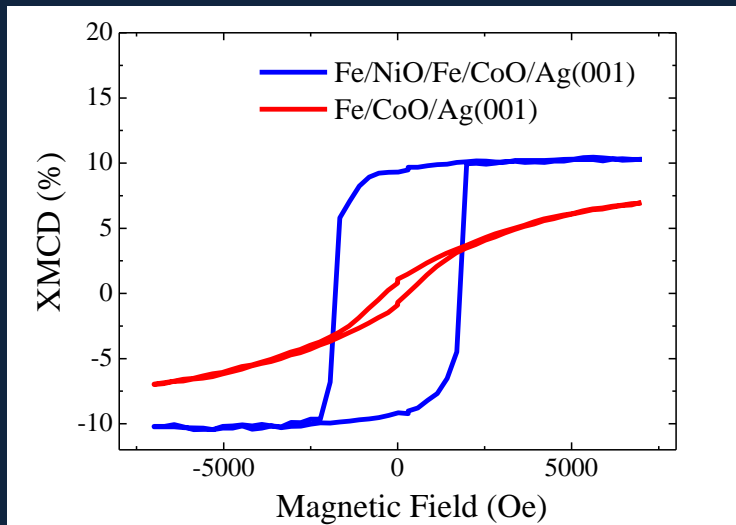
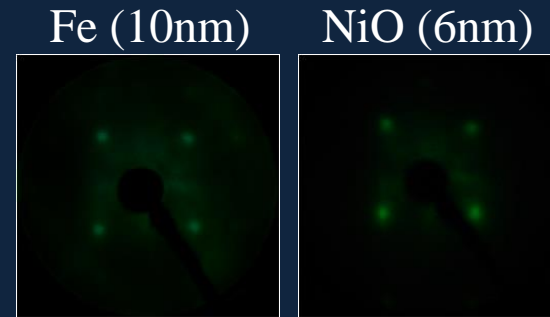
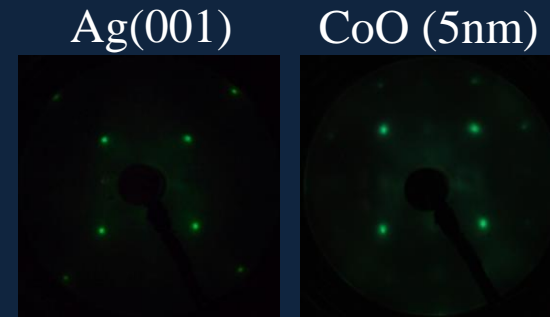
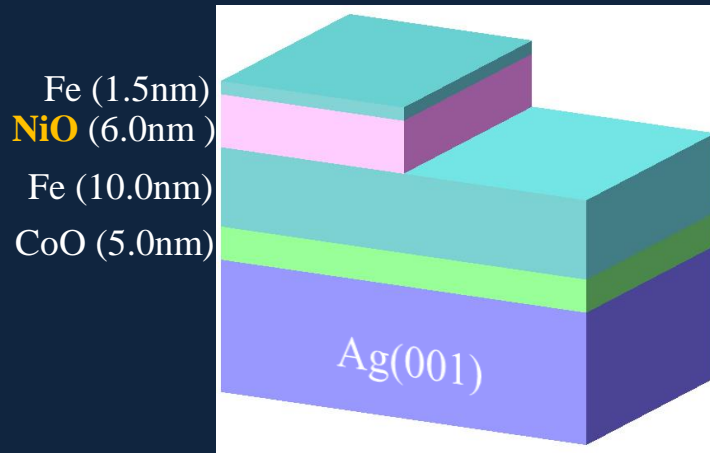
The FM rotation winds up an AFM domain wall, leading to the exchange bias.

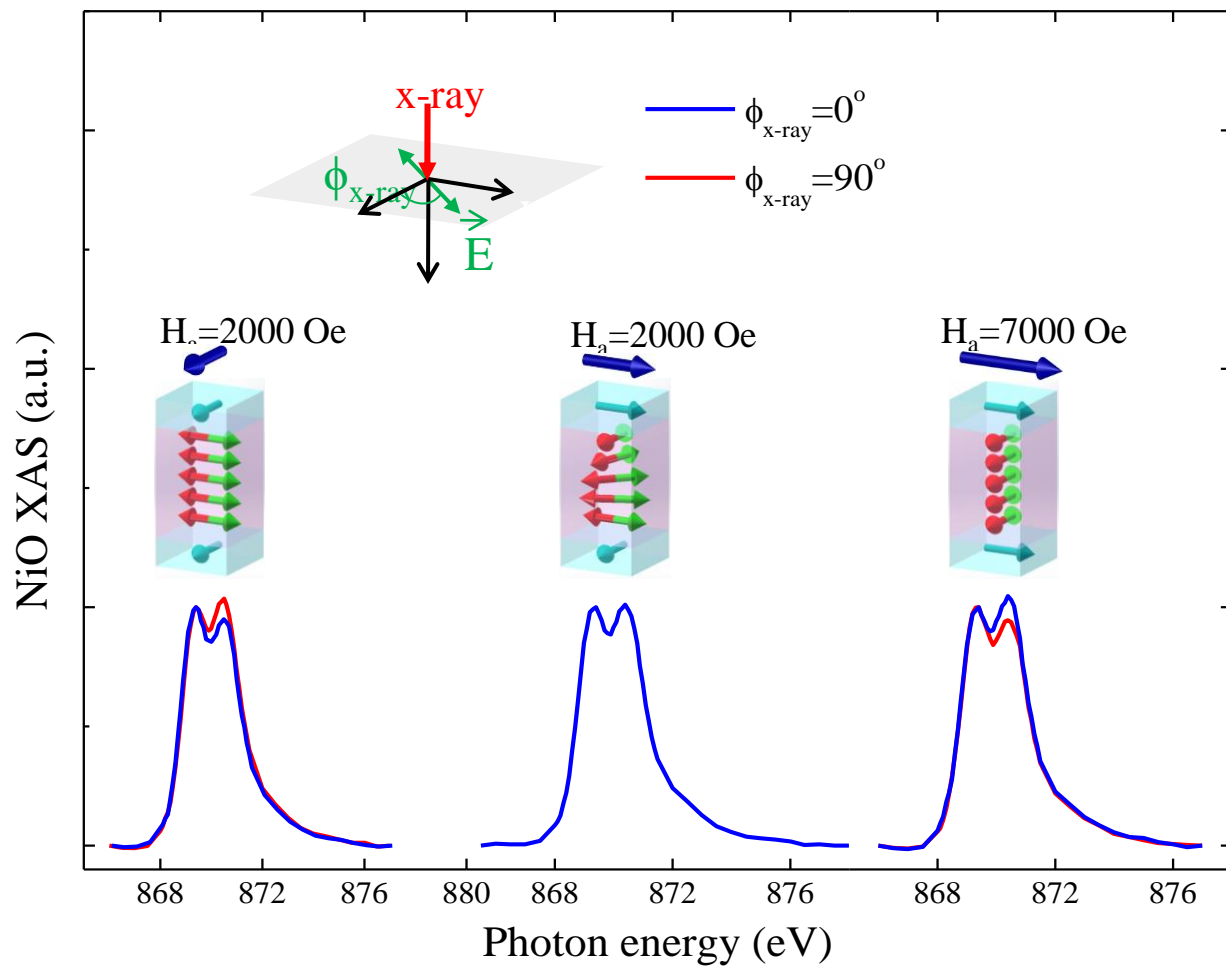


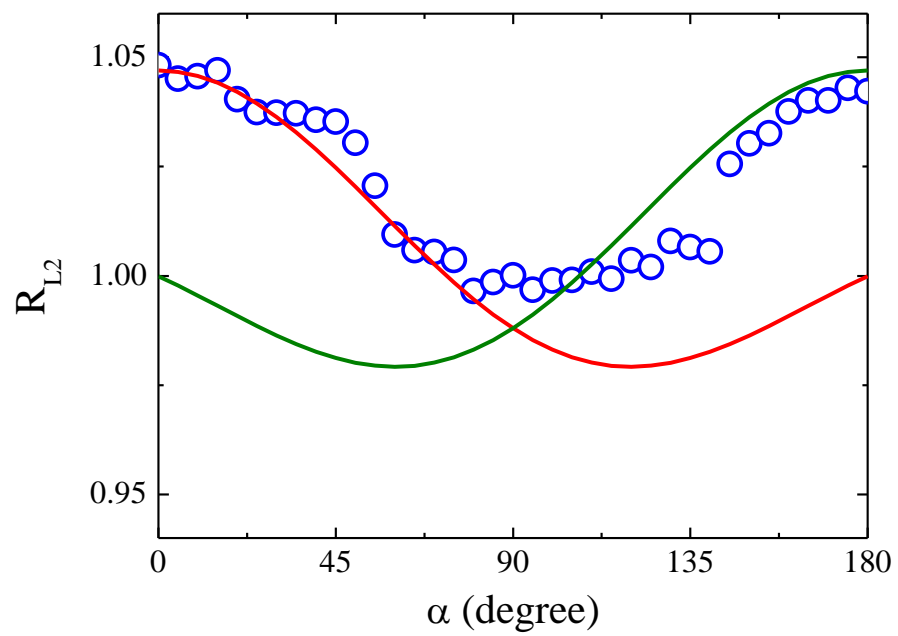
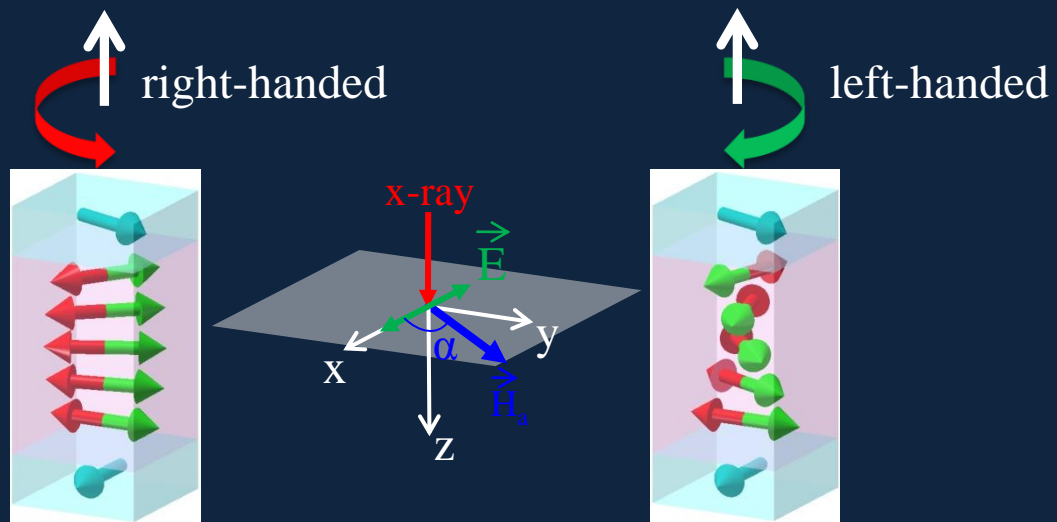
**Winding up**

**Fixed**

# Chirality switching of AFM NiO domains







# Summary

- **Artificial Skyrmions can be created by embedding a magnetic vortex into a perpendicular magnetized background.**
- **The Skyrmion number can be controlled by the relative orientation of the vortex core to the perpendicularly magnetized background.**
- **Artificial Skyrmion exhibits a topological effect in the core annihilation process.**
- **NiO domain wall exhibits a topological effect by switching its chirality.**



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